

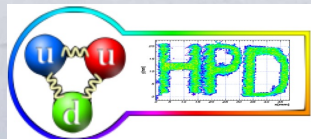


MINISTERUL CERCETĂRII,  
INOVĂRII ȘI DIGITALIZĂRII



*An overview  
on some global trends observed in heavy ion collisions  
based on experimental results from AGS up to LHC energies  
and  
on similarities between pp and Pb-Pb collisions at LHC*

*Results obtained in collaboration with A. Pop, C. Andrei, I. Berceanu, A. Lindner, M. Tarzila*

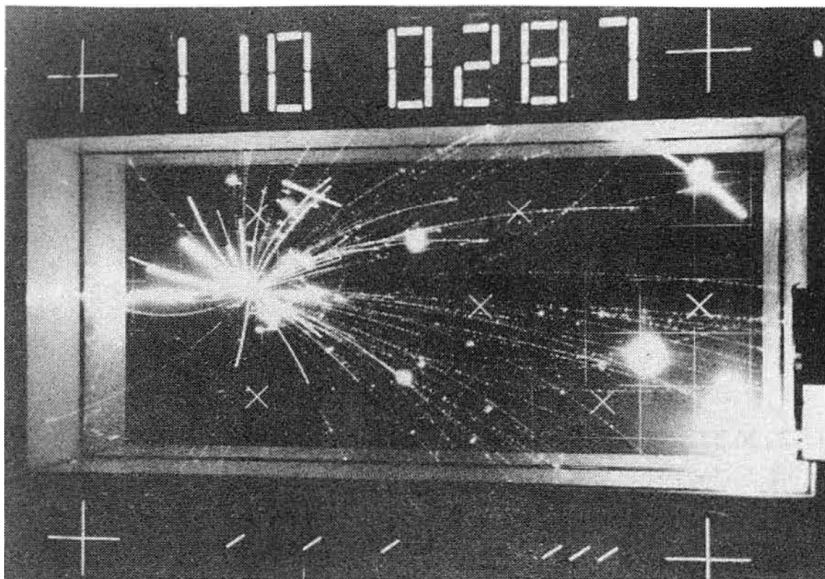


# Outline

- *Introduction*
- $\langle p_T \rangle / [(dN/dy)/S_{\perp}]^{1/2}$  *centrality and collision energy dependence*
- $[(dN/dy)/S_{\perp}]^{1/2}$  *scaling*
- $\langle dE_T/dy \rangle / \langle dN/dy \rangle - \langle dN/dy \rangle / S_{\perp}$  *correlation*
- *The slope of  $\epsilon_{Bj} \cdot \tau - \langle dN/dy \rangle / S_{\perp}$  correlation - energy dependence*
- $(dN/dy)^{\text{(strange and multi strange)}} / (dN/dy) - (dN/dy) / S_{\perp}$  *correlation*
- *Similar studies for pp collisions and comparison with Pb-Pb collisions*
- *Concluding remark*

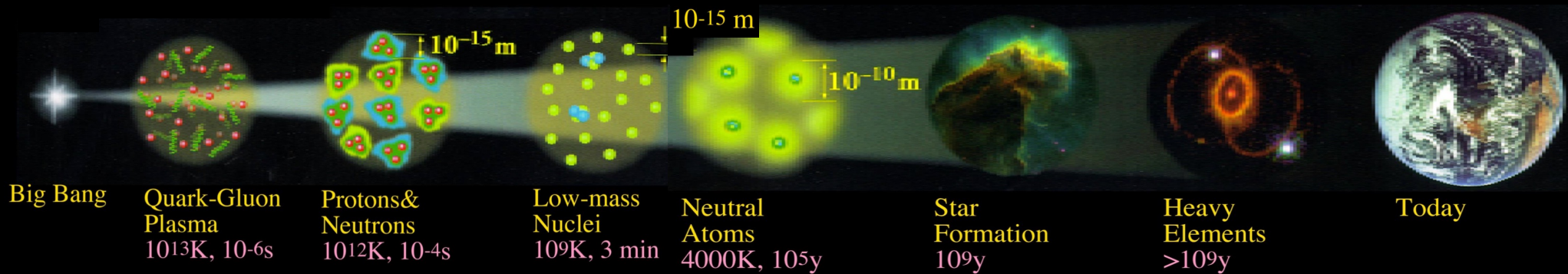
## *50<sup>th</sup> anniversary of high energy heavy-ion*

- *The high-energy heavy-ion program at LBL has started in summer 1974 (CERN Courier, June 1974)*
- *A University of Frankfurt group has exposed their AgCl detectors to various heavy-ion beams at energies from 250 MeV/A to 2.1 GeV/A. The observed peaks in the angular distributions of light fragments that moved with beam energy in a manner suggestive of these particles arising from shock waves, causing considerable excitement in the nuclear science community.*
- *After being used for several high energy experiments, the LBL streamer chamber used in the collision of 1.8-GeV/nucleon Ar on a lead oxide target, evidenced charged particle multiplicities of over 100 in such reactions.*



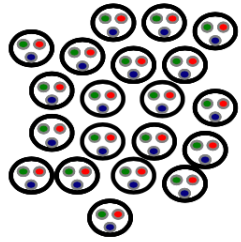
<https://escholarship.org/uc/item/8bw3436f>

# Could we unravel the History of the Universe

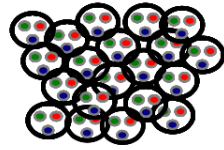


**based on experiments  
in terrestrial laboratories ?**

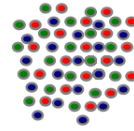
# How to produce extreme states of nuclear matter in terrestrial laboratories ?



