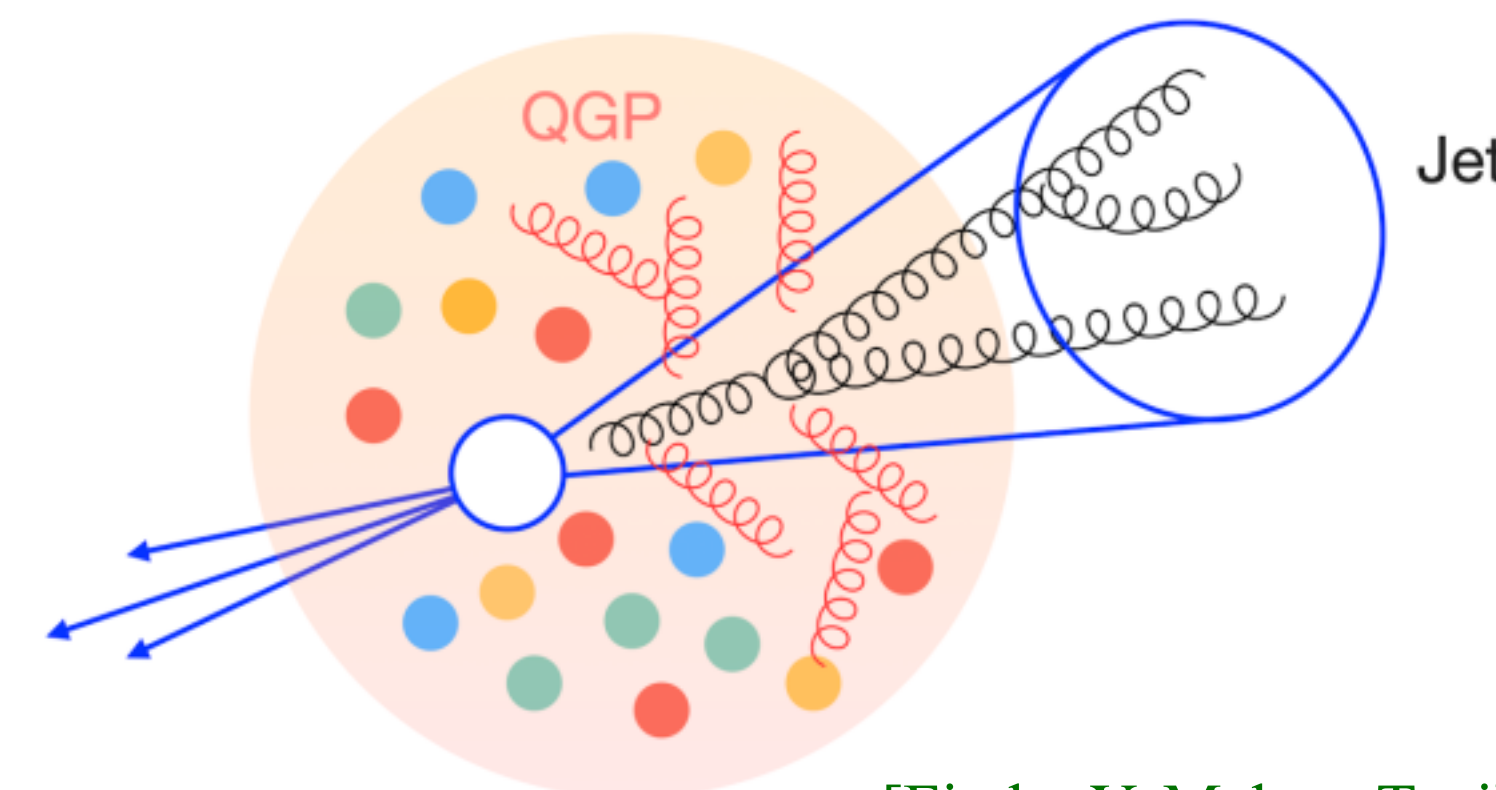
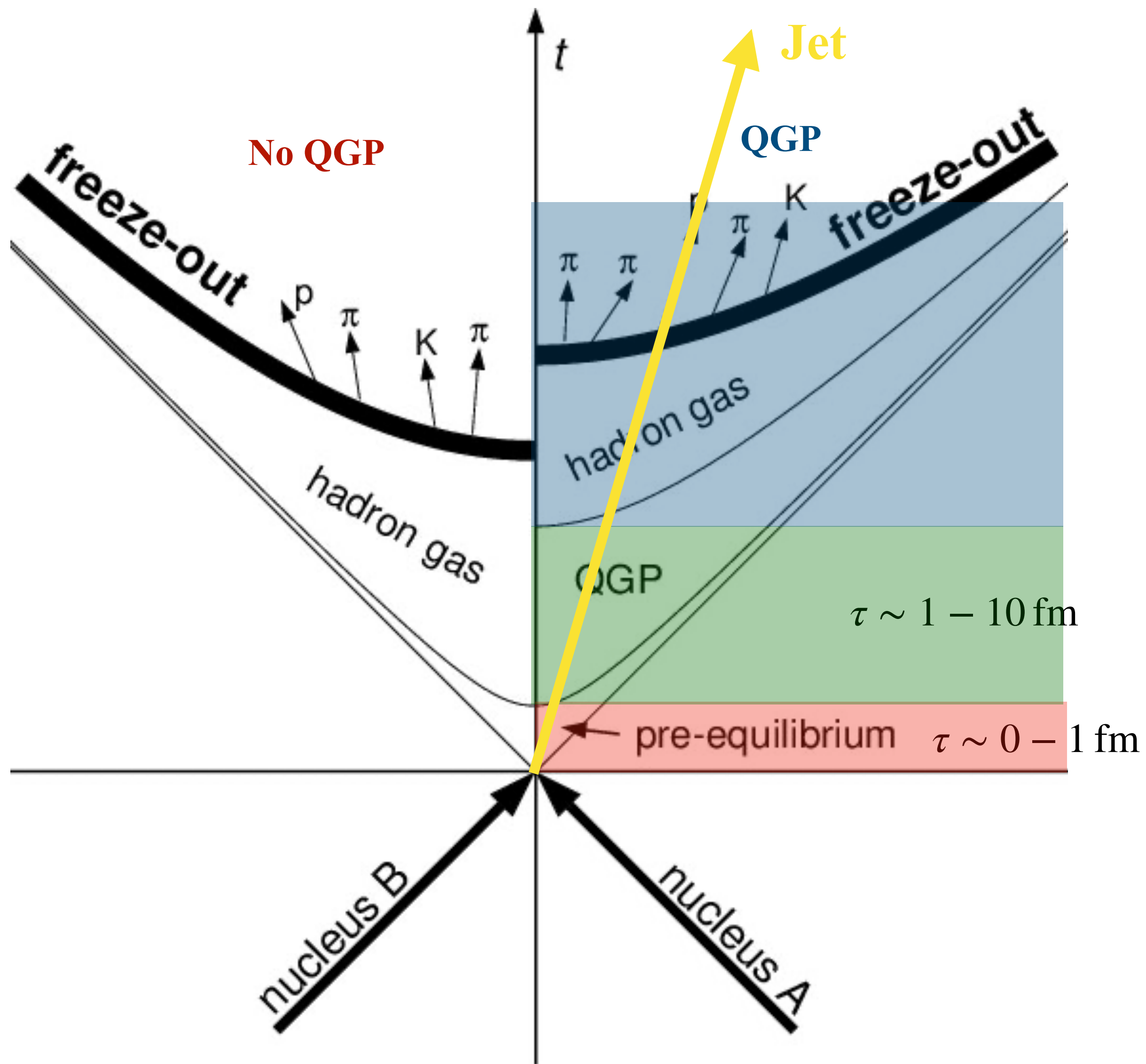


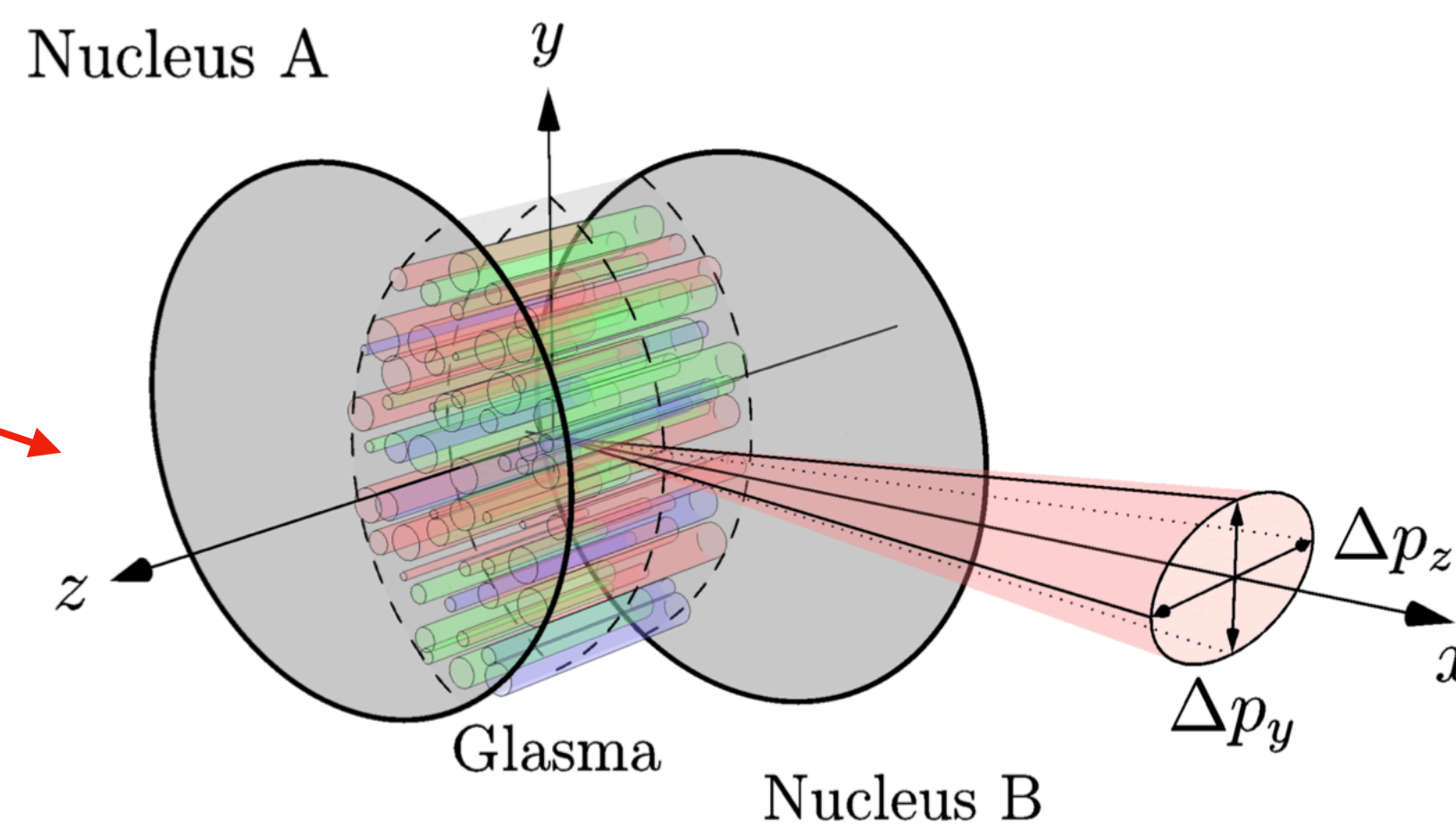
Jet anisotropies in jet quenching

4th September 2024, Muenster

João Barata, BNL



[Fig by Y. Mehtar-Tani]



[A. App, D. I. Muller, D. Schuh, 2009.14206]

The Landscape



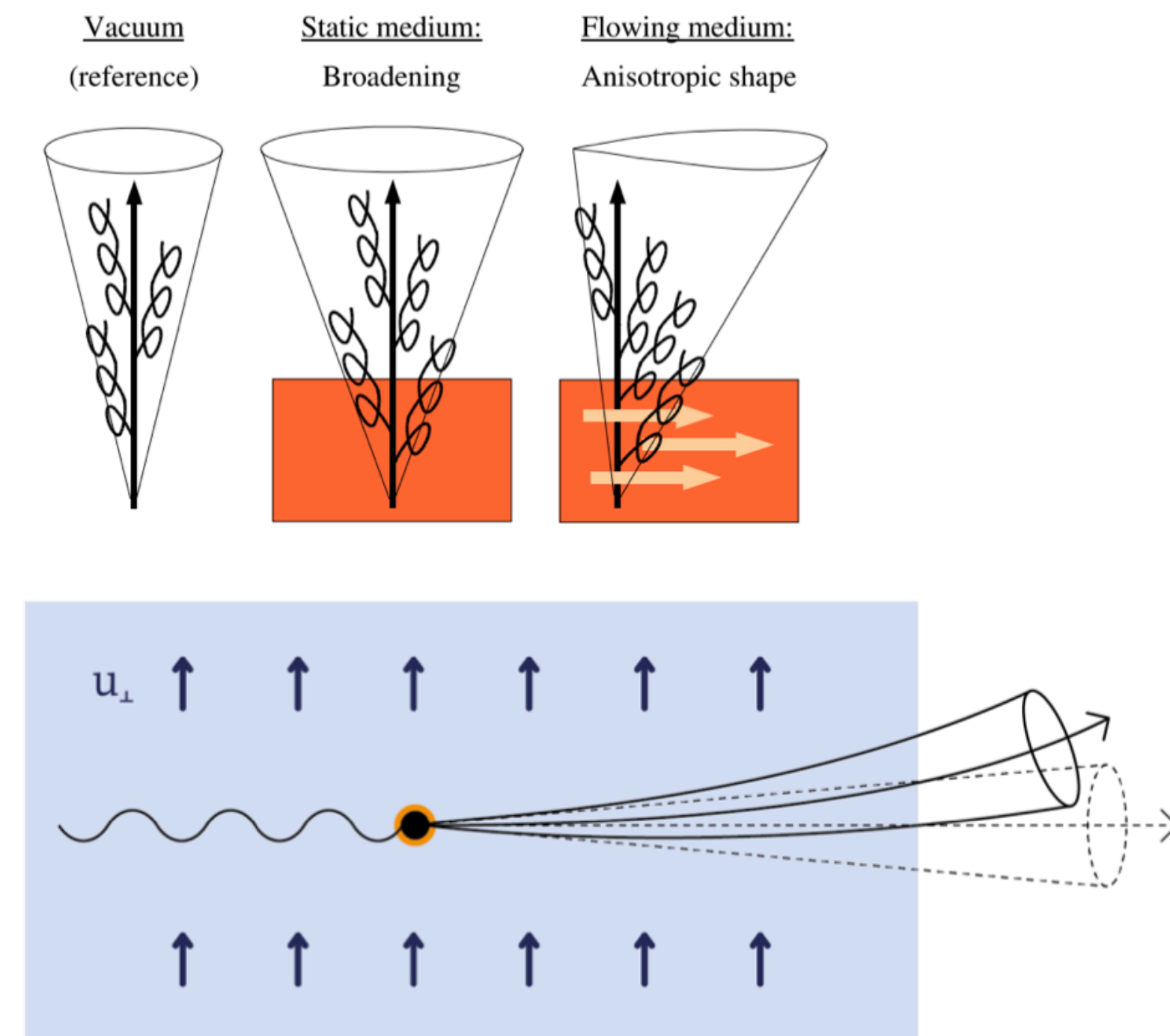
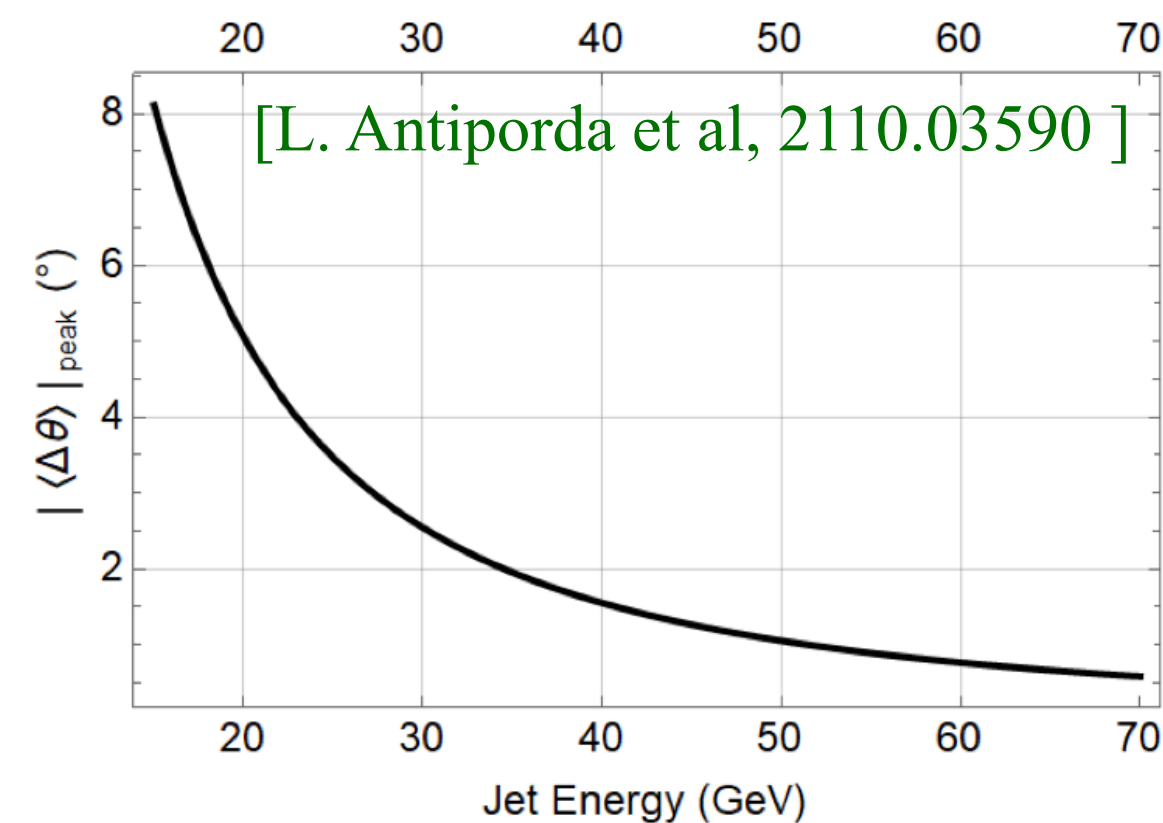
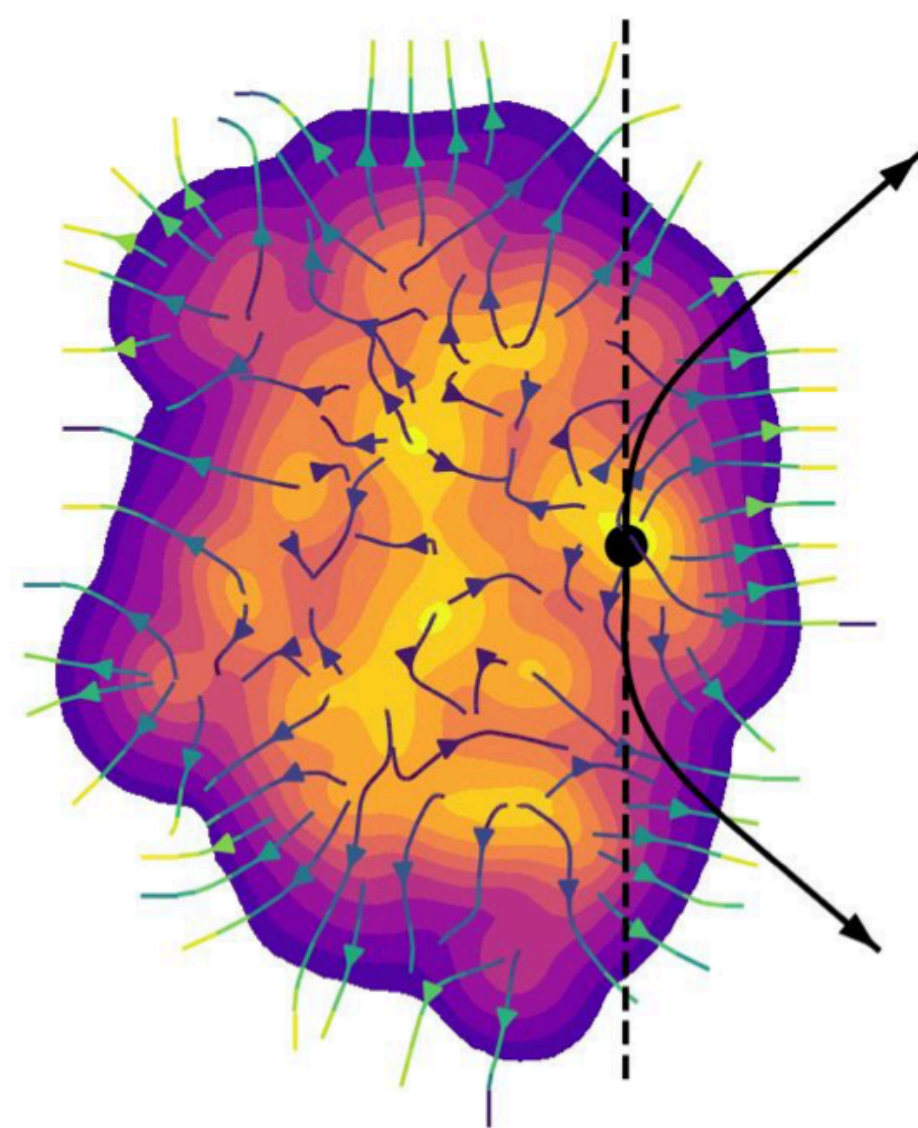
The Landscape

Jet evolution in hydrodynamical phase

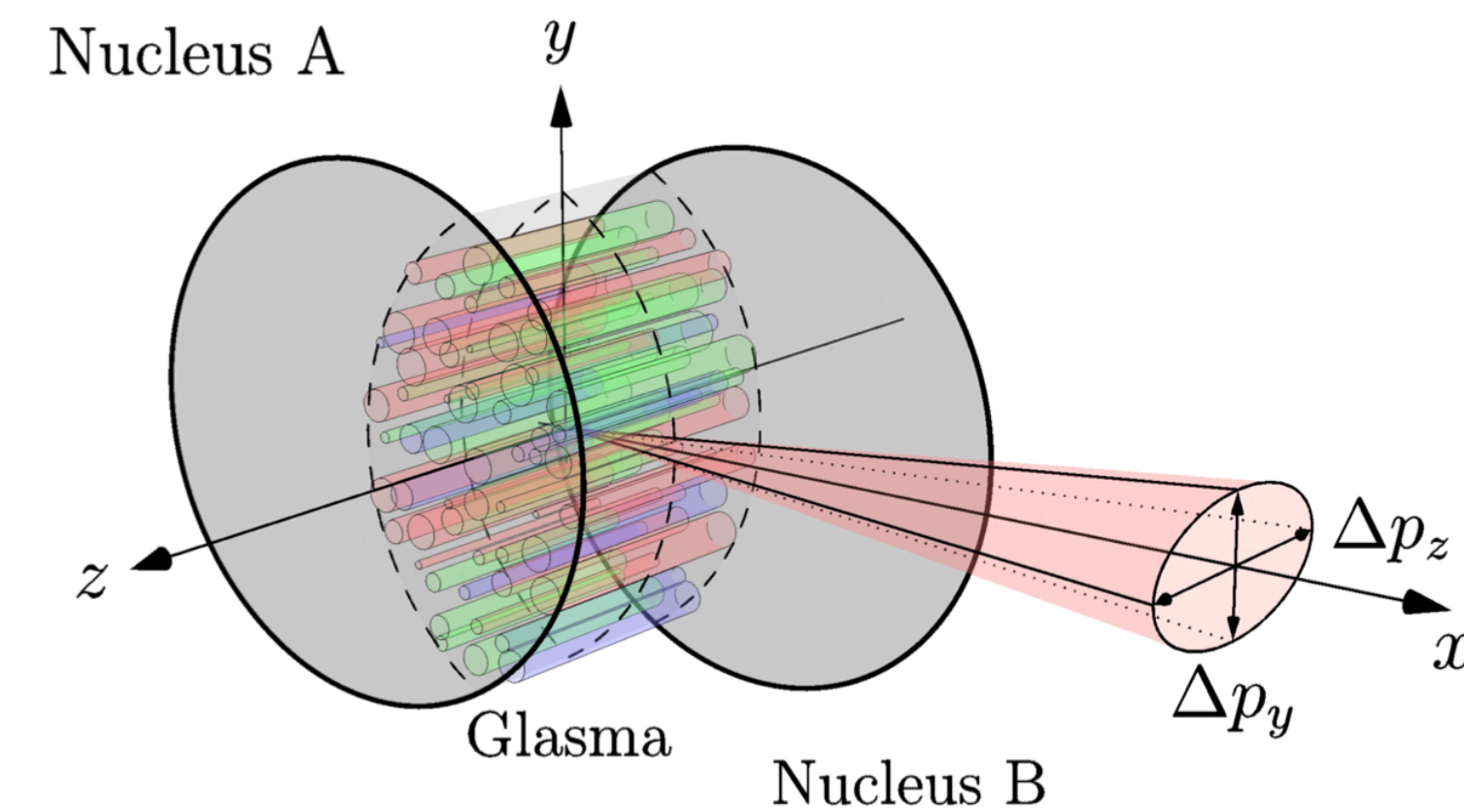
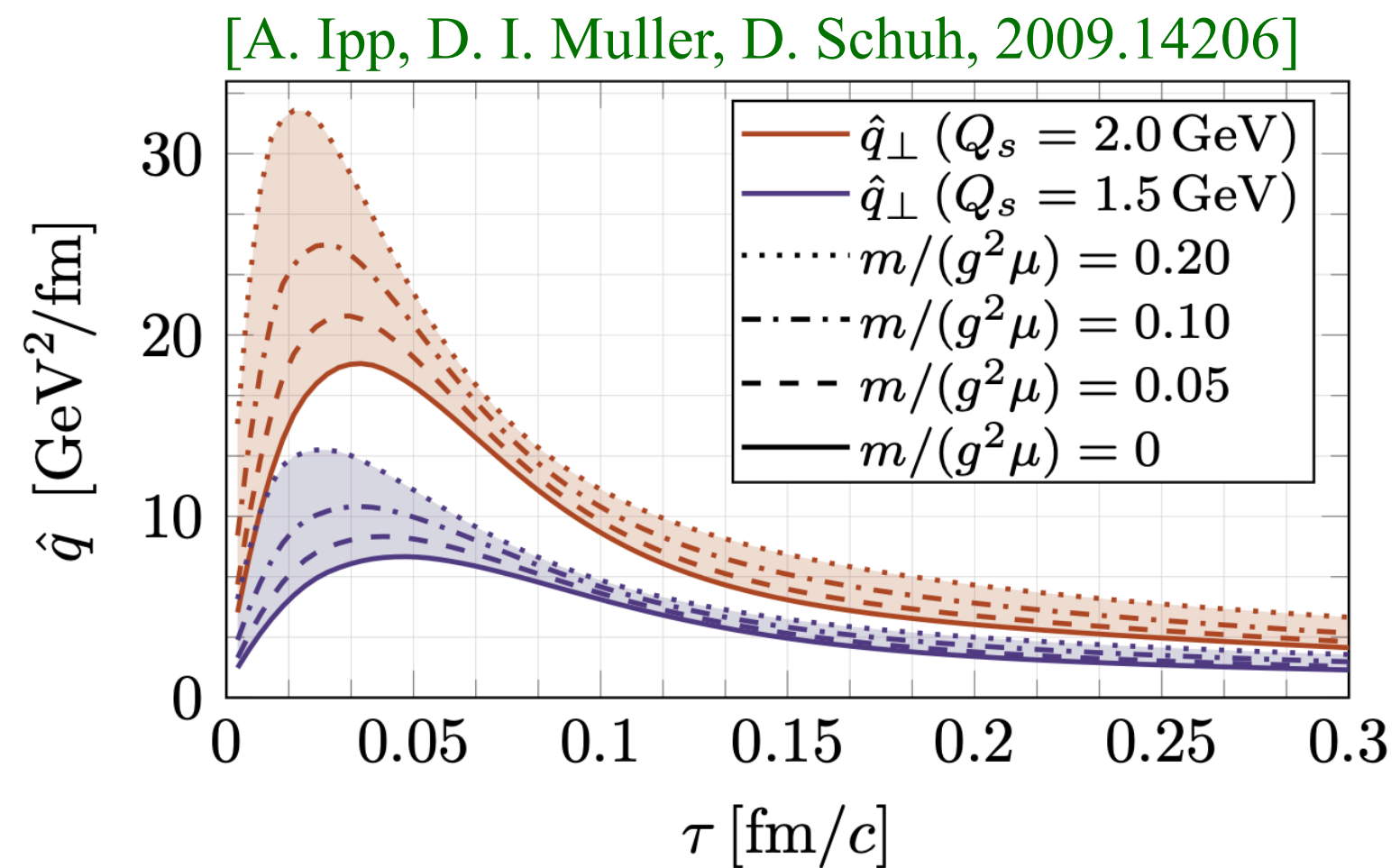
Flowing matter: 2104.09513 [N=1], 2207.07141 [Resummation], 2406.14628 [Gluon radiation], 2309.00683 [Flowing anisotropic matter], ...

Matter gradients: 2104.09513 [N=1], 2202.08847 [Resummation], 2210.06519 [Kinetic Th.], 2304.03712 [Gluon Radiation], 2204.05323 [Broadening]

Jet observables/Pheno: 2110.03590 [Jet drift], 2308.01294 [Jet substructure], DREENA, ...



The Landscape



Jet evolution in the early stages

Momentum broadening in early stages: [Works by Avramescu et al](#); [Lindebauer et al](#); [Muller et al](#), ...

Radiative spectrum in “glasma”: [2306.20307 \[Photons\]](#), [2303.03914 \[Gluon branching\]](#), [2407.04774 \[Quark antenna\]](#), [2406.07615 \[Spatial correlations\]](#)

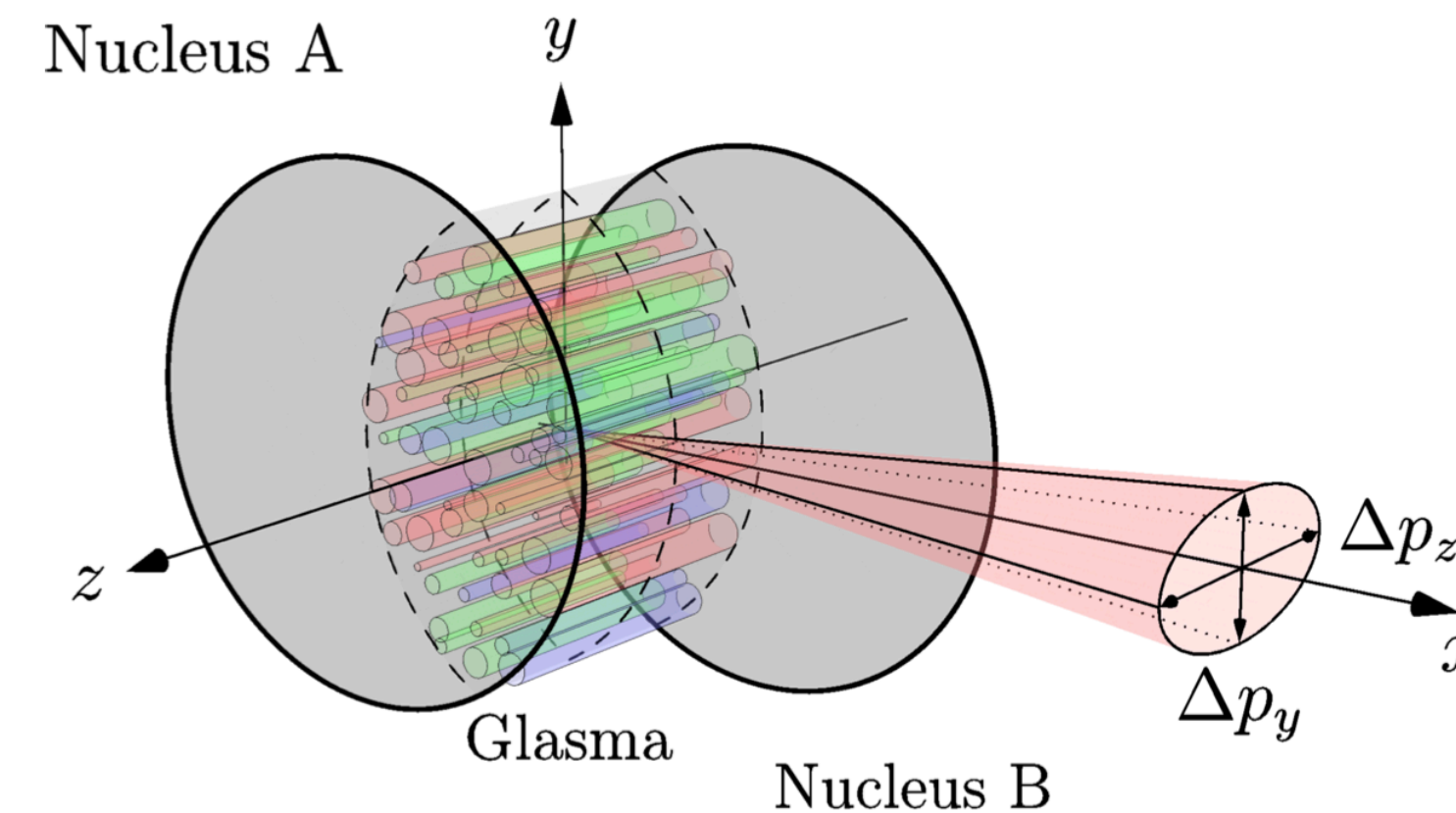
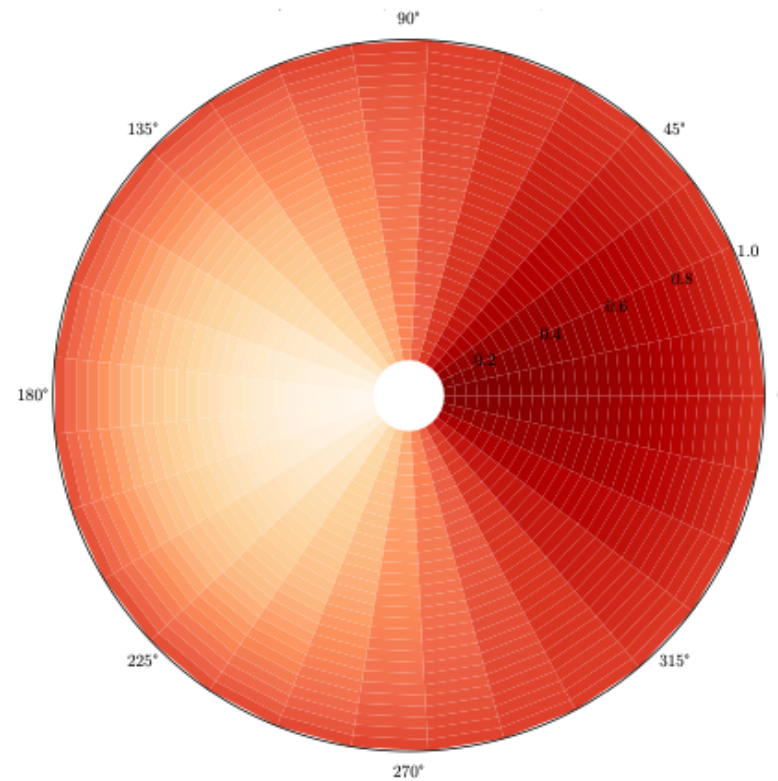
A lot of recent pheno and theory activity, but why is this important ? Can we measure ?

The Questions



The Questions

1) How about existing measurements (e.g. Jet FF, **jet radial profile**)? What insights do they provide ? Do we need new tools ? **Yes**



2) Are there any new (orthogonal) observables that we need to consider ?

A few, but not studied in great detail in theory or pheno; small exp. interest

3) Based on our physics goals, can we already identify a particular set of observables ?

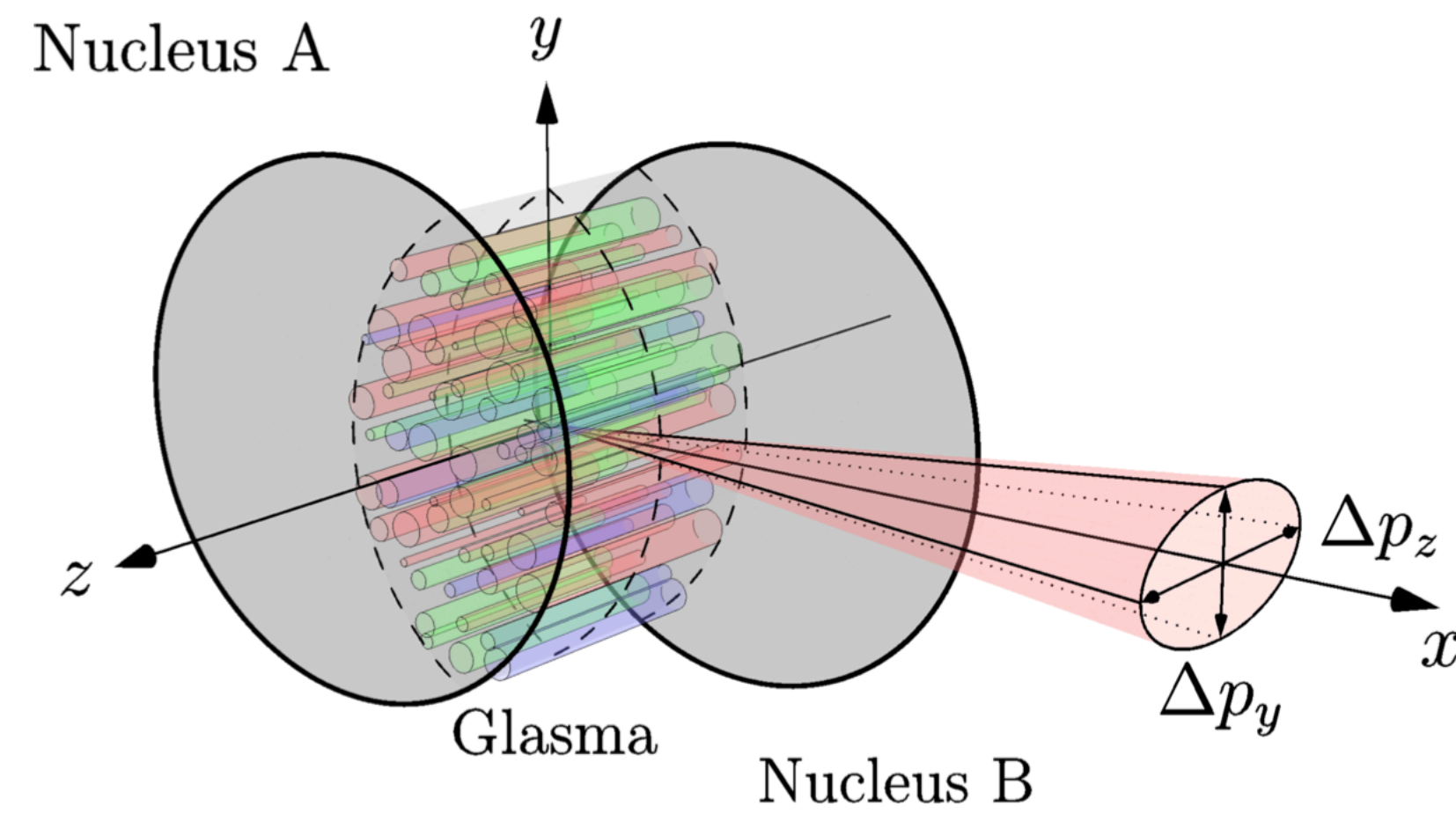
Yes: azimuthally unintegrated jet substructure obs. , spin dependent obs., ...

Why is this hard(er) ? Soft-hard correlation, more differential obs., less theoretical control, ...