Strongly Bound Clusters  
Hadrons



Phase transition



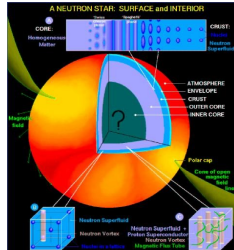
Weakly interacting  
Quarks and Gluons

A phase transition is expected at:

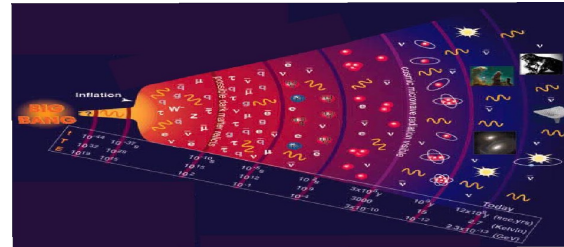
$$\rho_B \sim \Lambda_{QCD}^3 \sim 1 \text{fm}^{-3}$$

$$T \sim \Lambda_{QCD} \sim O(10^{12} \text{ K})$$

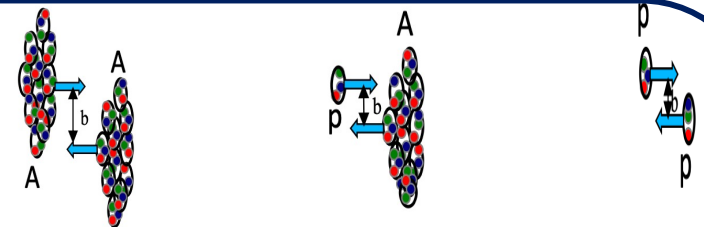
~ 16 km  
~ 10<sup>16</sup> sec



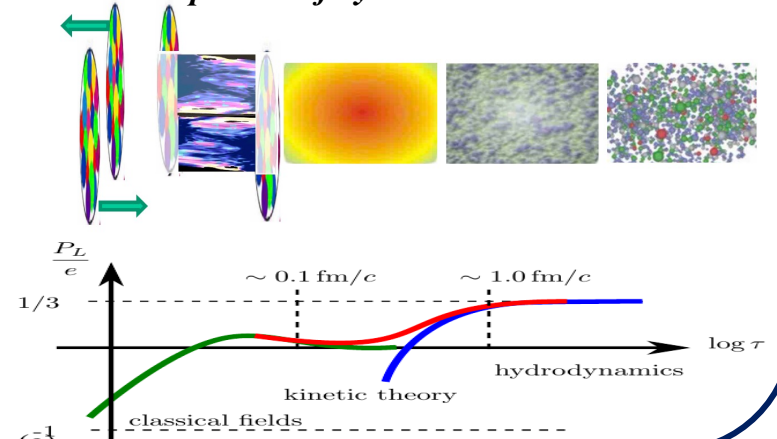
~ 10 km  
~ 10<sup>-6</sup> sec



~ 6 fm  
~ 10<sup>-22</sup> sec



Snap shots of dynamical evolution



A. Mazeliauskas, Nucl.Phys. A 00(2018)1  
QM 2018

# Large scale facilities

**LHC:** Collider  
Pb+Pb @5020GeV/A



**RHIC:** Collider  
Au+Au @ 200GeV/A



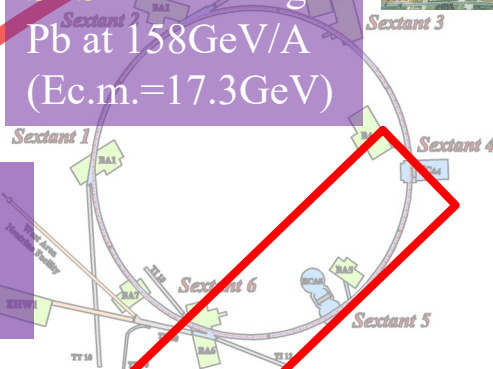
Hotter  
Denser  
Longer  
Bigger



?