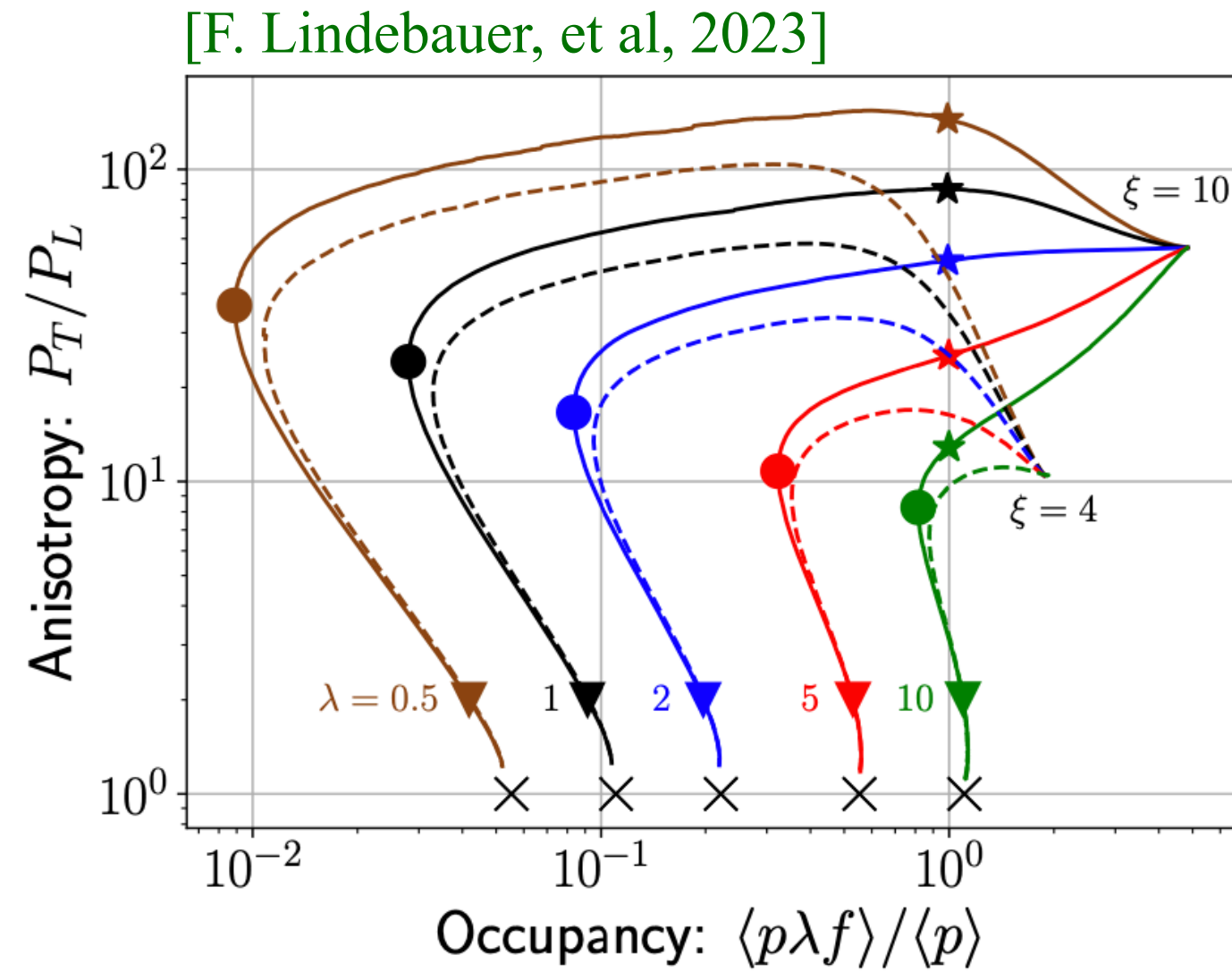
Another Landscape



Jet evolution in the Glasma



Why are the early stages “different”



- At early times there is a big pressure anisotropy

- This reflects in an anisotropic transport coefficient

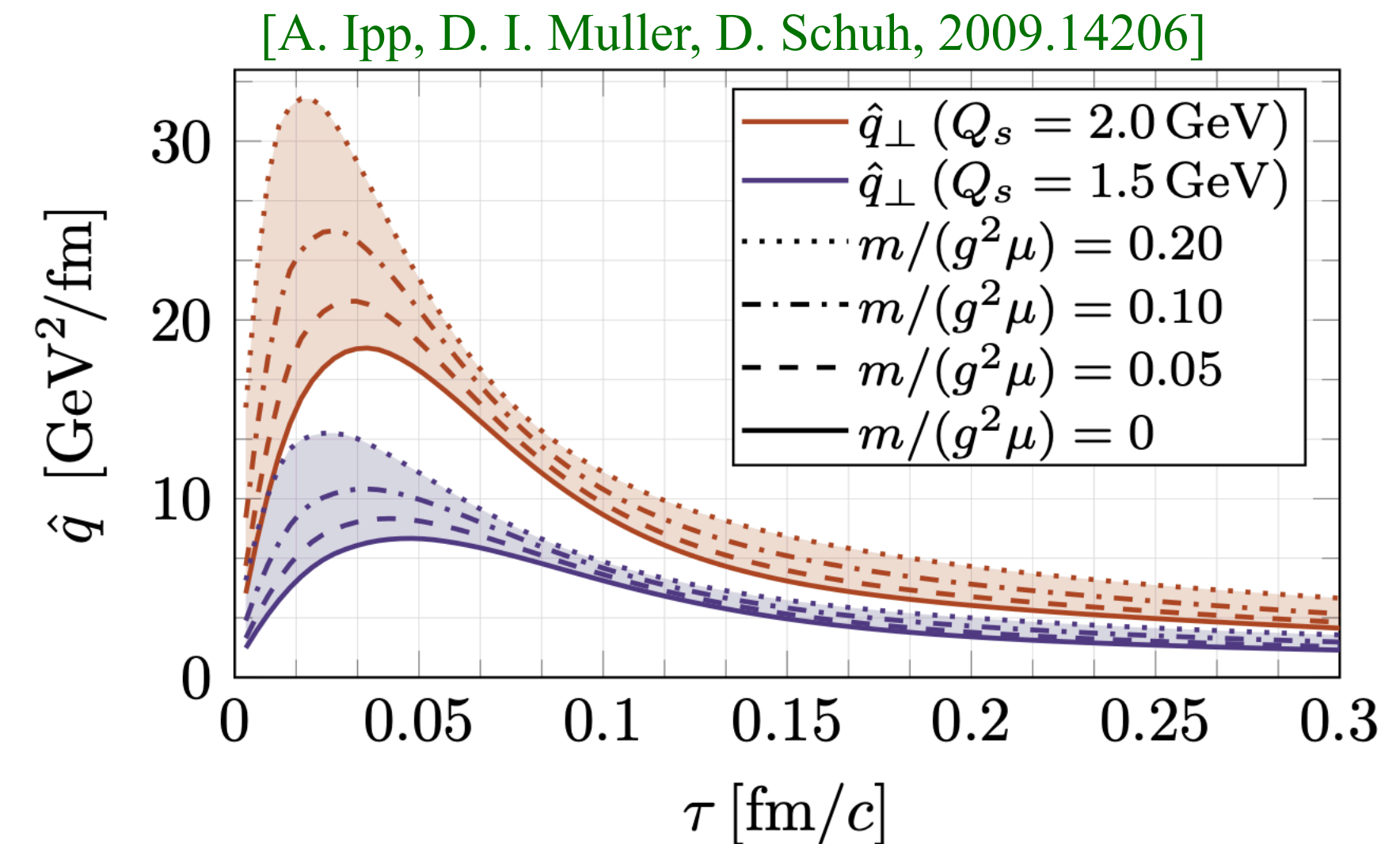
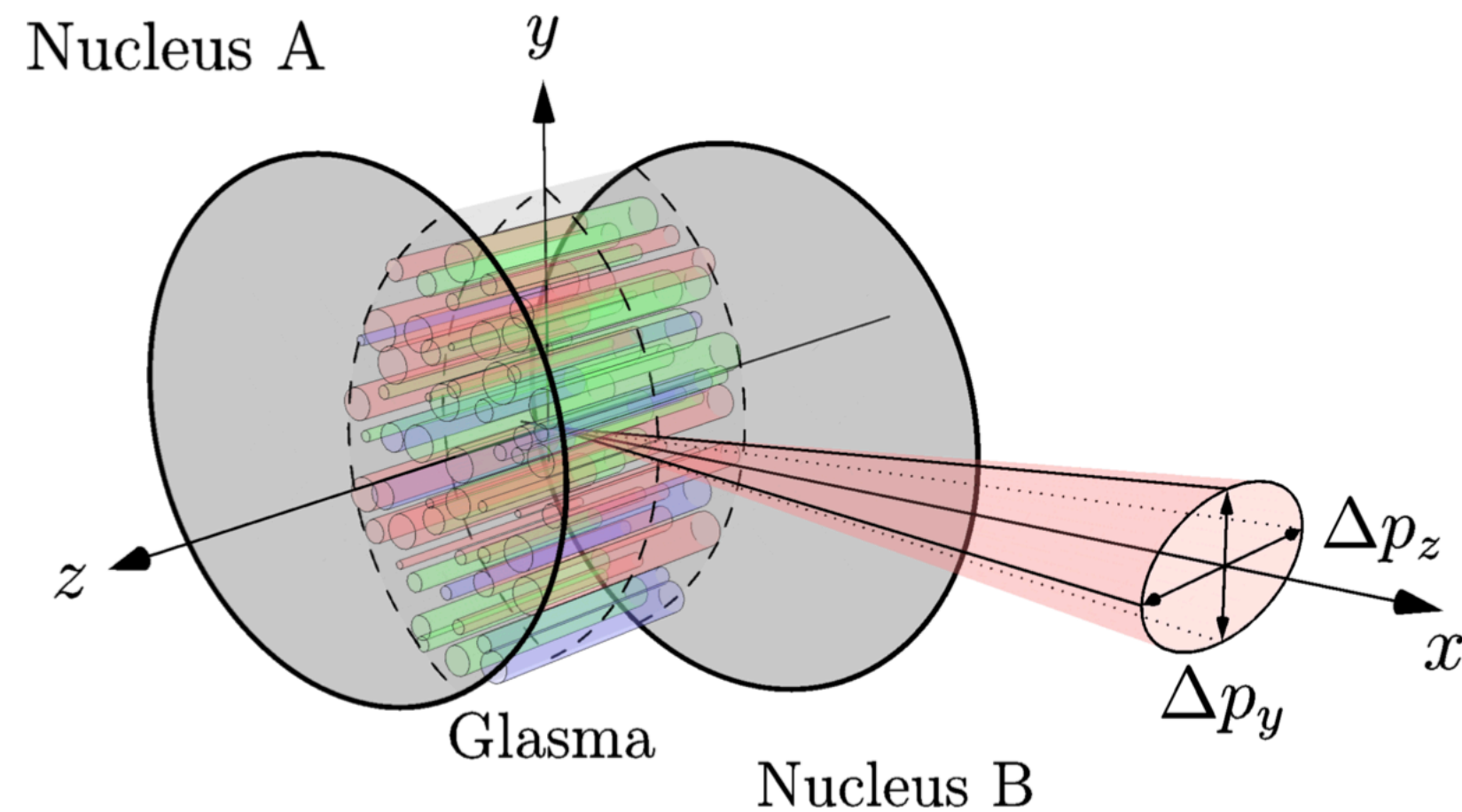
[Can be washed out by hydro expansion]

$$\hat{q}_x \gg \hat{q}_y$$

- Furthermore, the observed jet quenching coefficient seems to be much larger than the hydro one

[Observables integrate over jet path]

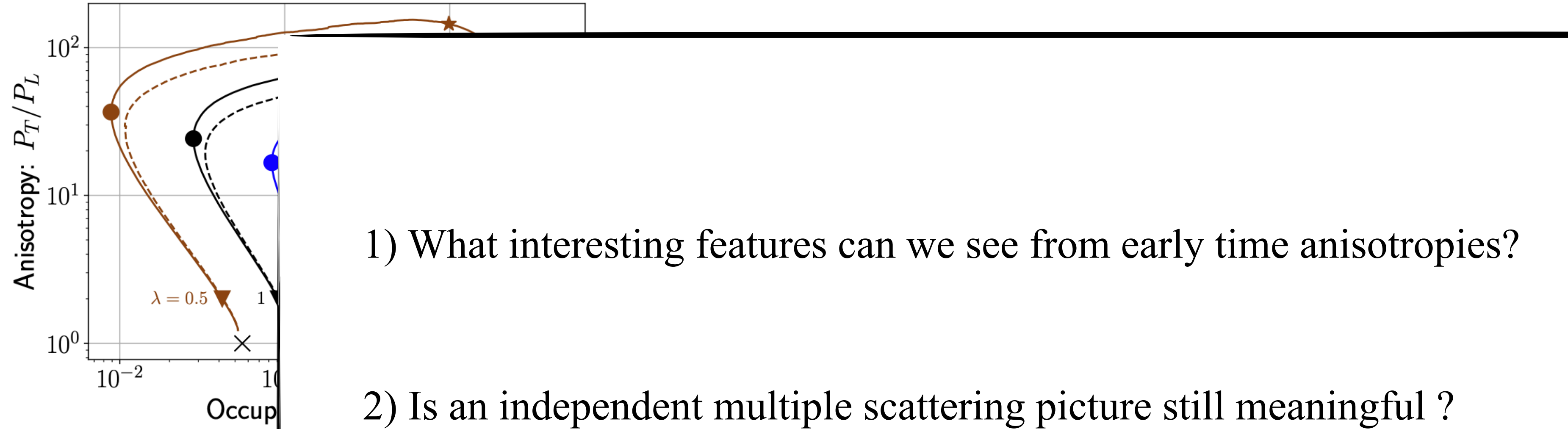
$$\hat{q}_{\text{early}} \gg \hat{q}_{\text{hydro}}, \tau_{\text{hydro}} \gg \tau_{\text{early}}$$



Why are the early stages “different”

- At early times there is a big pressure anisotropy

[F. Lindebauer, et al, 2023]



1) What interesting features can we see from early time anisotropies?

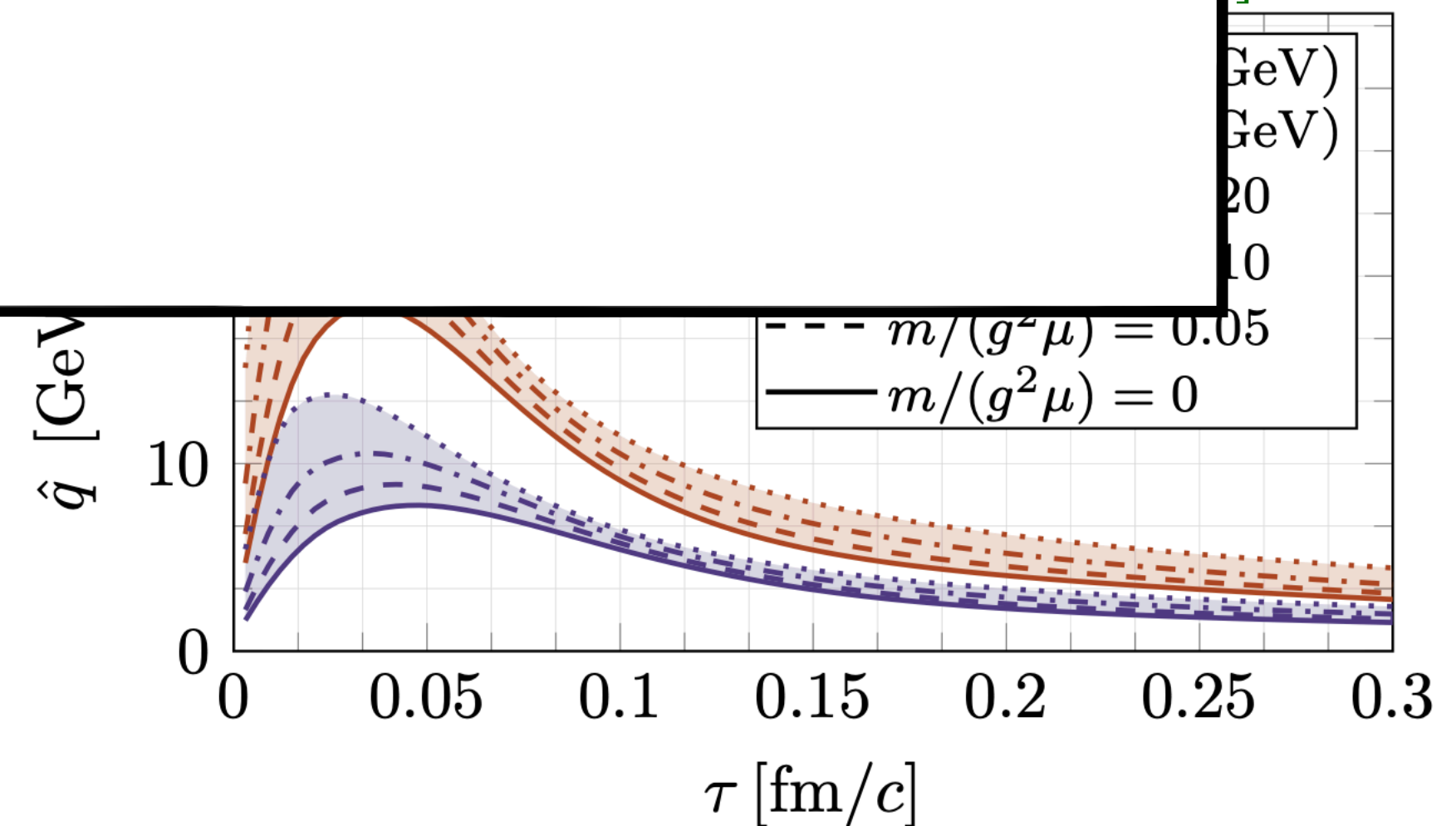
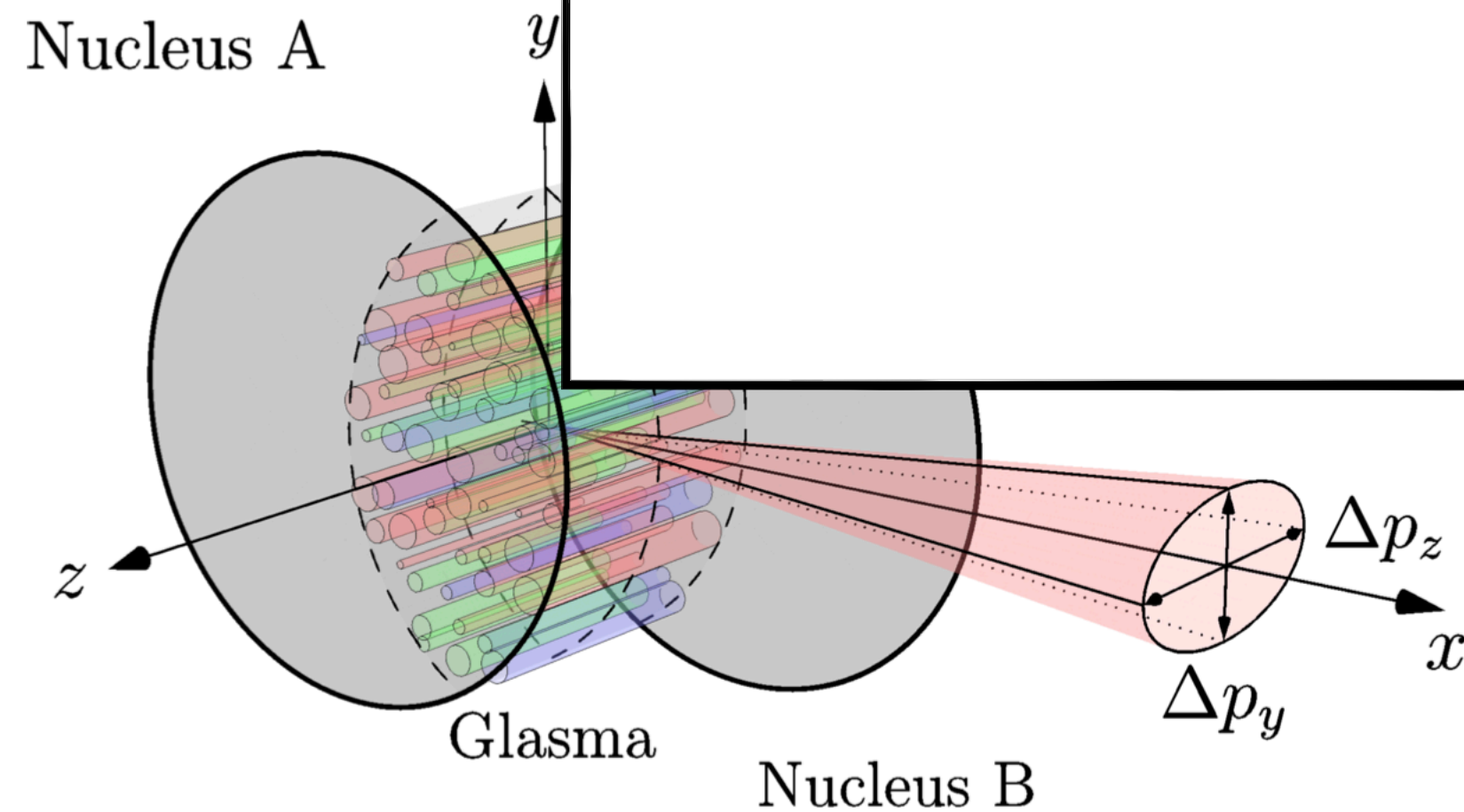
2) Is an independent multiple scattering picture still meaningful ?

coefficient
[d out by hydro expansion]

coefficient

observables integrate over jet path]