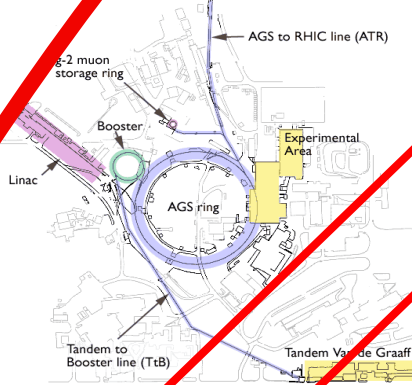
Click on the area of interest

**SPS:** Fixed Target  
Pb at 158GeV/A  
(Ec.m.=17.3GeV)

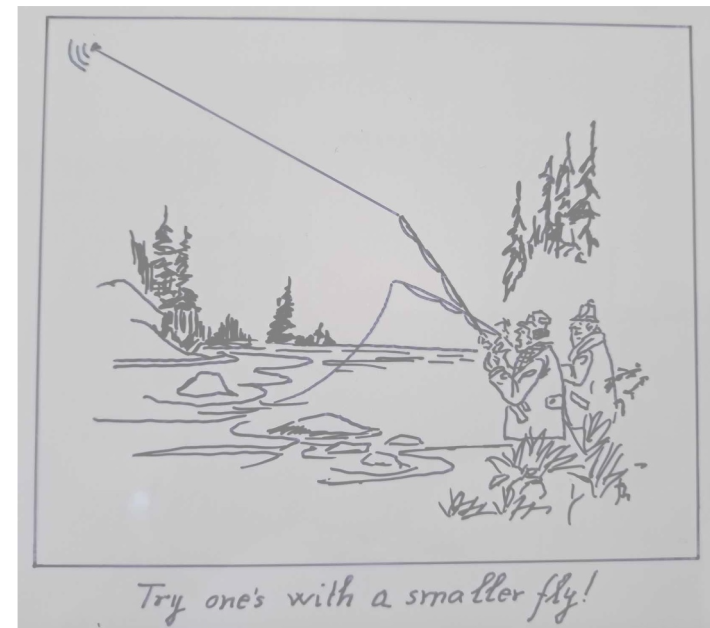


???

**AGS:** Fixed Target  
Au at 11.7GeV/A  
(Ec.m.=4.86GeV)



BES



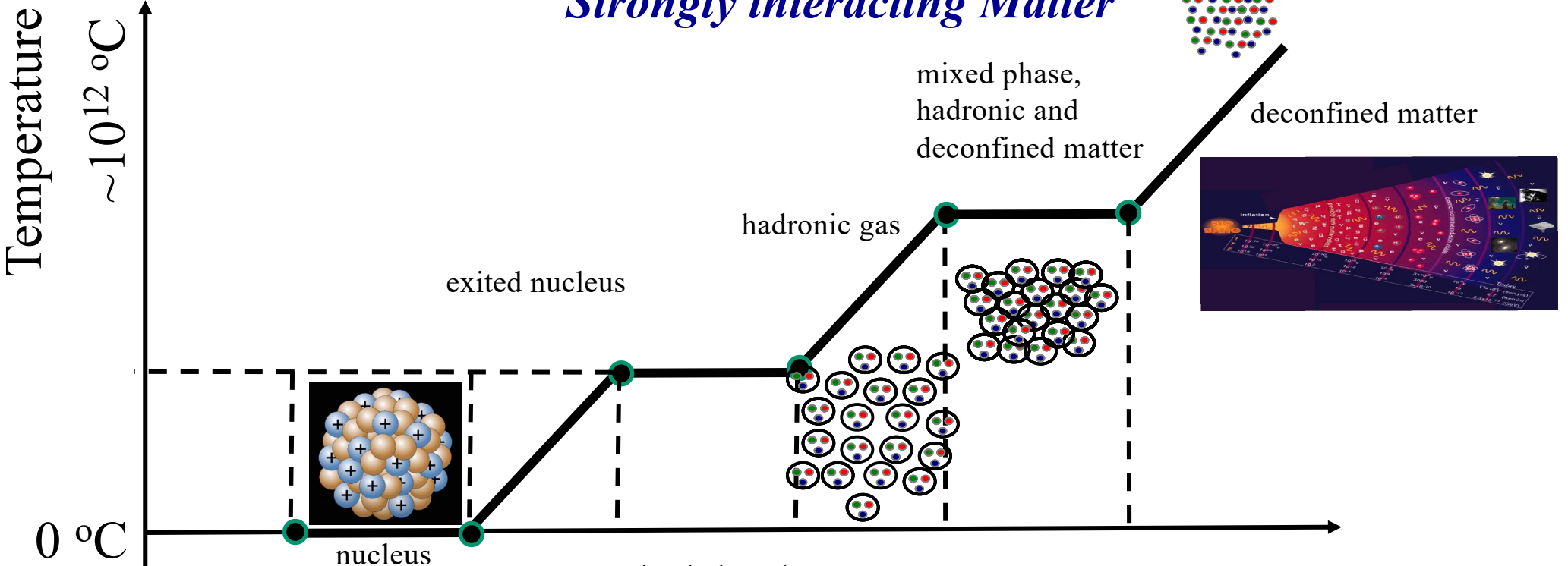
**Bevalac**  
Fixed Target  
1-2GeV/A



SIS 18

# Physics motivation

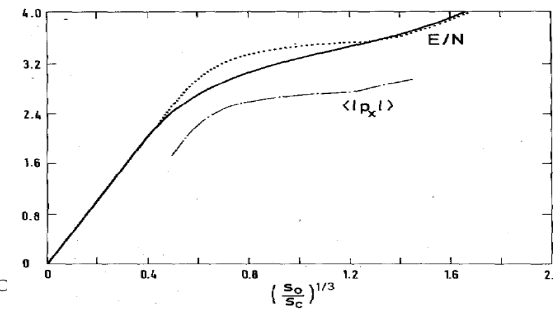
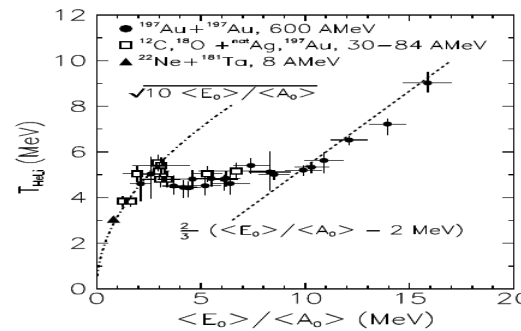
## Strongly interacting Matter



them.”) The elder Bohr, as a young graduate student in 1905, had written a prize-winning paper on the vibration of liquid drops of water. Seventy years later his son is being honored for work growing out of the liquid-drop picture.