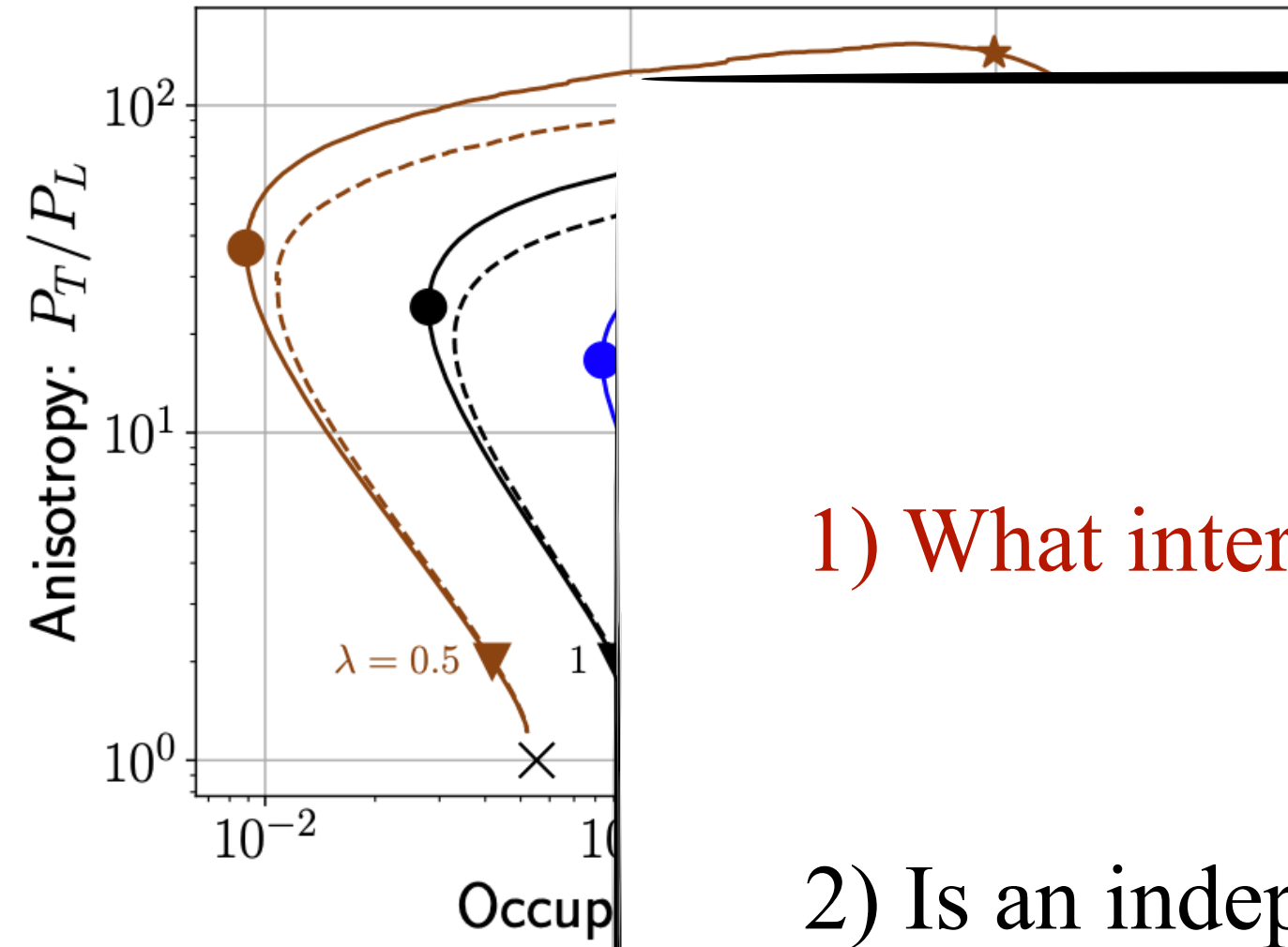
early



Why are the early stages “different”

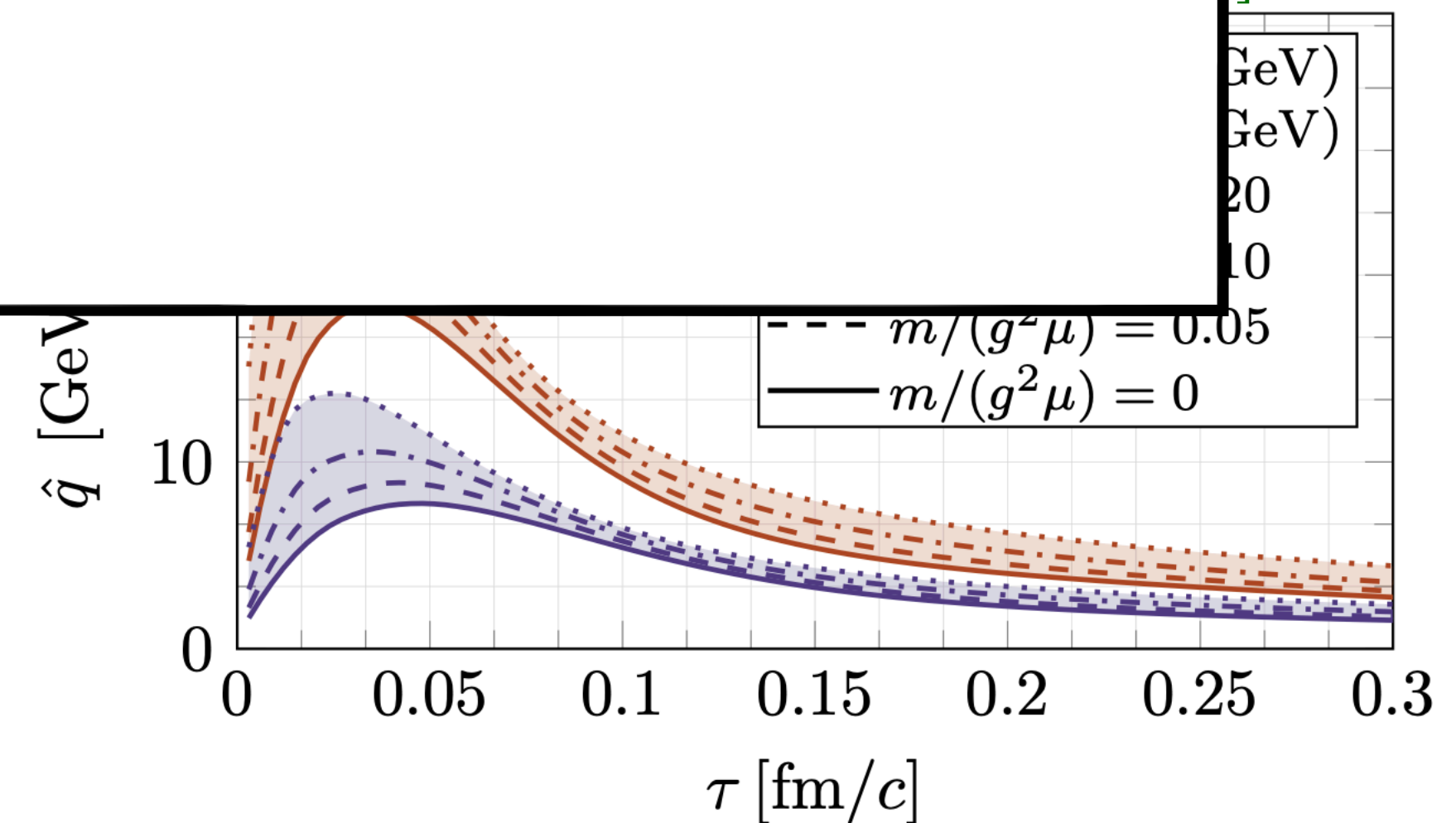
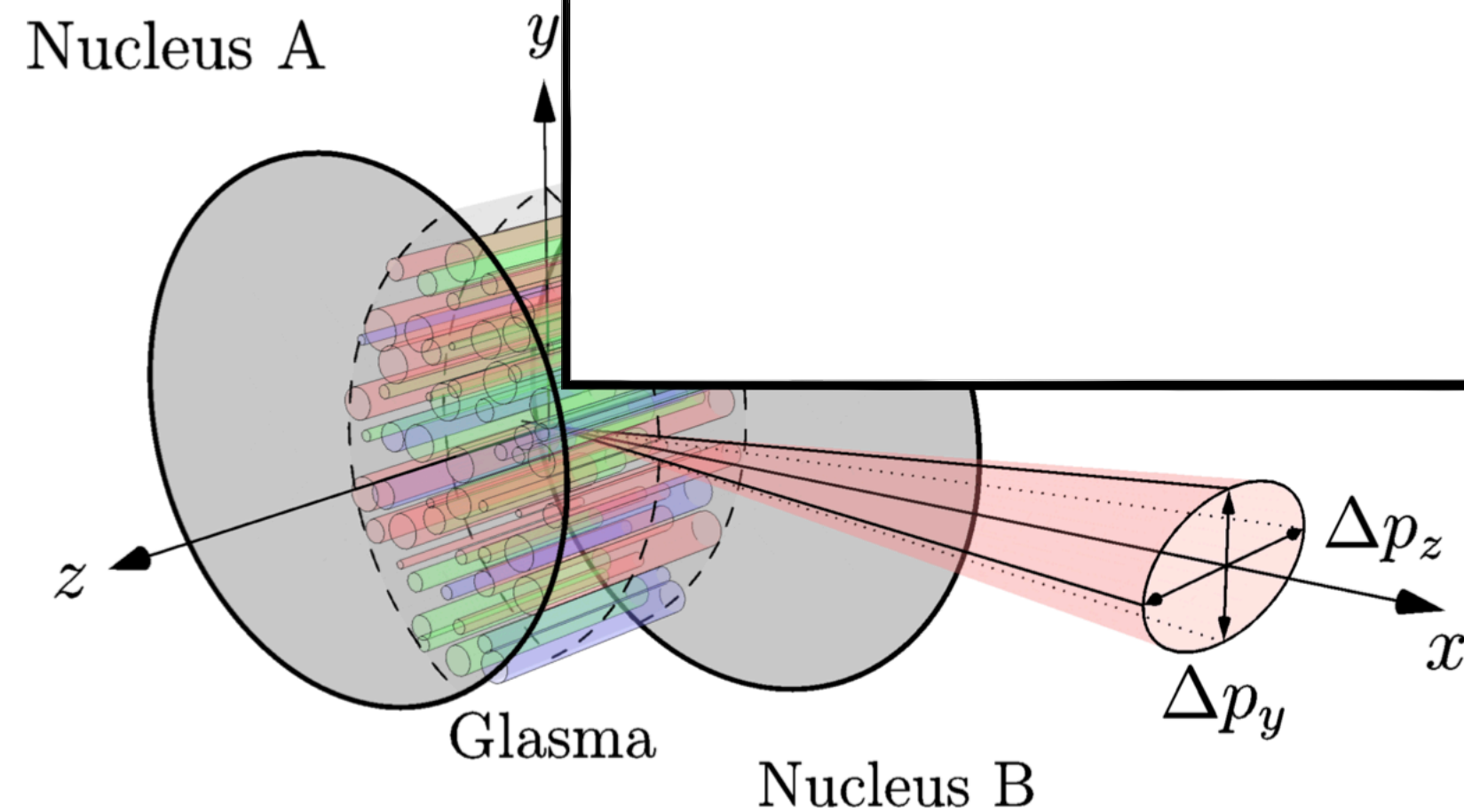
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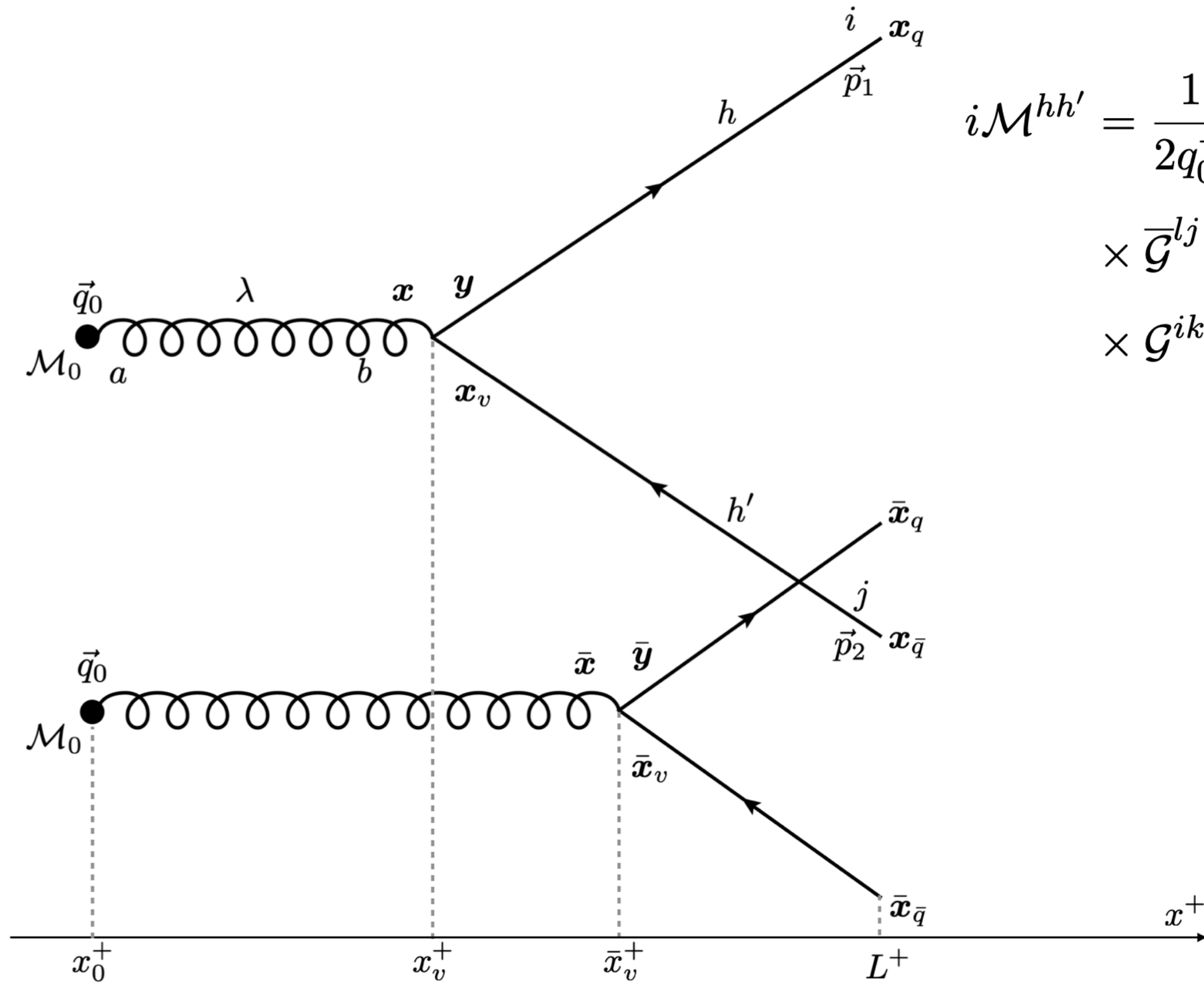
coefficient
[d out by hydro expansion]

coefficient

[observables integrate over jet path]

early

Similar to hydro gradients, evolution in the presence of anisotropic jet quenching coefficient leads to a non-trivial azimuthal structure. As an example consider g to $q\bar{q}$ in the presence of such a background



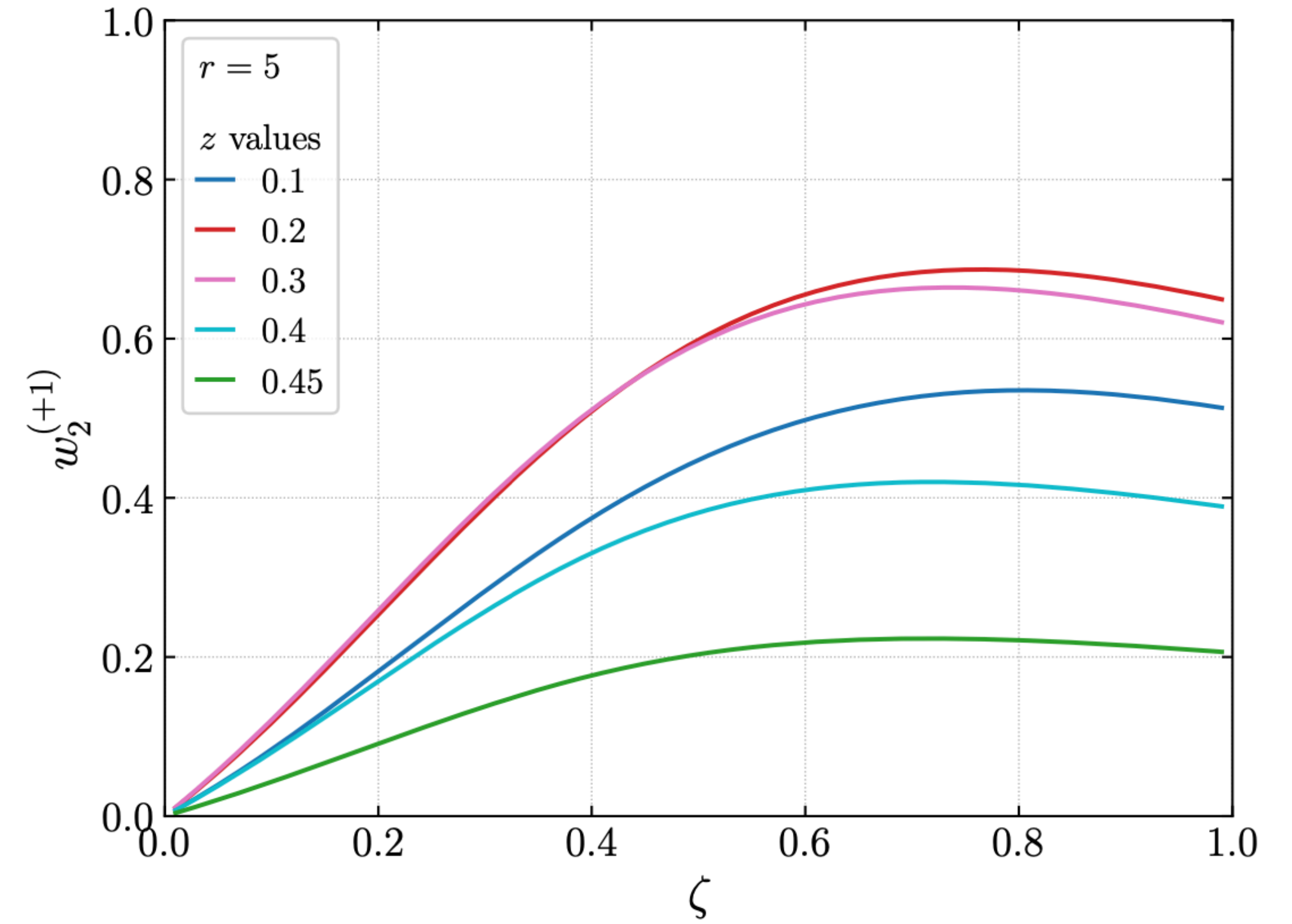
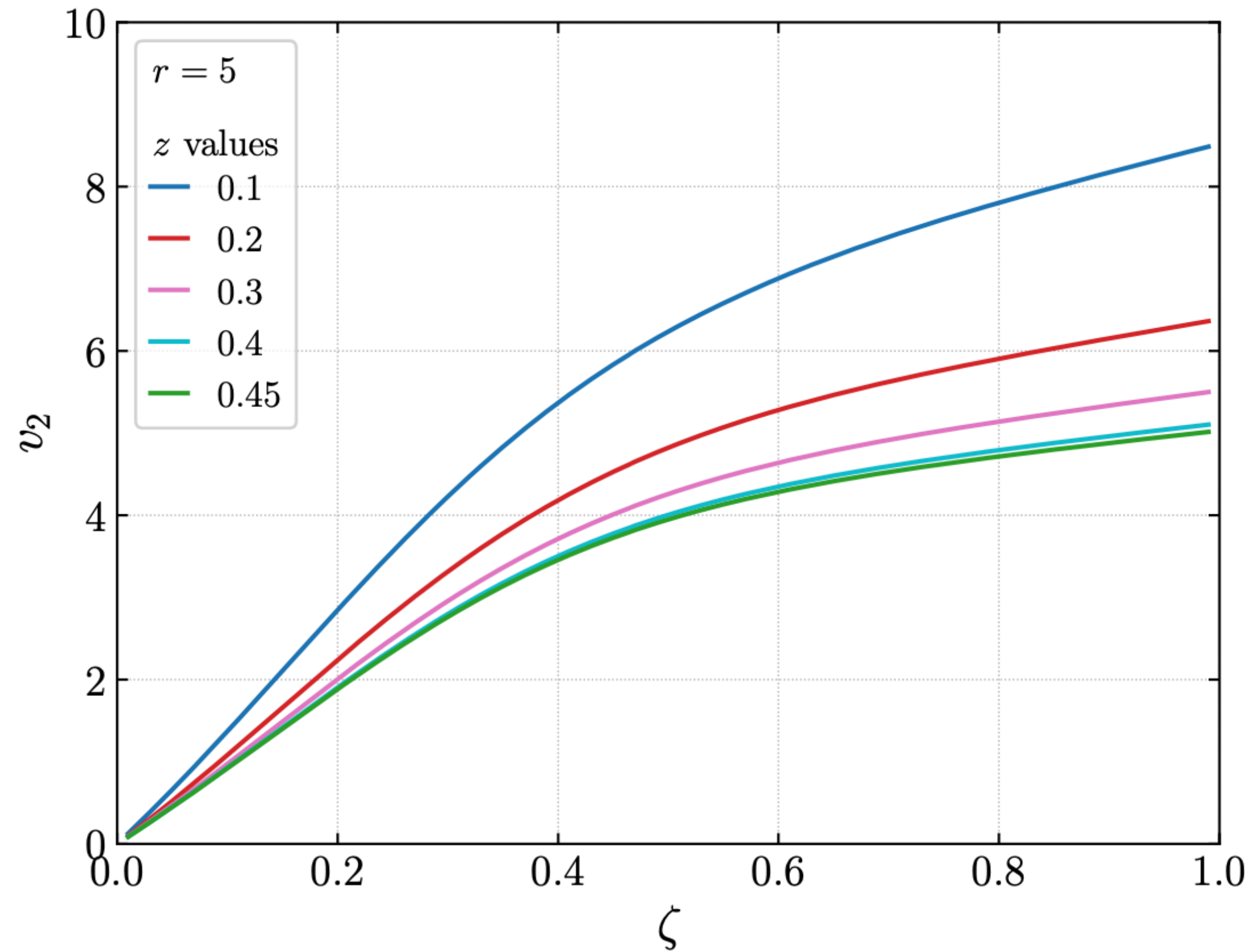
$$i\mathcal{M}^{hh'} = \frac{1}{2q_0^+} e^{i\frac{p_1^2}{2p_1^+}L^+} e^{i\frac{p_2^2}{2p_2^+}L^+} \sum_{\lambda=\pm 1} \int_{x_v^+} \int_{\mathbf{x}_v, \mathbf{x}_q, \mathbf{x}_{\bar{q}}, \mathbf{x}_g, \mathbf{x}, \mathbf{y}} \int_{x_g^+} \mathcal{M}_0^{\lambda, a}(q_0^+, x_g^+, \mathbf{x}_g) e^{-i\mathbf{p}_1 \cdot \mathbf{x}_q} e^{-i\mathbf{p}_2 \cdot \mathbf{x}_{\bar{q}}} \\ \times \bar{\mathcal{G}}^{lj}(L^+, \mathbf{x}_{\bar{q}}; x_v^+, \mathbf{x}_v | p_2^+) \left(igt_{kl}^b V^{\lambda hh'}(z, \mathbf{x}, \mathbf{y})\right) \mathcal{G}_A^{ba}(x_v^+, \mathbf{x}_v - \mathbf{x}; x_g^+, \mathbf{x}_g | q_0^+) \\ \times \mathcal{G}^{ik}(L^+, \mathbf{x}_q; x_v^+, \mathbf{x}_v - \mathbf{y} | p_1^+),$$

$$V^{\lambda hh'}(z, \mathbf{x}, \mathbf{y}) = 2\gamma^{\lambda h}(z)\delta_{h, -h'}\epsilon_\lambda \cdot i(\delta(\mathbf{x})\delta'(\mathbf{y}) + z\delta'(\mathbf{x})\delta(\mathbf{y}))$$

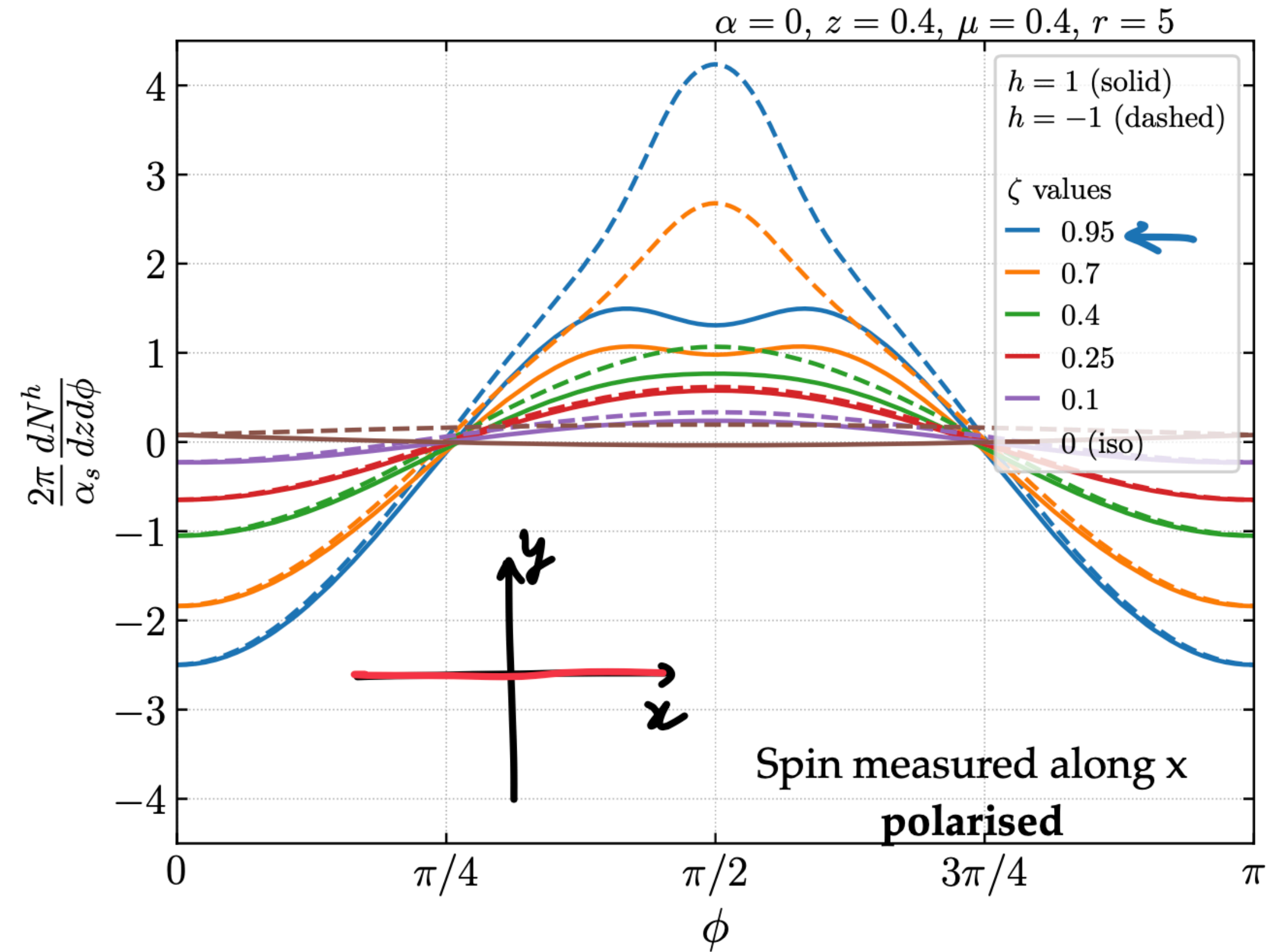
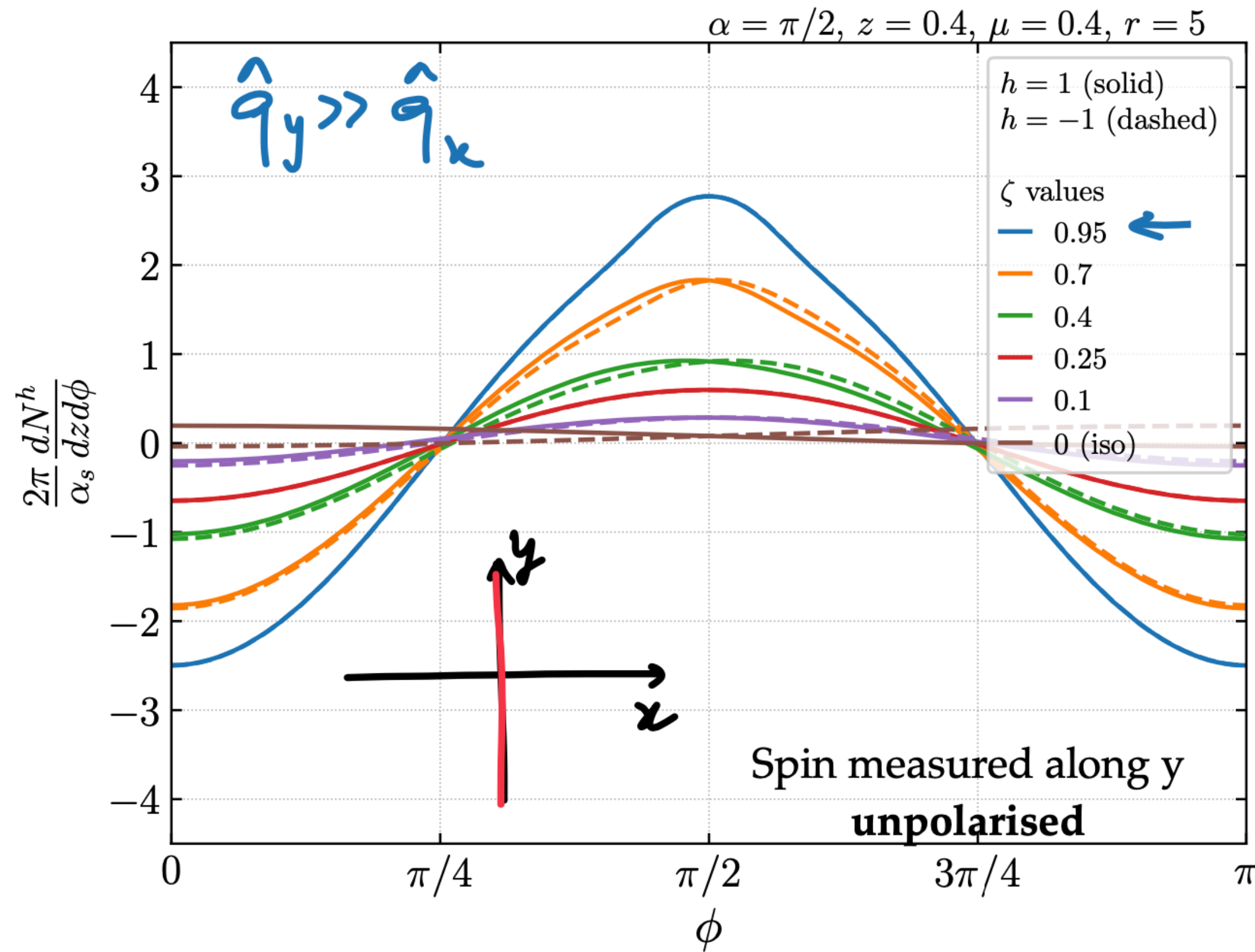
$$\hat{q} = (\hat{q}_x, \hat{q}_y \gg \hat{q}_x)$$

$$\frac{2\pi}{dN^h/dz} \frac{dN^h}{dzd\phi} = 1 + \sum_{n=1}^{\infty} v_n \cos(n\phi) + \sum_{n=1}^{\infty} w_n^{(h)} \sin(n\phi)$$

$$\zeta = \frac{\sqrt{\hat{q}_y} - \sqrt{\hat{q}_x}}{\sqrt{\hat{q}_y} + \sqrt{\hat{q}_x}}, \quad r = L^+ \frac{(\sqrt{\hat{q}_y} + \sqrt{\hat{q}_x})}{2\sqrt{q_0^+}}$$



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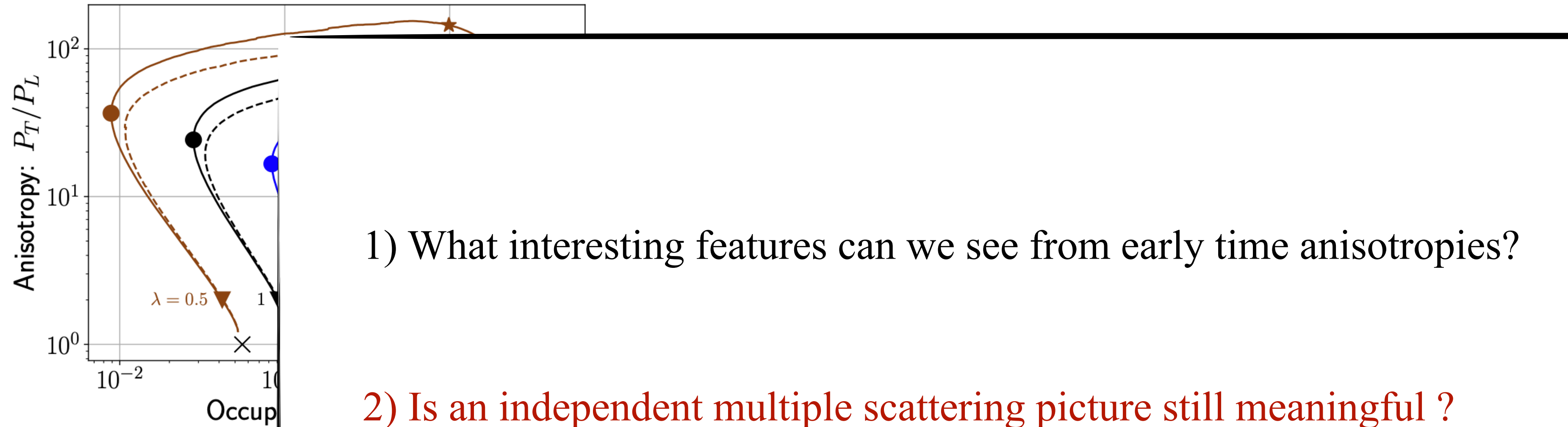


Similar observations in [S. Hauksson, E. Iancu, 2303.03914]

Why are the early stages “different”

- At early times there is a big pressure anisotropy

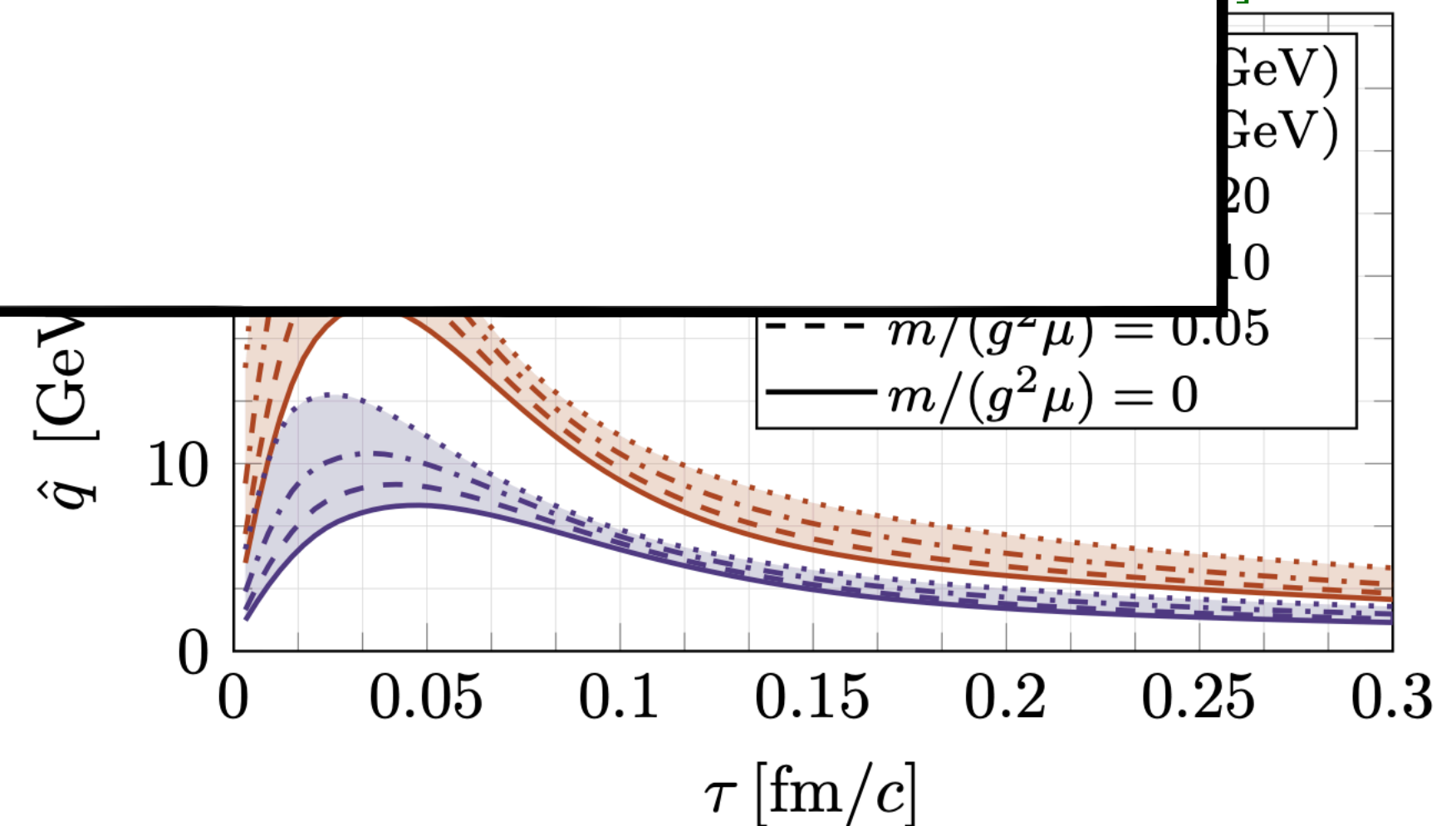
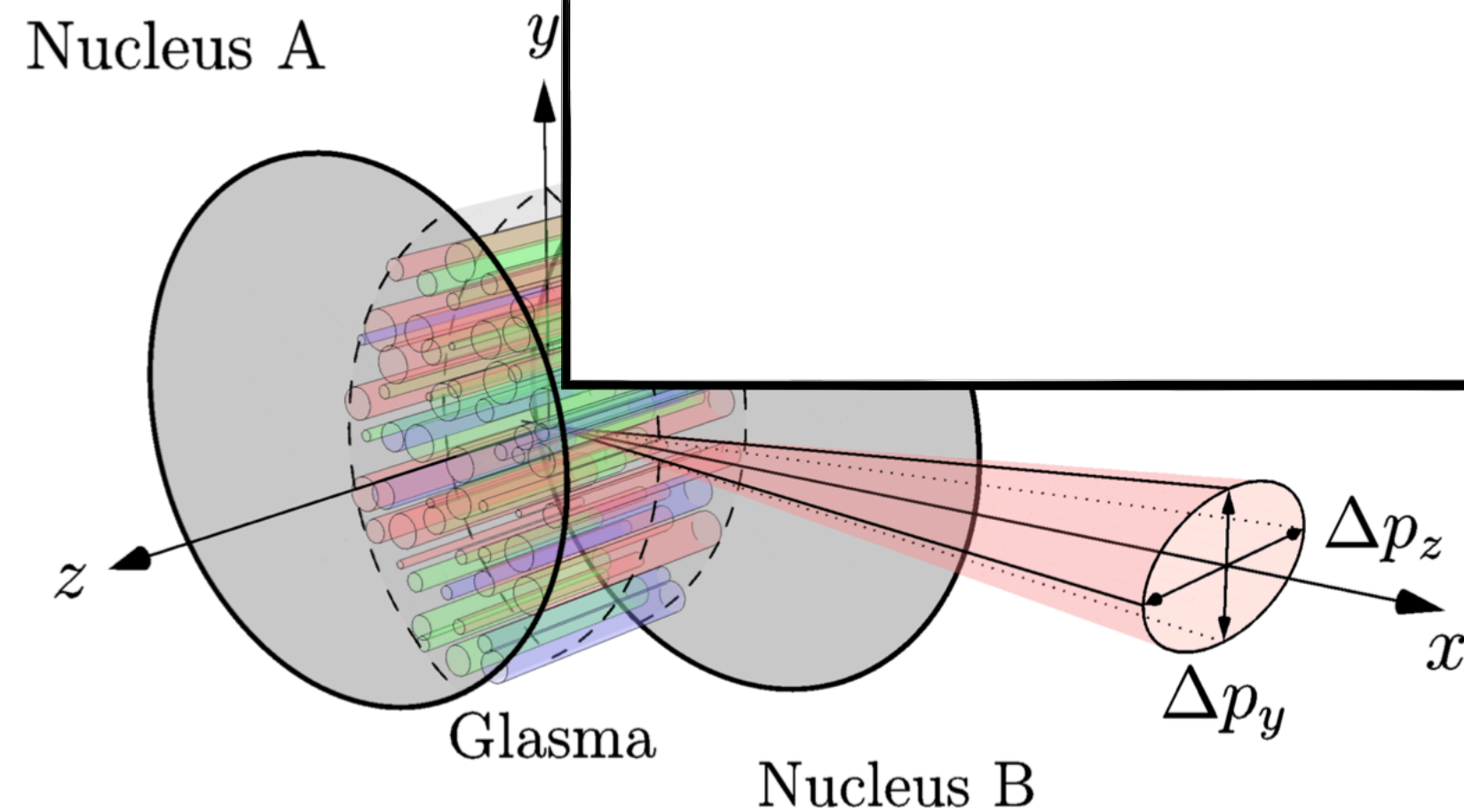
[F. Lindebauer, et al, 2023]

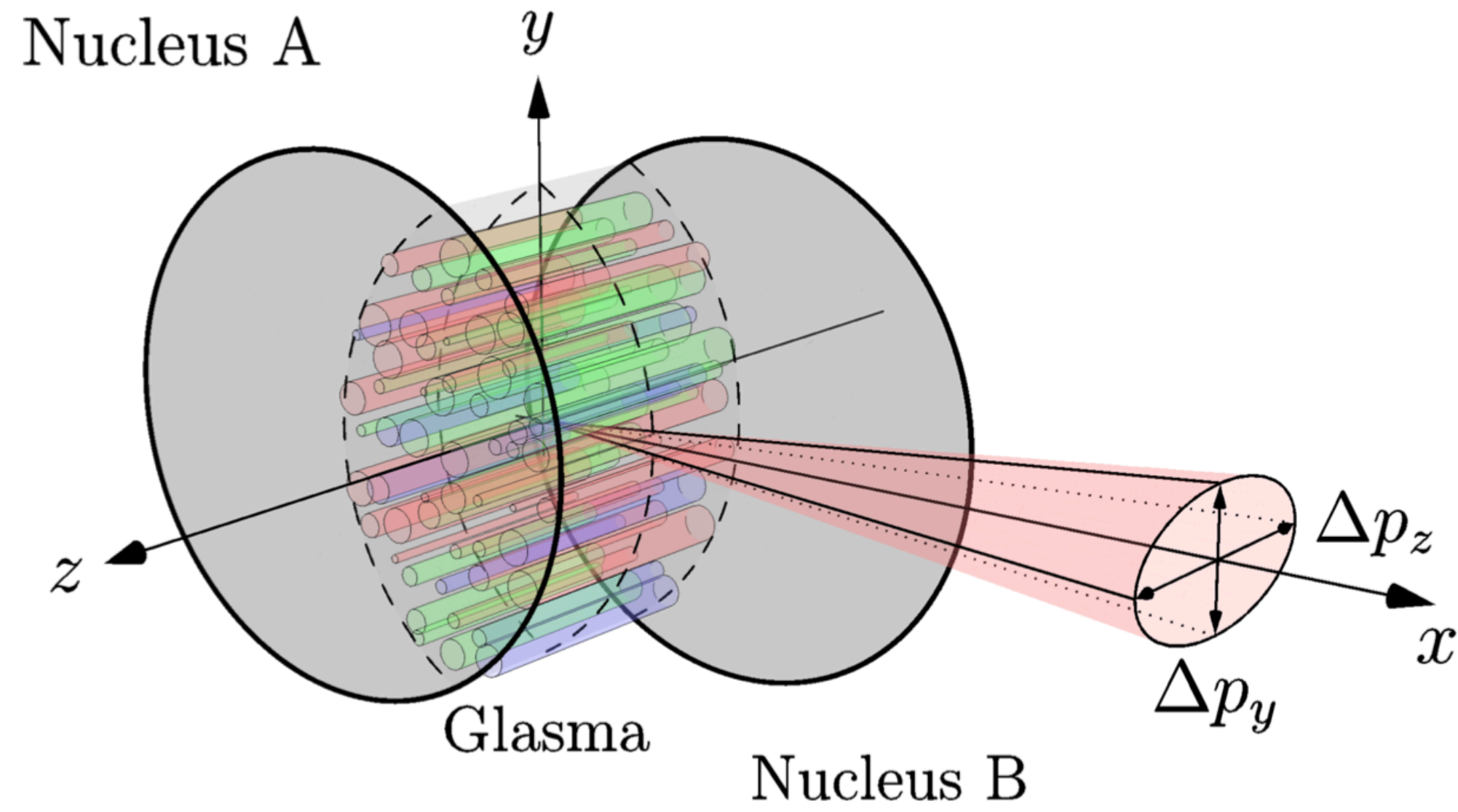


1) What interesting features can we see from early time anisotropies?