mixed phased

thermal energy



*J.Pochodzalla et al.,  
 ALADIN Coll.,  
 arXiv:[nucl-ex]9607004*

*J.-P. Blaizot and J.-Y. Ollitrault,  
 Phys.Lett 191B(1987)21*

# Theory predictions

## String percolation

*T.S.Biro, H.B.Nielsen and J.Knoll, Nucl.Phys. B245(1984)449*  
*J.Dias de Deus and C. Pajares, Phys.Lett. B695(2011)211*  
*I. Bautista et al., Revista Mexicana de Fisica 65(2019)197*

$$\frac{dN}{dy} = F(\eta)\bar{N}^s\mu$$

$\eta \equiv (r_0/R)^2\bar{N}^s$  - transverse string density;  $\bar{N}^s$  - the average number of strings  
 $\mu$  - string multiplicity

$$F(\eta) \equiv \sqrt{\frac{1-e^{-\eta}}{\eta}}$$

$$\langle p_T^2 \rangle = \langle p_T^2 \rangle_1 / F(\eta) \quad \langle p_T^2 \rangle_1 \text{ - average string transverse momentum}$$

$$\sqrt{\langle p_T^2 \rangle} / \sqrt{\langle dN/dy \rangle / S_{\perp}} \sim 1 / \sqrt{(1-e^{-\eta})}$$

$$\langle p_T^2 \rangle / [(\langle dn/dy \rangle / S_{\perp})] \propto \langle p_T^2 \rangle_1 r_0^2 / \mu (1-e^{-\eta})$$

## CGC

*Local parton-hadron duality picture*  
*and dimensionality argument*

- *Y.L.Dokshitzer, V.A.Khoze and S.Troian, J.Phys.G 17 (1991) 1585*  
 - *T. Lappi, Eur.Phys.J. C71 (2011) 1699*  
 - *E. Levin and A.H. Rezaeian, Phys.Rev.D 83 (2011)114001*

$$\langle p_T \rangle / \sqrt{\langle dN/dy \rangle / S_{\perp}} \sim \frac{1}{n\sqrt{n}}$$

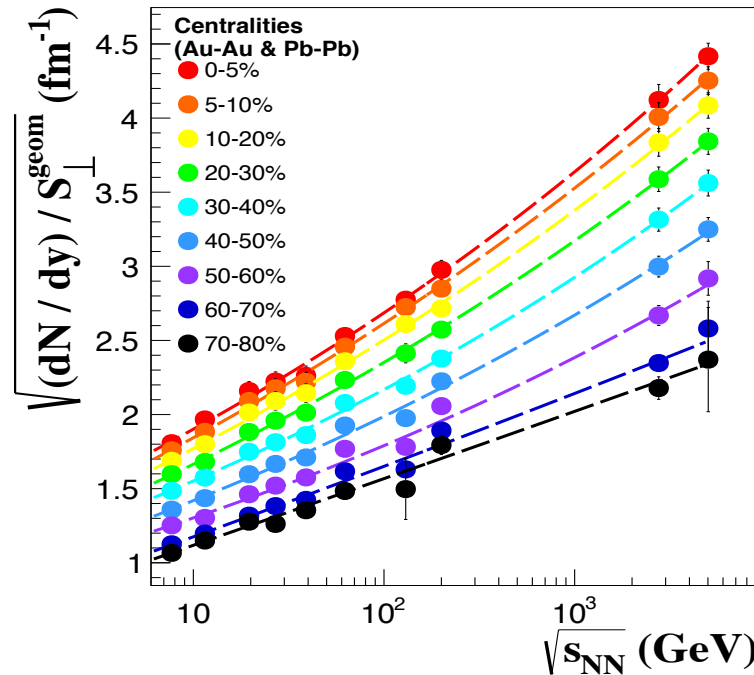
$n$  - no. of charged particles  
 from a gluon fragmentation



$$\langle p_T \rangle / \sqrt{\langle dN/dy \rangle / S_{\perp}}$$

decreases as a function of:

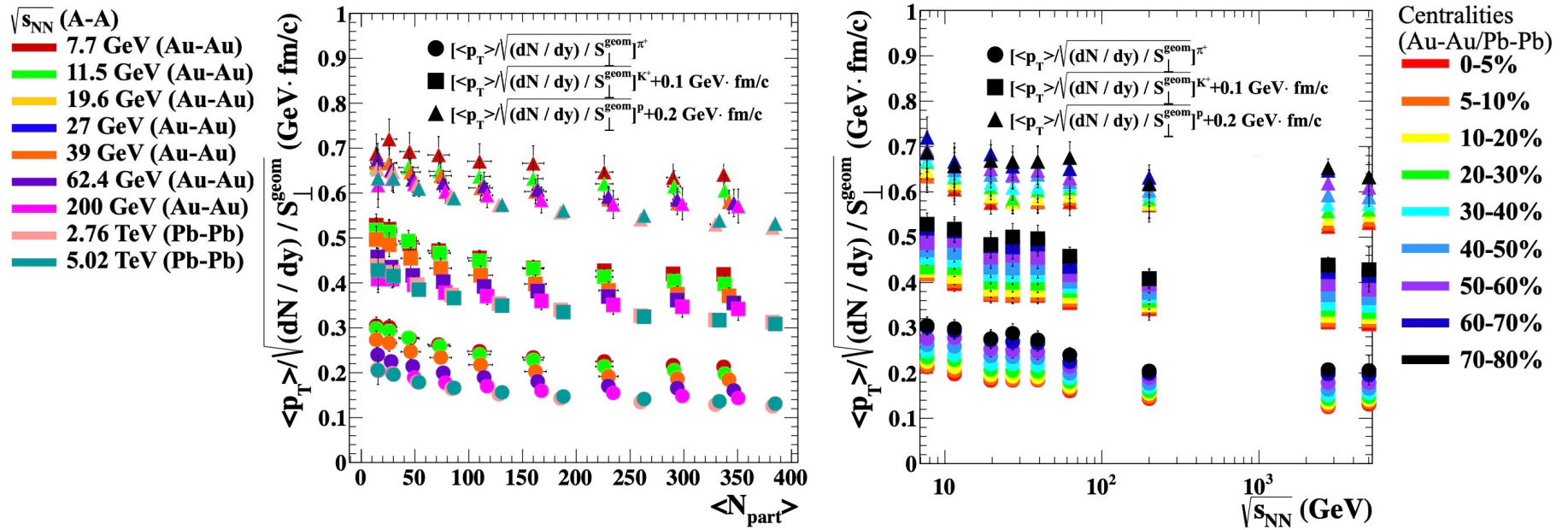
- collision energy
- centrality



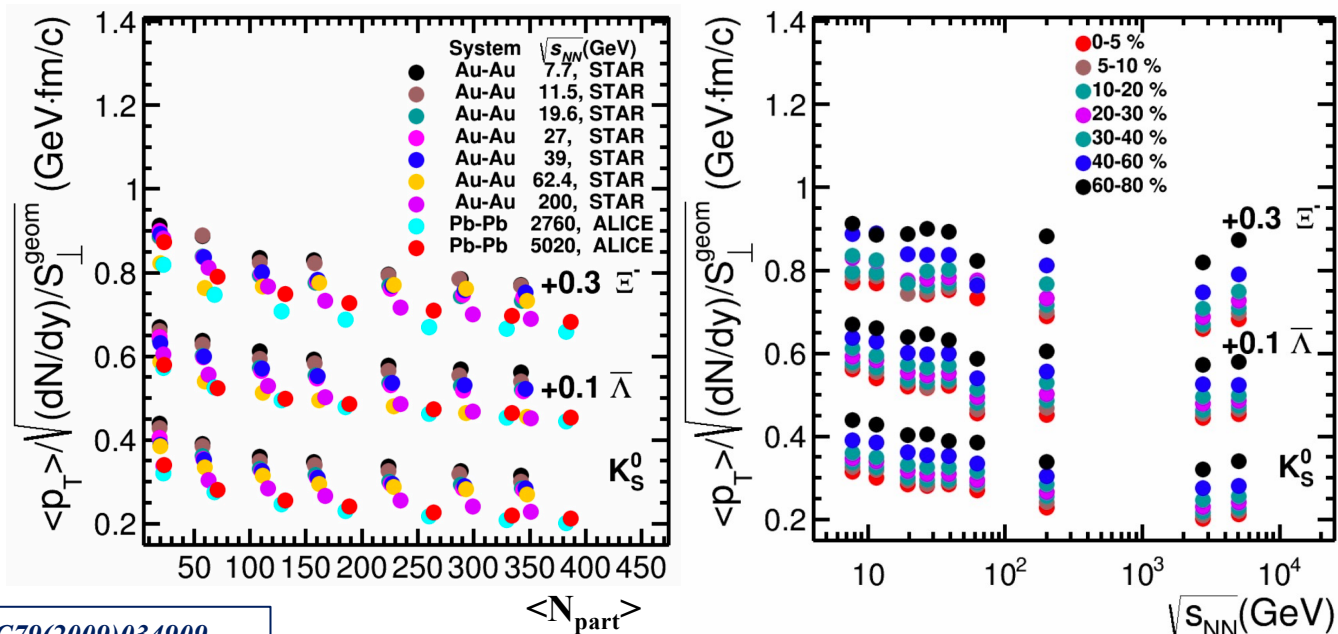
*M.Petrovici, A.Lindner and A.Pop, Phys. Rev. C 98(2018)024904*



# Experimental results



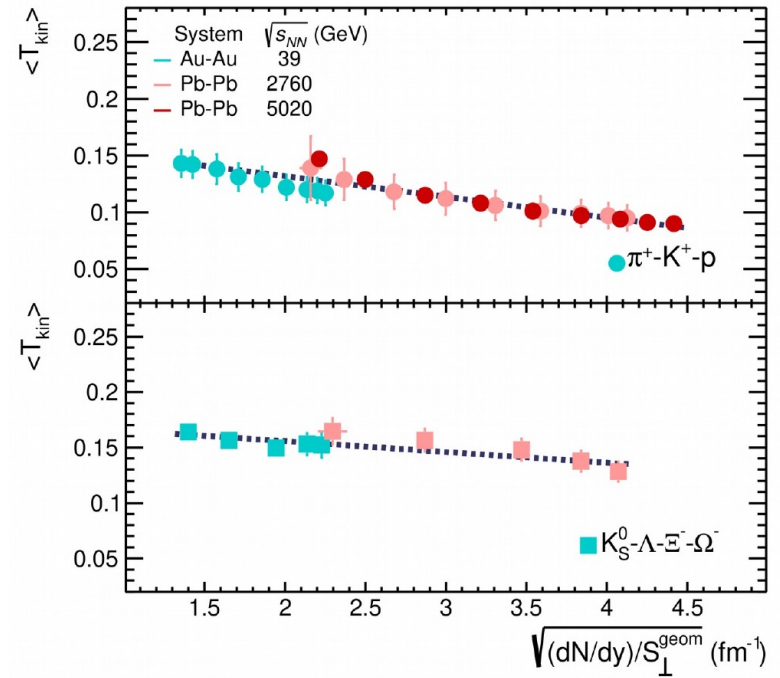
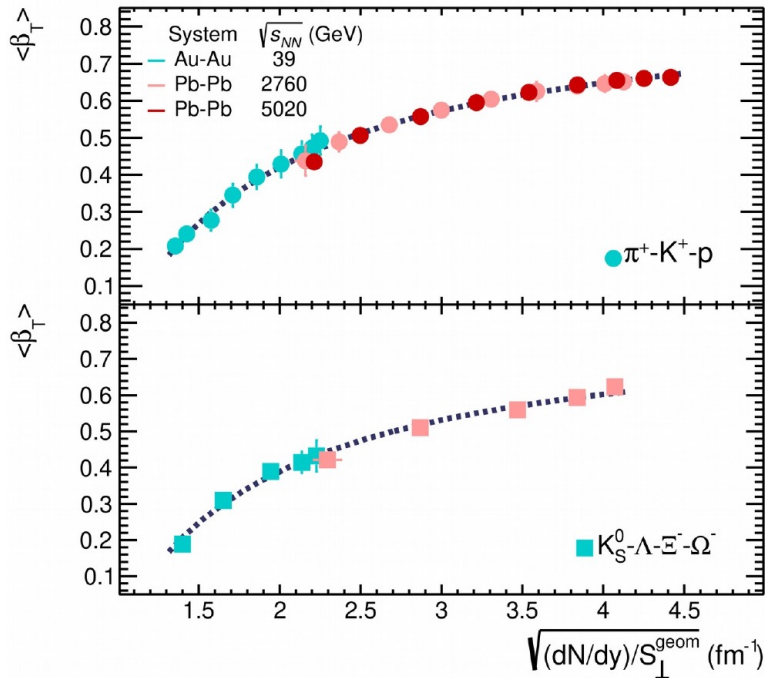
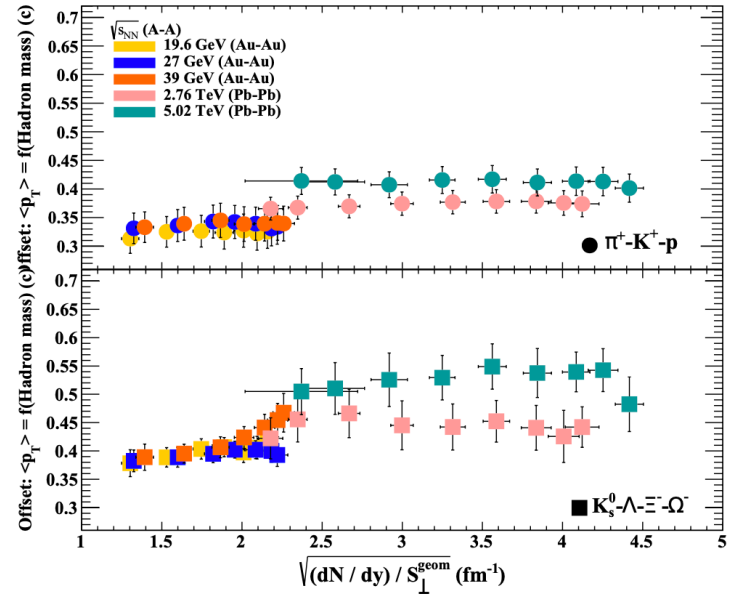
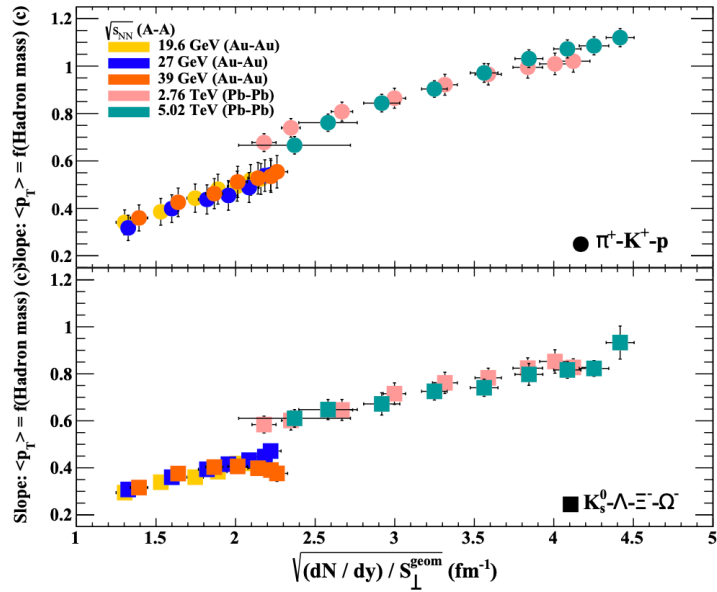
M. Petrovici, A. Lindner and A. Pop, Phys. Rev. C 98(2018)024904