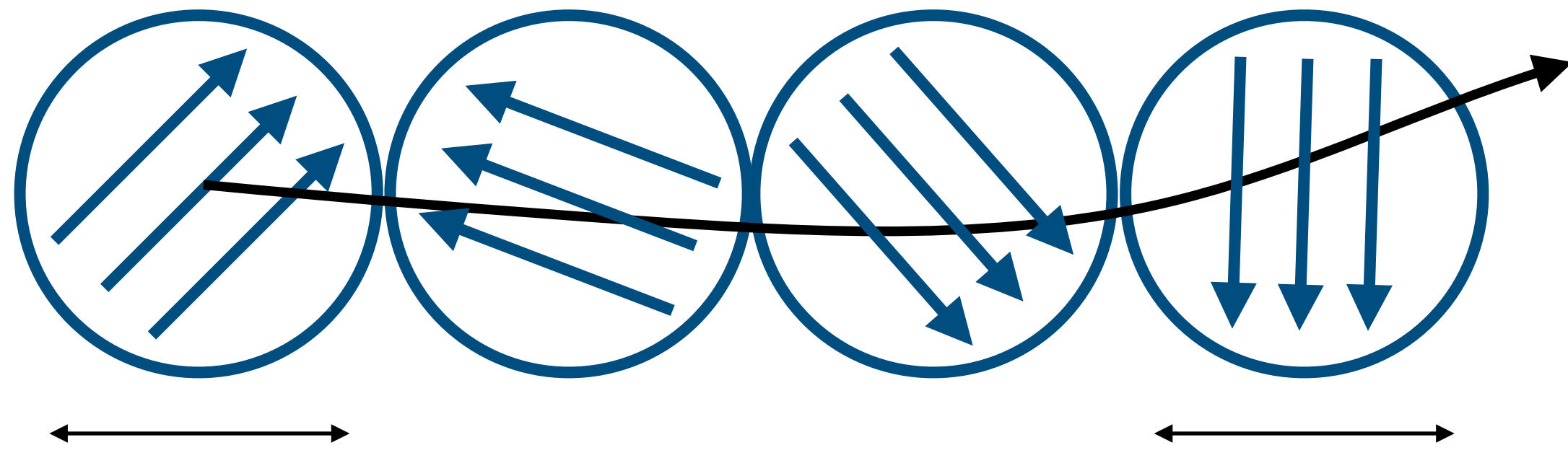
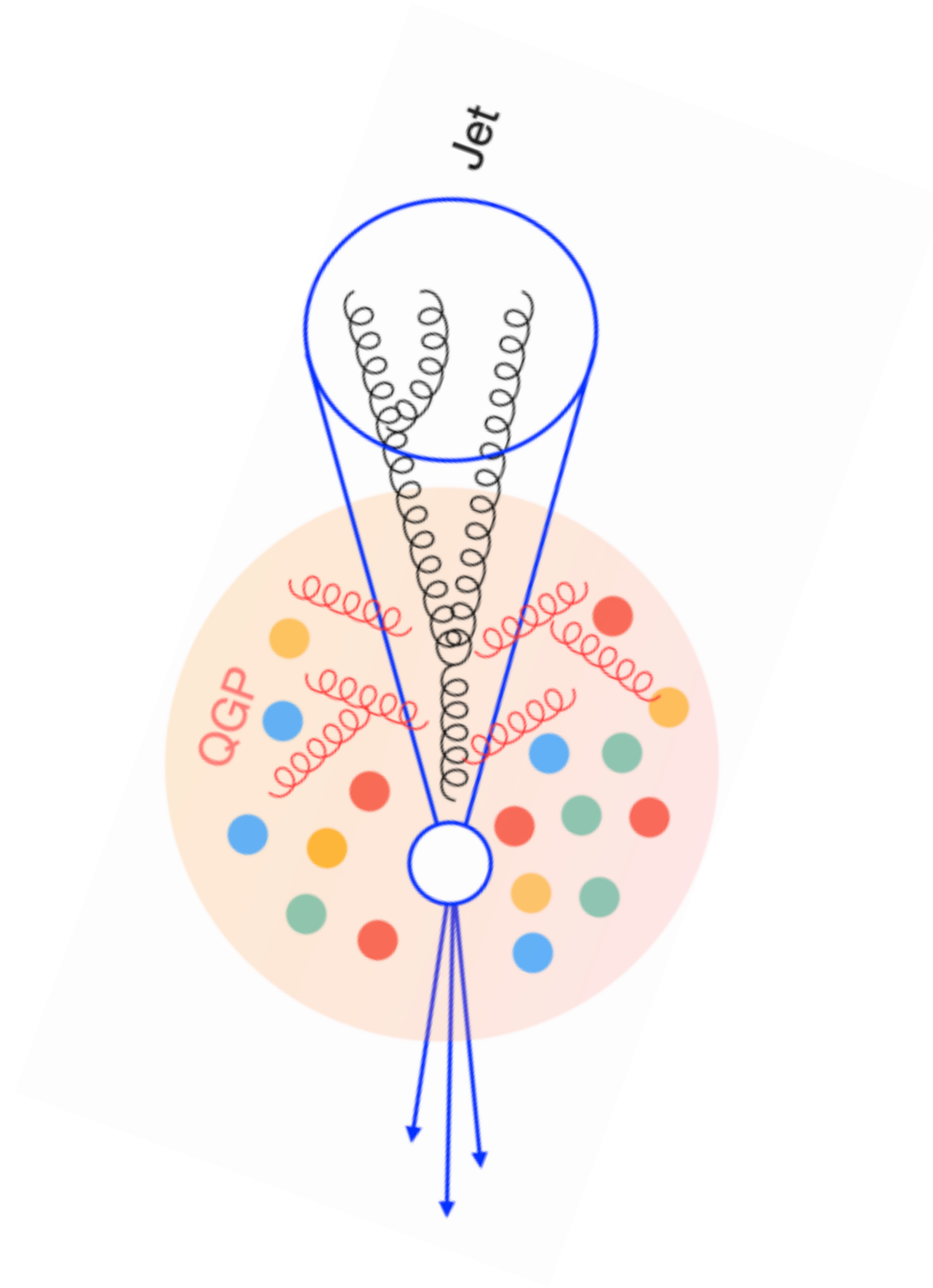
2) Is an independent multiple scattering picture still meaningful ?

coefficient
[d out by hydro expansion]
coefficient
observables integrate over jet path]
early



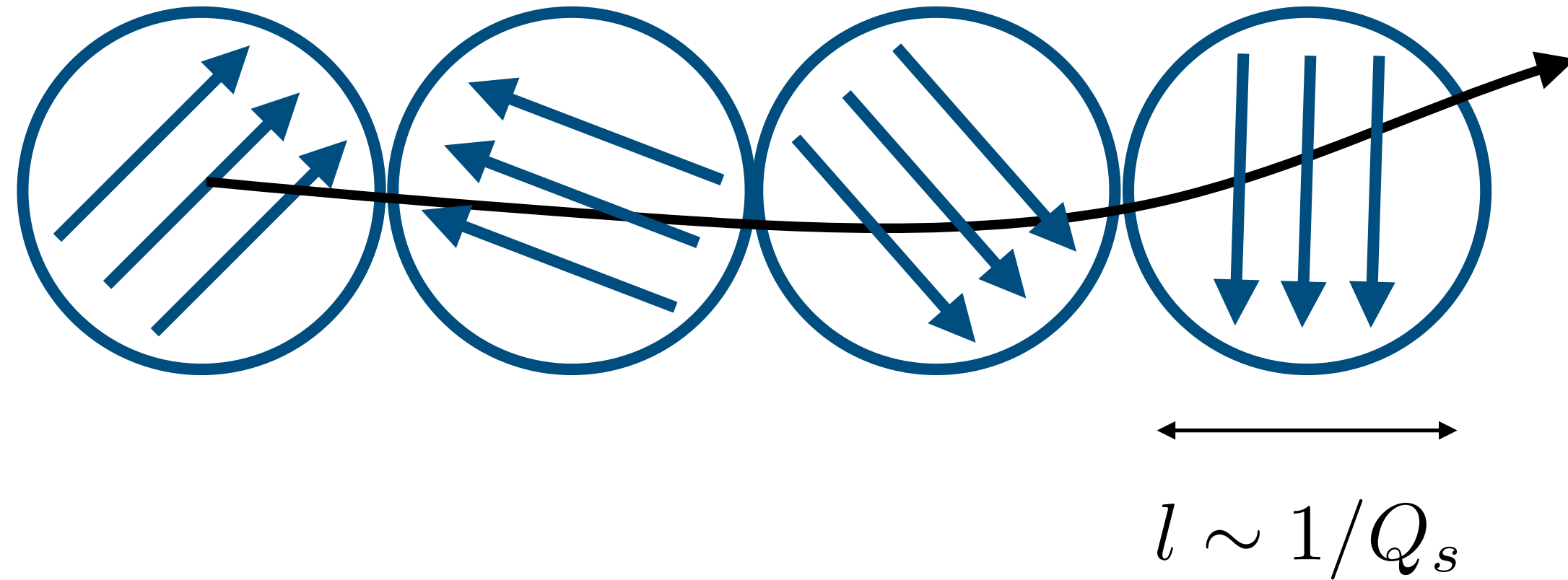


??
=



Synchrotron like scenario

$$l \sim 1/Q_s$$

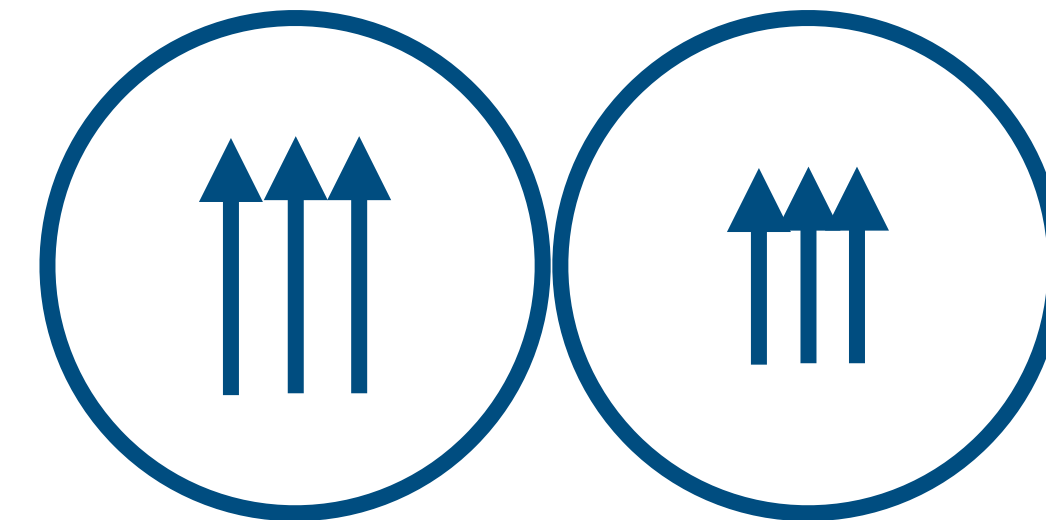


$$A_{\text{coh}}^{a\mu}(\mathbf{x}, z) = \delta_{\mu}^0 \mathbf{x} \cdot \mathbf{E}^a(z) = \begin{cases} \delta^{\mu 0} \mathbf{x} \cdot \mathbf{E}_1^a, & 0 \leq z < \ell \\ \delta^{\mu 0} \mathbf{x} \cdot \mathbf{E}_2^a, & \ell \leq z < 2\ell \\ \delta^{\mu 0} \mathbf{x} \cdot \mathbf{E}_3^a, & 2\ell \leq z < 3\ell \\ \vdots & \end{cases}$$

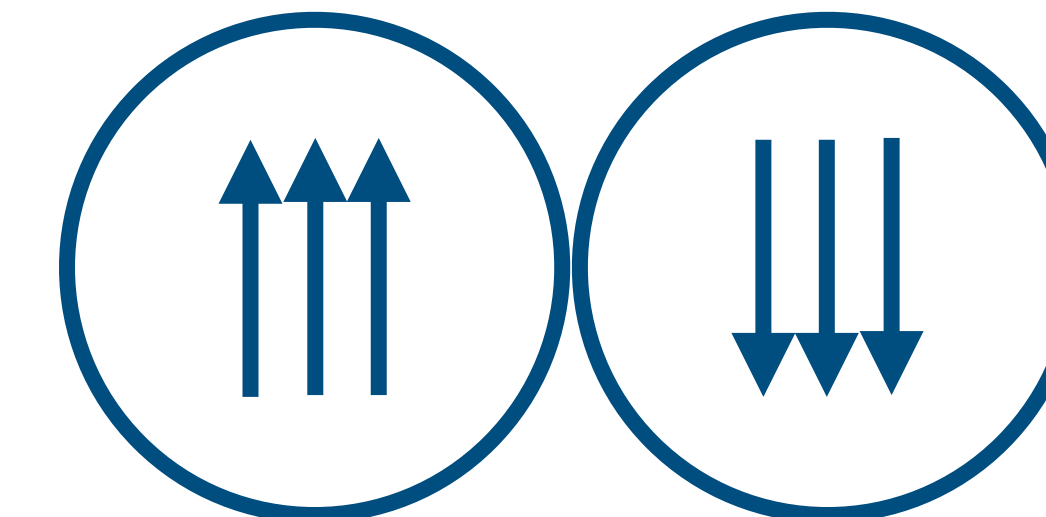
$$\langle A(x)A(y) \rangle \neq \delta(x^+ - y^+)$$

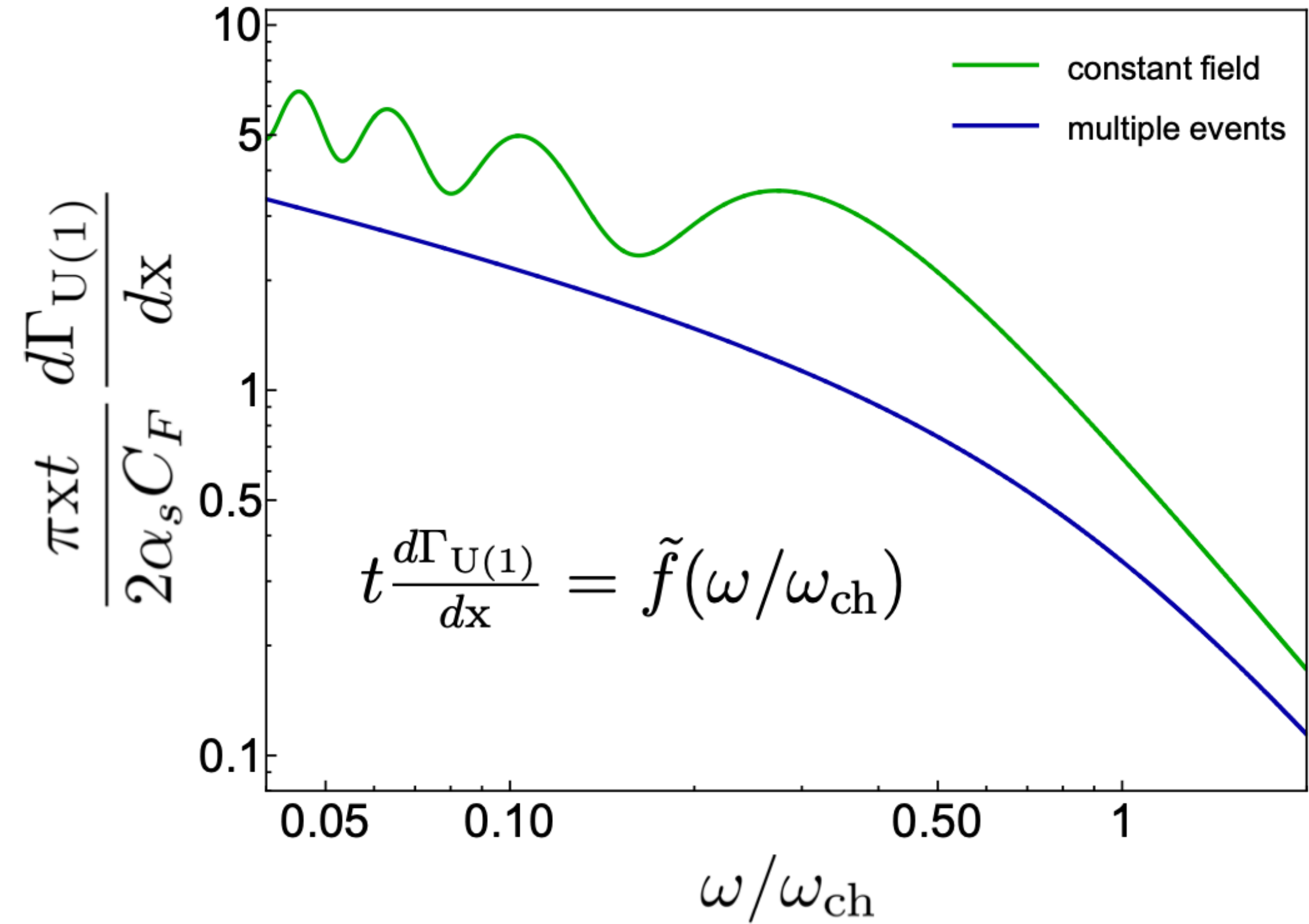
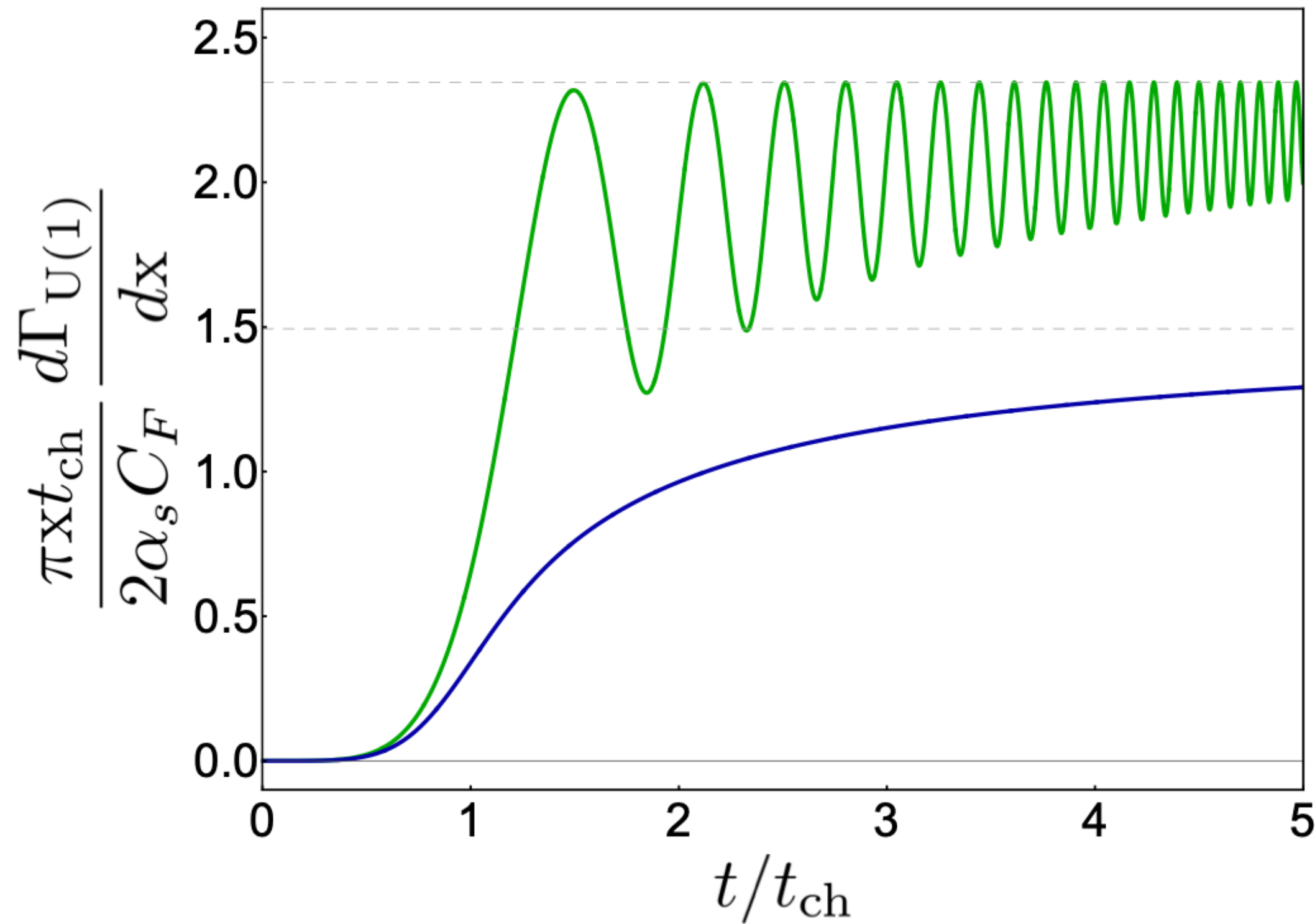
We consider two medium models:

- $\langle f(E_{1x}^a, E_{2x}^a, E_{3x}^a, \dots) \rangle = \int_{E_1} e^{-E_{1x}^2/E_0^2} \int_{E_2} e^{-E_{2x}^2/E_0^2} \dots f(E_{1x}^a, E_{2x}^a, E_{3x}^a, \dots)$



- $\langle f(E_{1x}, E_{2x}, E_{3x}, \dots) \rangle = \left[\prod_n \int_{E_n} \frac{dE_{nx}}{2} (\delta(E_{nx} - E_0) + \delta(E_{nx} + E_0)) \right] f(E_{1x}, E_{2x}, E_{3x}, \dots)$





At early and late times, we find :

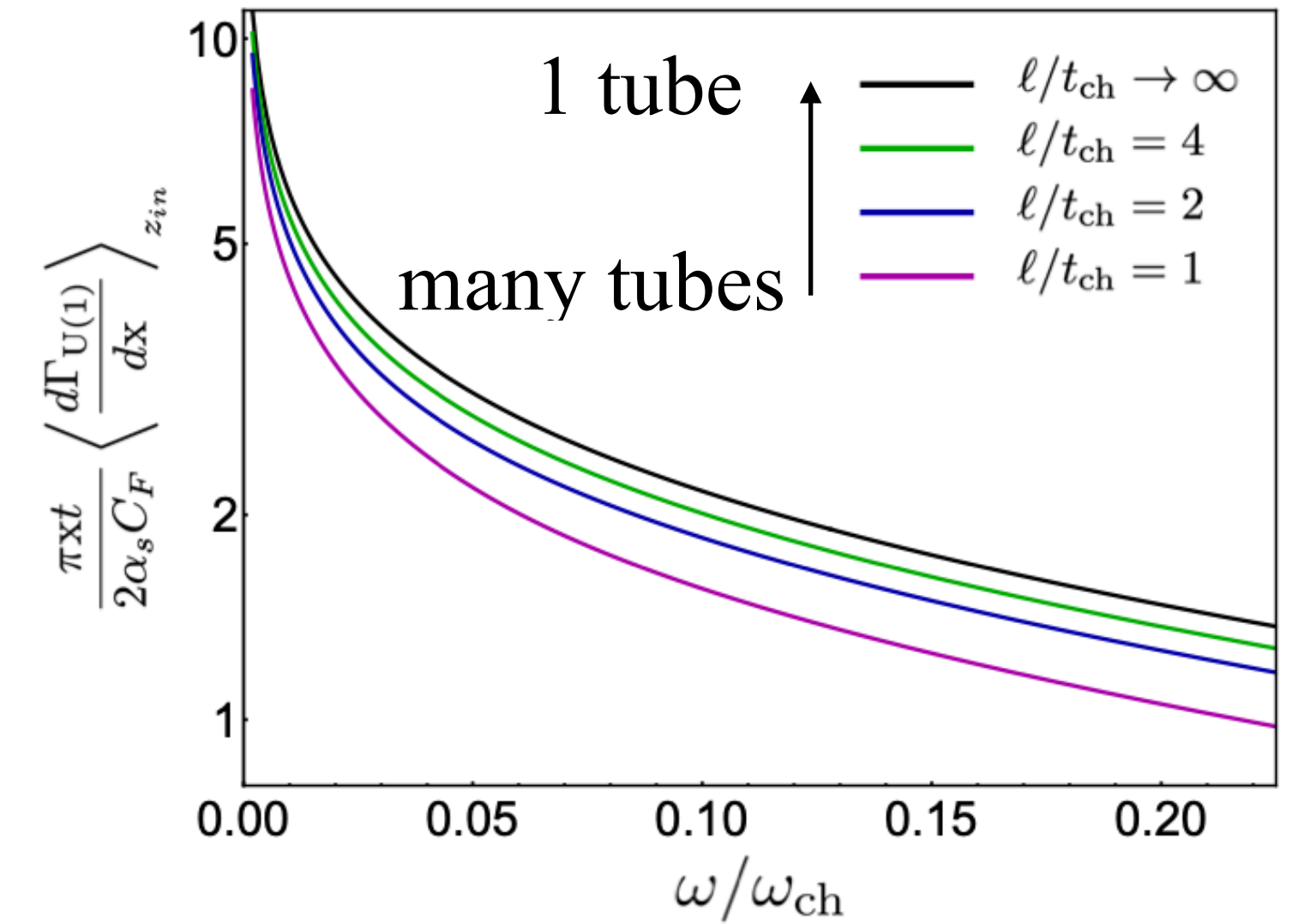
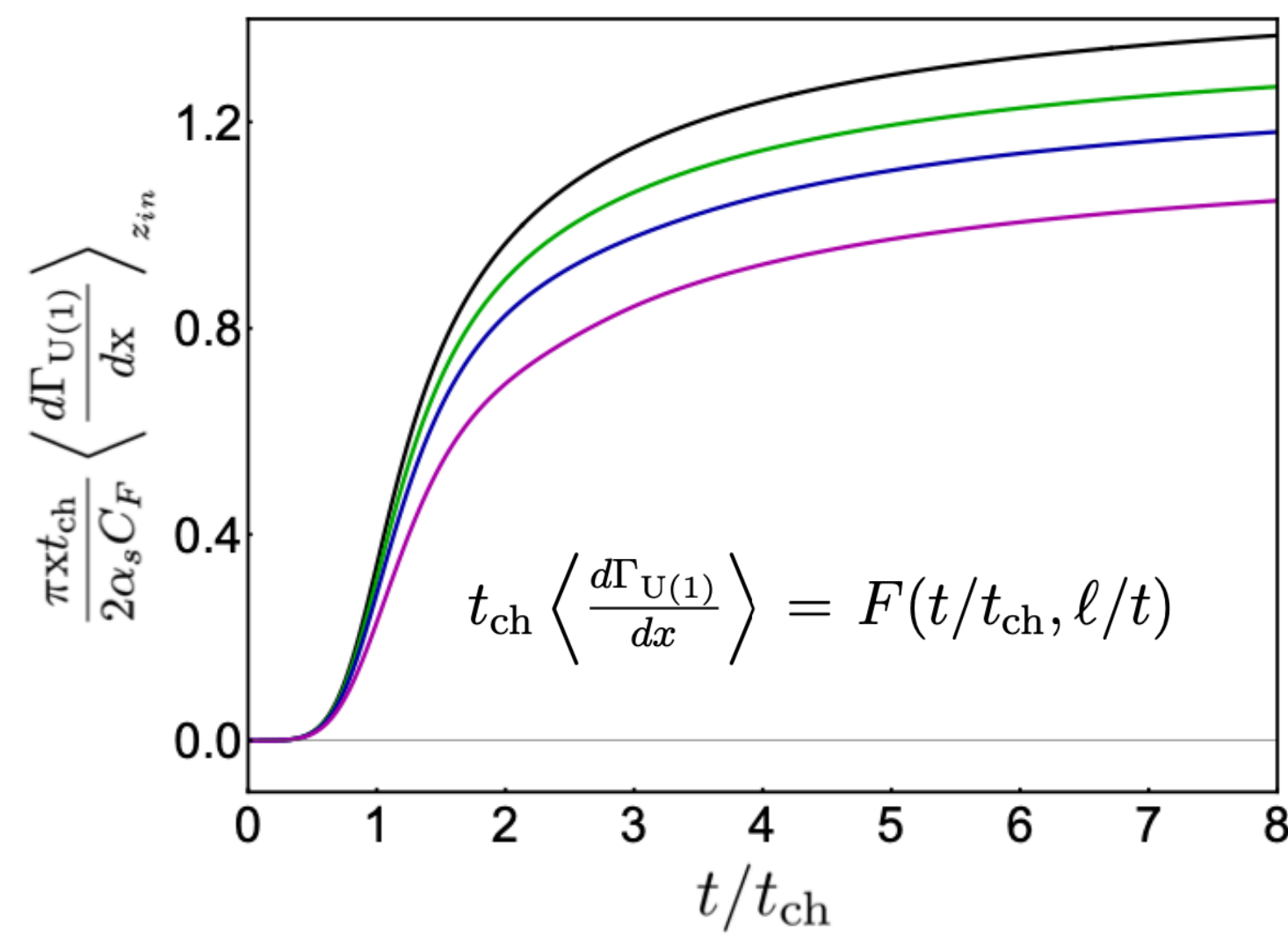
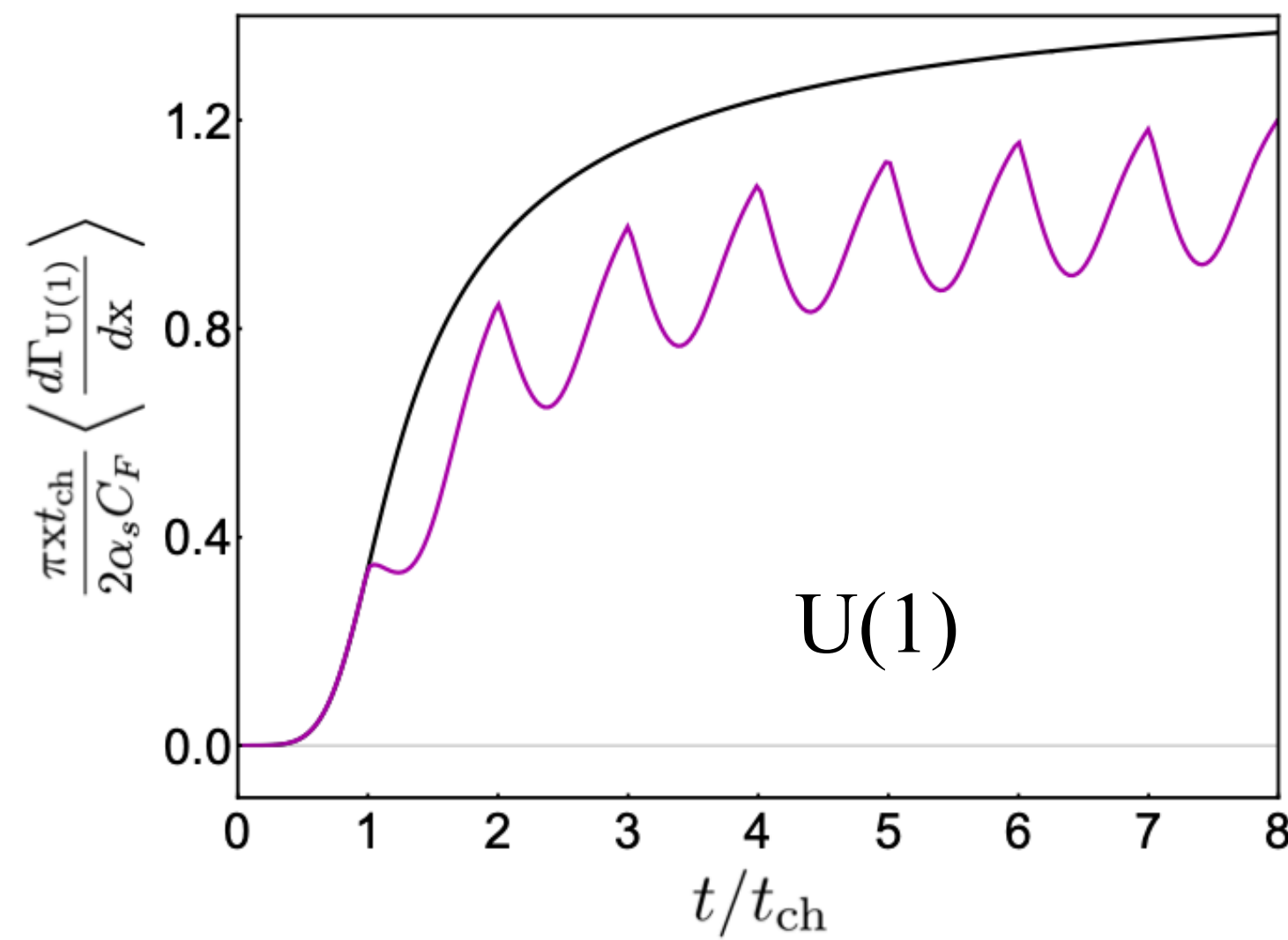
$$\frac{d\Gamma_{U(1)}}{dx} \simeq \frac{7\alpha_s C_F}{24\pi} \frac{E^4}{\omega^2} \frac{t^5}{5!}$$

$$\frac{d\Gamma_{U(1)}}{dx} \Big|_{t \rightarrow \infty} = 3^{1/6} \Gamma\left(\frac{2}{3}\right) \frac{\alpha_s C_F}{\pi} E^{2/3} \omega^{-1/3}$$

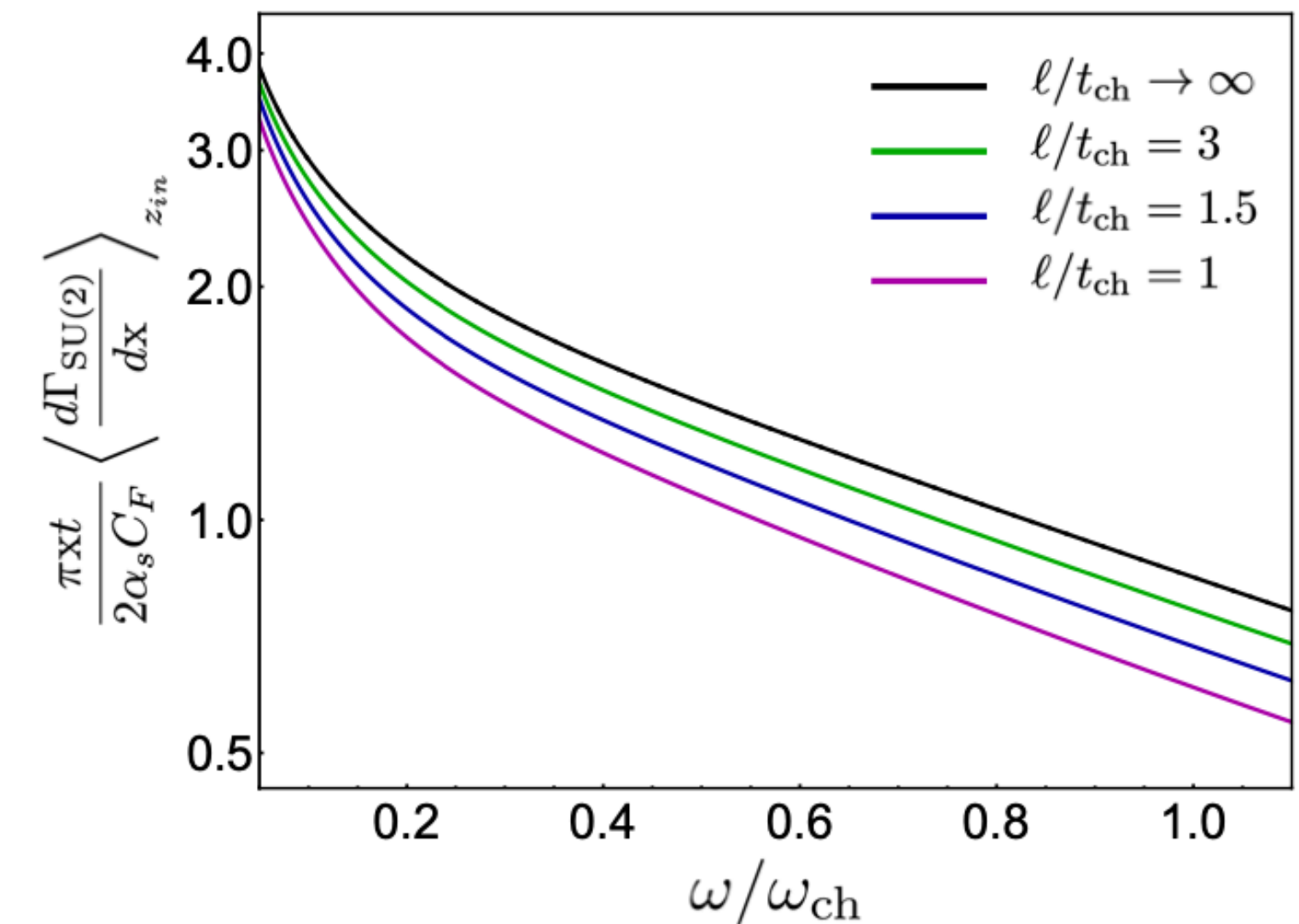
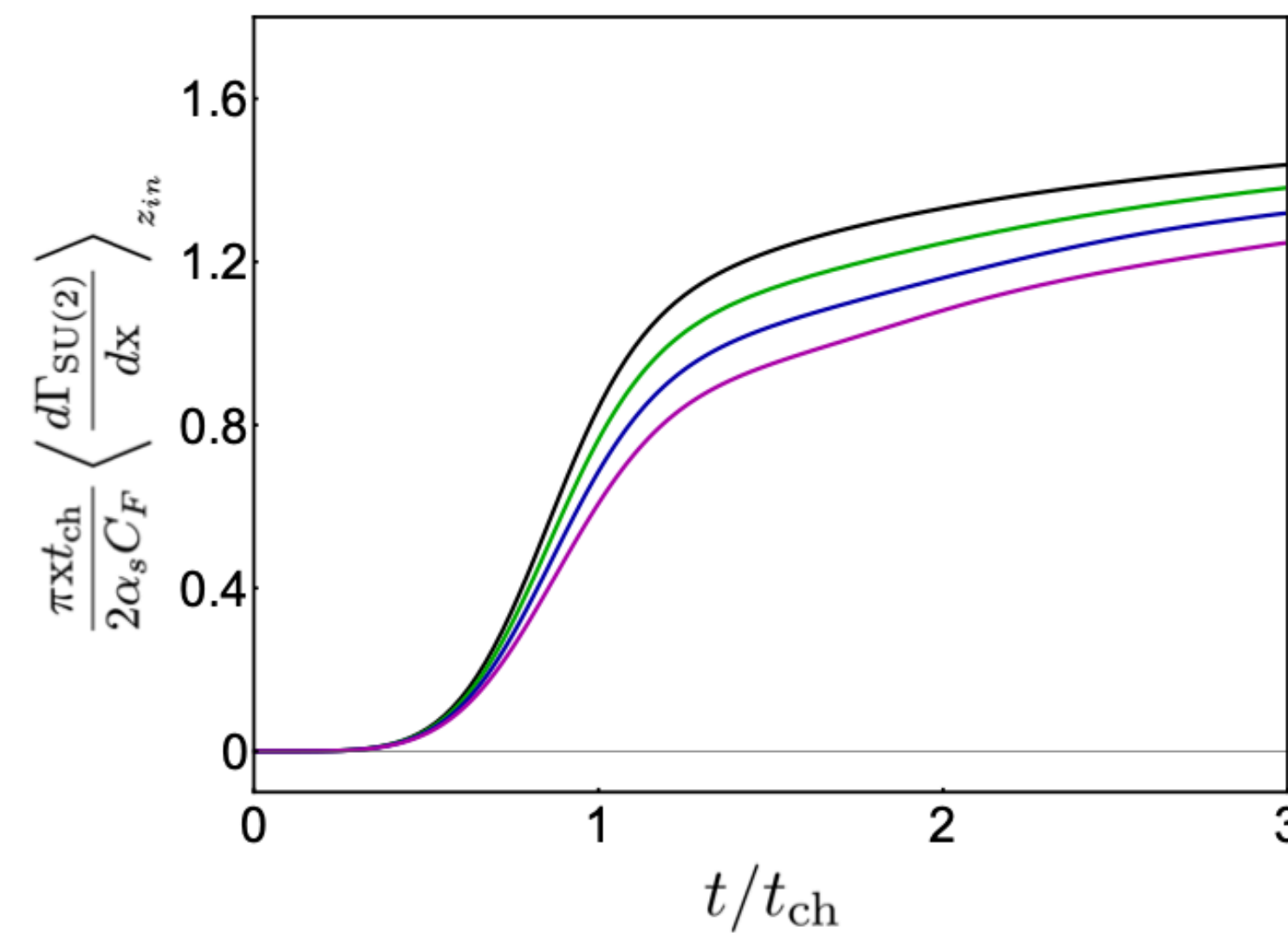
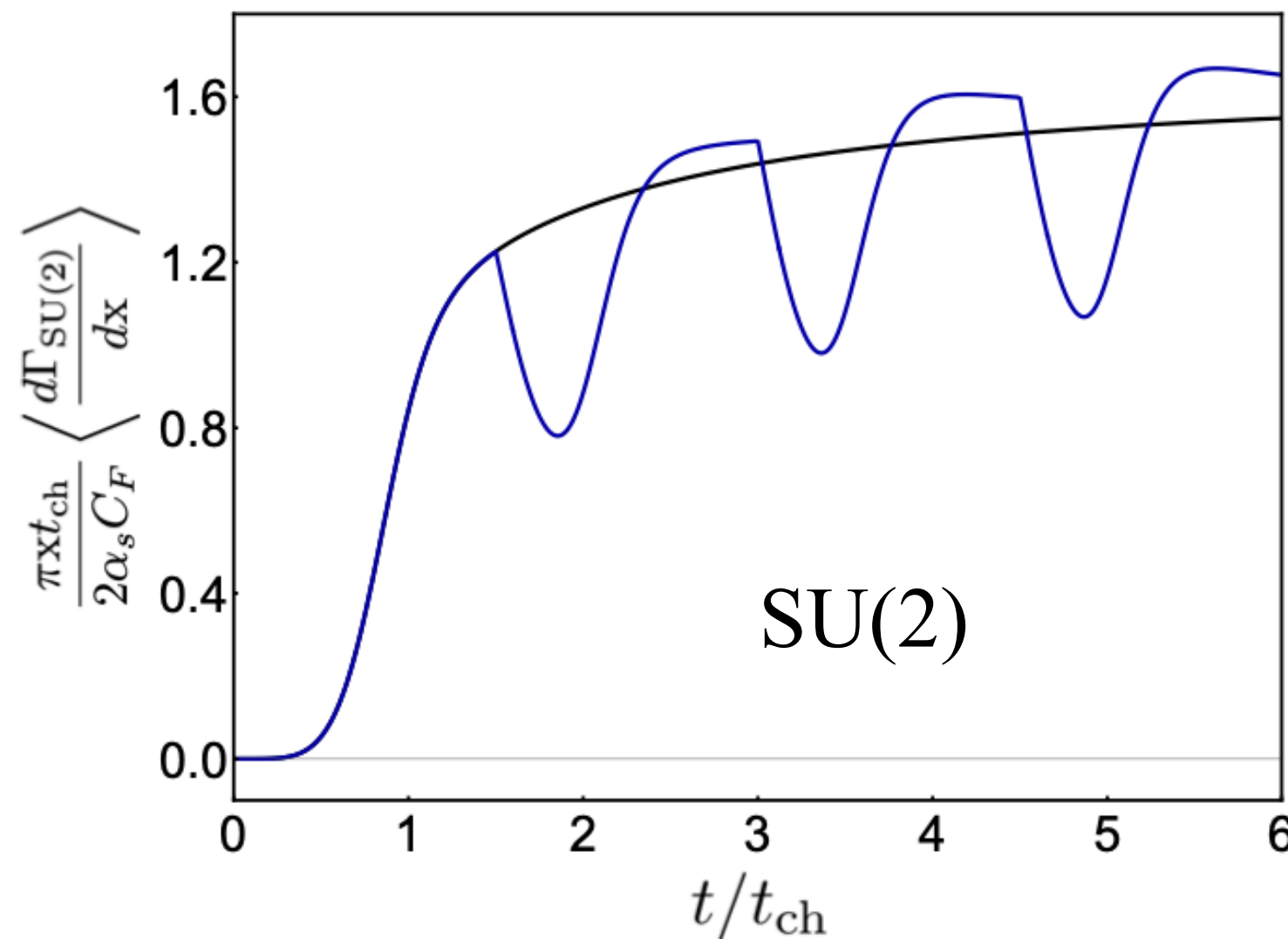
There is an emergent time scale : $t_{ch} = (24\omega/E^2)^{1/3}$ which can also be obtained from