STAR Collaboration, Phys. Rev. C 79(2009)034909  
 ALICE Collaboration, Phys. Rev. C 88(2013)044910  
 STAR Collaboration, Phys. Rev. C 96(2017)044904  
 ALICE Collaboration, Nucl. Phys. A 967(2017)421

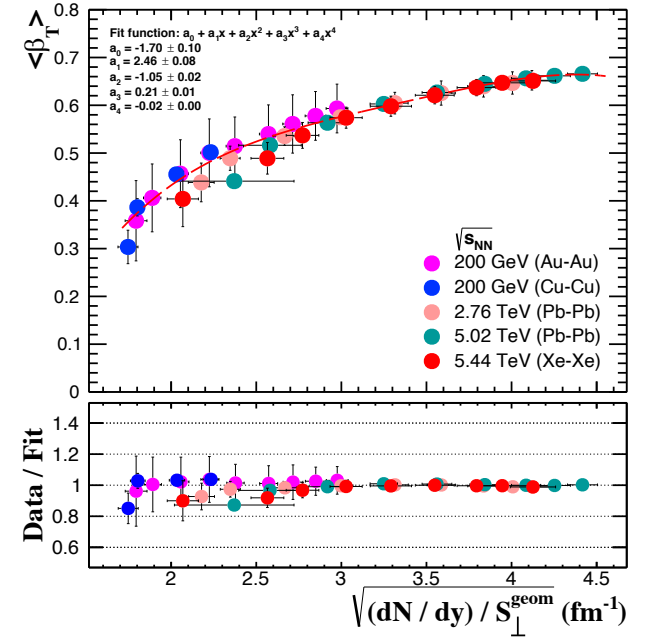
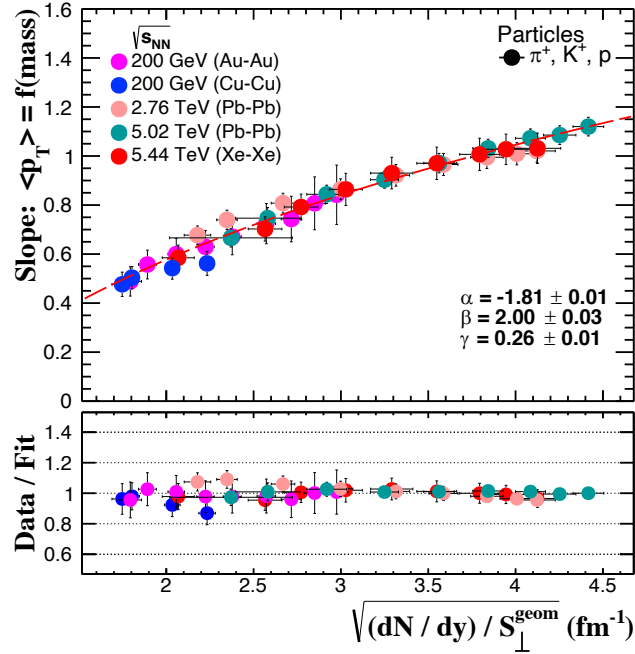
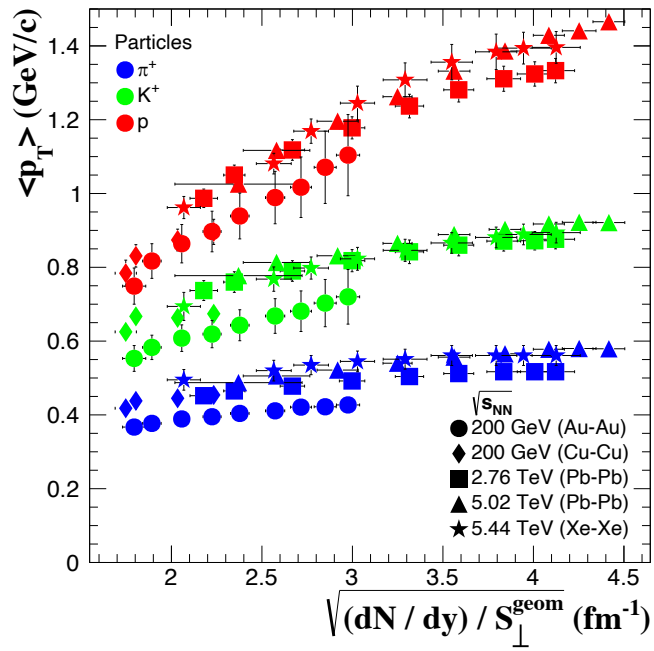
M. Petrovici and A. Pop, EuNPC 2022