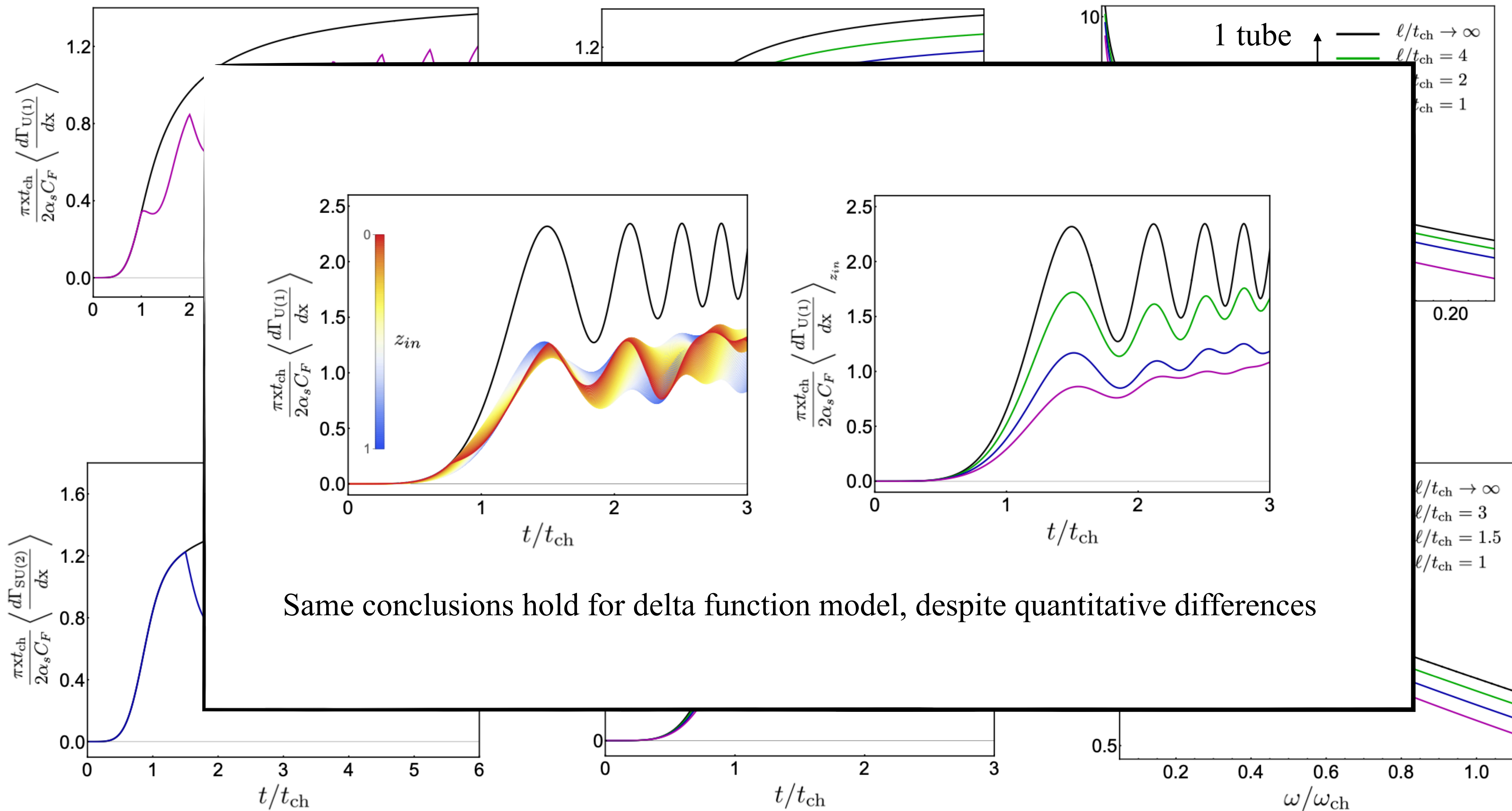
$$t_{ch} \sim \sqrt{\omega/\hat{q}(t_{ch})}$$

$$\langle \hat{q} \rangle_{E, z_{in}} = cE_0^2 l$$

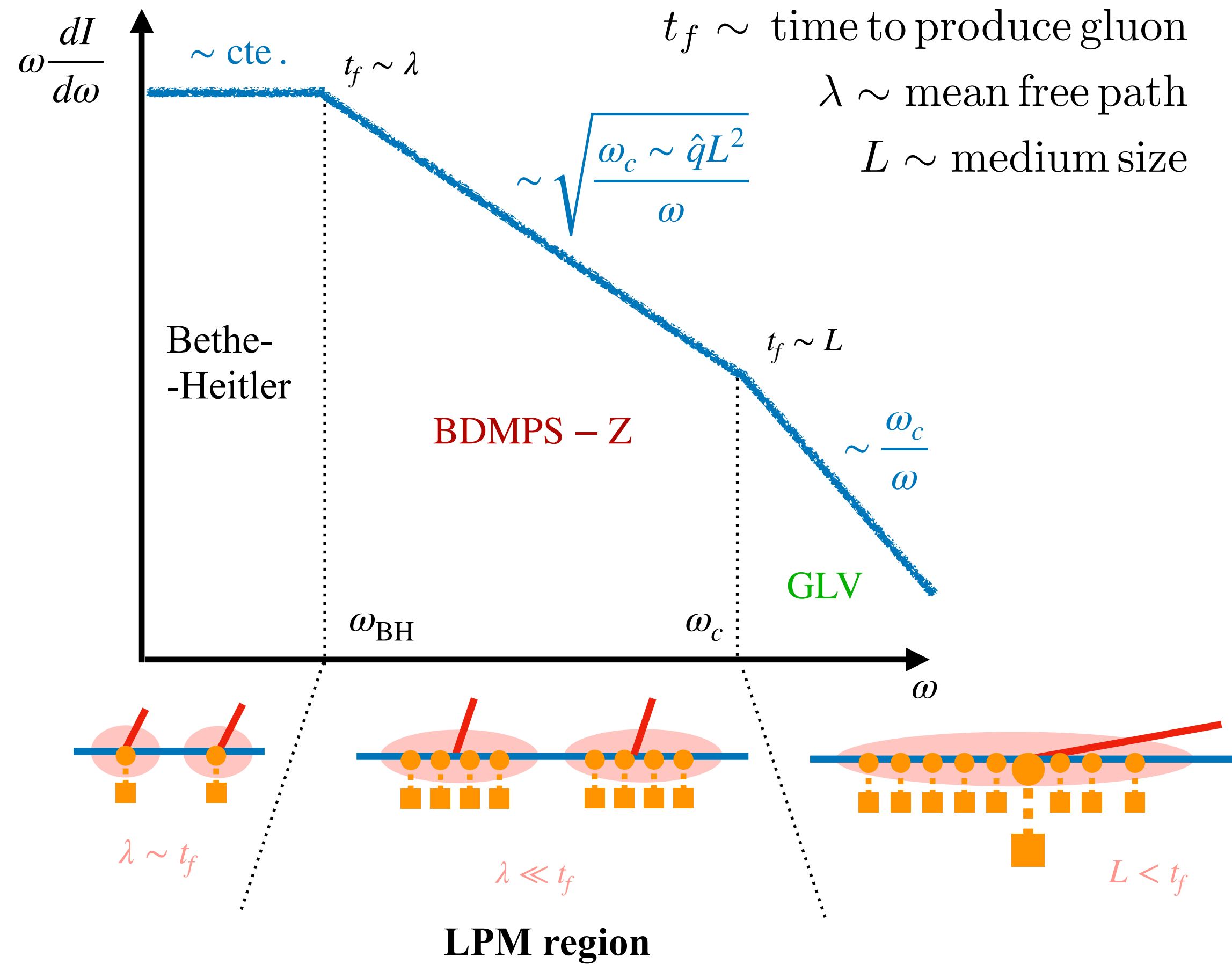


More flux tubes = longer formation time = lower rate

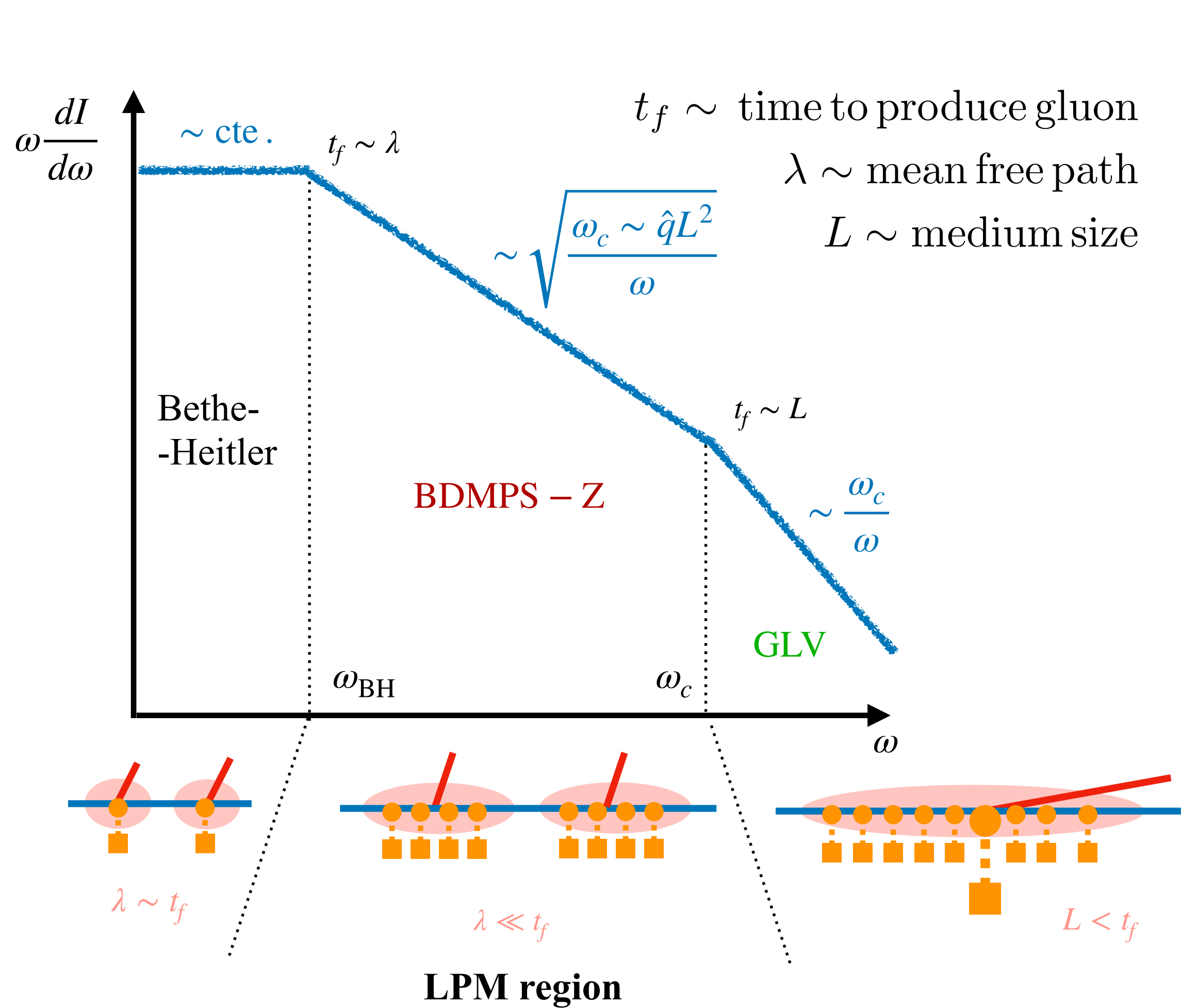




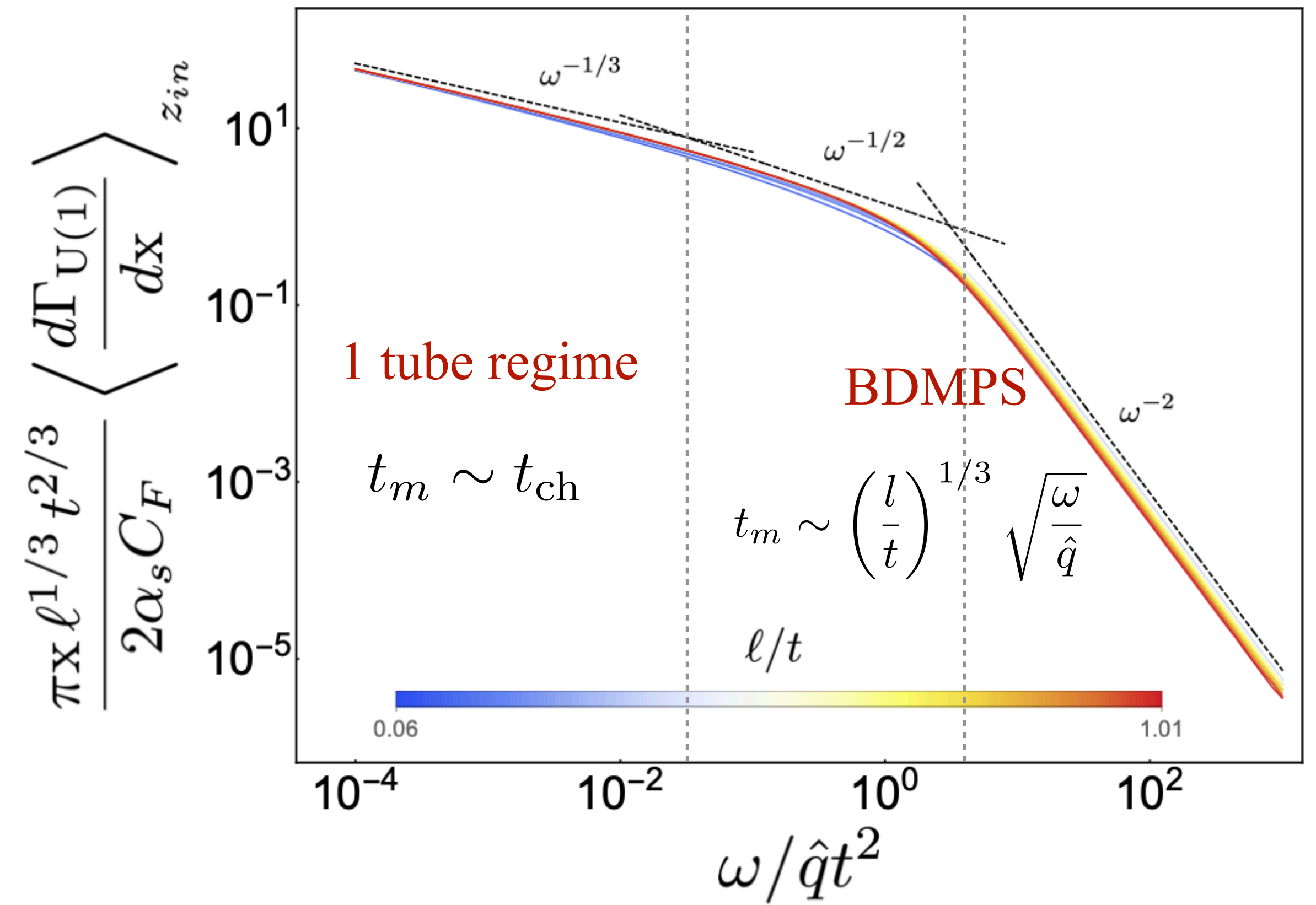
How different is this model from the standard multiple scattering picture ?



How different is this model from the standard multiple scattering picture ?



Fixed $\langle \hat{q} \rangle_{U(1)} = \frac{1}{8\pi^2} E_0^2 \ell$



$t \left\langle \frac{d\Gamma_{U(1)}}{dx} \right\rangle_{z_{in}} \sim (t/\ell)^{1/3} (\hat{q}t^2/\omega)^{1/3} F(\omega/(\hat{q}t^2), l/t)$

Some ideas



How about existing measurements ? What insights do they provide ? Do we need new tools ?

- Most existing jet substructure observables should be insensitive to this physics (not differential enough)
- There are proposals to go beyond, but not fully studied; can not get over full overall evolution effects

For any physical “model” what we want/need: 1) Azimuthal/polarized observable



2) Picks up time constrained evolution



Boosted W s time delayed decays + azimuthal/polarized dependent observable ?