# $[(dN/dy)/S_{\perp}]^{1/2}$ scaling



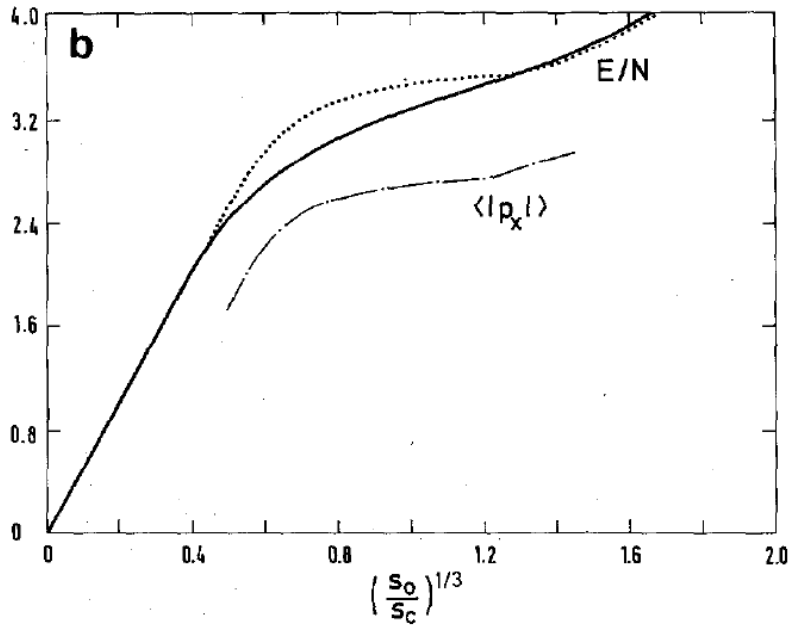
*M. Petrovici et al., Phys. Rev. C 98(2018)024904*  
*M. Petrovici and A. Pop, EuNPC 2022*  
*A. Pop and M. Petrovici, arXiv:2402.19115[hep-ph]*

# $[(dN/dy)/S_{\perp}]^{1/2}$ scaling



M. Petrovici, A. Lindner and A. Pop, AIP Conf.Proc. 2076 (2019) 1, 040001

# Signature for phase transition ?



*J.-P. Blaizot and J.-Y. Ollitrault,  
Phys.Lett 191B(1987)21*

$$E/N \sim \epsilon/s = E_{fo}/S_{fo} ; s(T_0) = a(1/R_0^3)(dN/dy)$$

$$\frac{dN}{dy} \simeq \frac{3}{2} \frac{dN}{dy}^{(\pi^+ + \pi^-)} + 2 \frac{dN}{dy}^{(p + \bar{p}, \Xi^- + \bar{\Xi}^+)} + \frac{dN}{dy}^{(K^+ + K^-, \Lambda + \bar{\Lambda}, \Omega^- + \bar{\Omega}^+)} + 2 \frac{dN}{dy}^{K_S^0} + 2 \frac{dN}{dy}^{(\Sigma^+ + \Sigma^-)}$$

$$\frac{dE_T}{dy} \simeq \frac{3}{2} \left( \langle m_T \rangle \frac{dN}{dy} \right)^{(\pi^+ + \pi^-)} + 2 \left( \langle m_T \rangle \frac{dN}{dy} \right)^{(p + \bar{p}, \Xi^- + \bar{\Xi}^+)} + \left( \langle m_T \rangle \frac{dN}{dy} \right)^{(K^+ + K^-, \Lambda + \bar{\Lambda}, \Omega^- + \bar{\Omega}^+)} + 2 \left( \langle m_T \rangle \frac{dN}{dy} \right)^{K_S^0} + 2 \left( \langle m_T \rangle \frac{dN}{dy} \right)^{(\Sigma^+ + \Sigma^-)}$$

$$\langle m_T \rangle - > \langle m_T \rangle - m_0 \quad - \text{ for baryons}$$

for RHIC energies:

$$\langle m_T \rangle - > \langle m_T \rangle + m_0 \quad - \text{ for antibaryons}$$

$$\langle m_T \rangle \quad - \text{ for other particles}$$

- AGS si SPS  
• S. Chatterjee et al., Advances in High Energy Physics 2015, 349013 (2015).

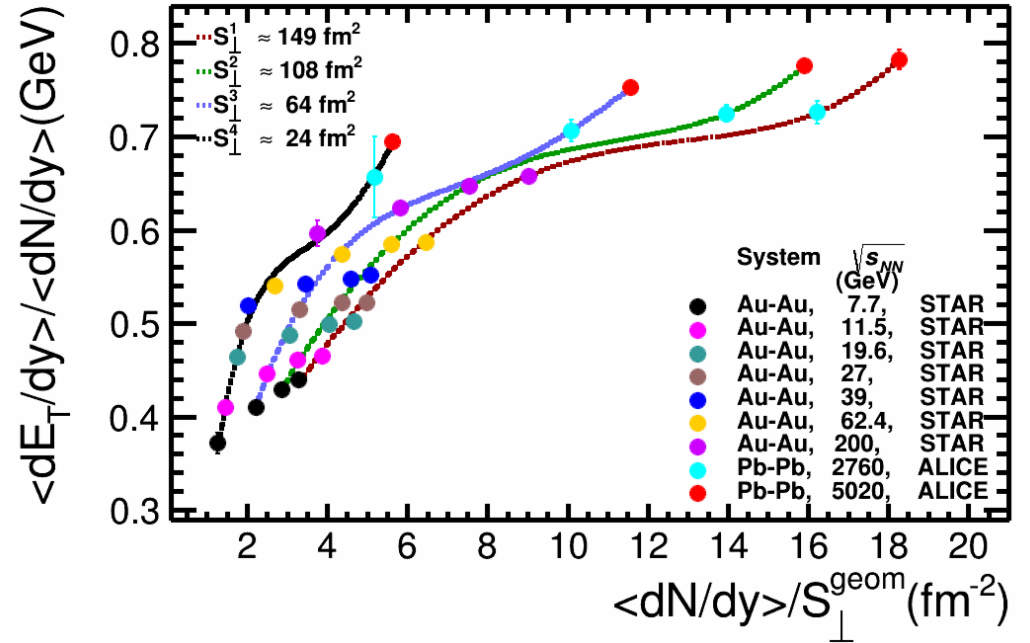
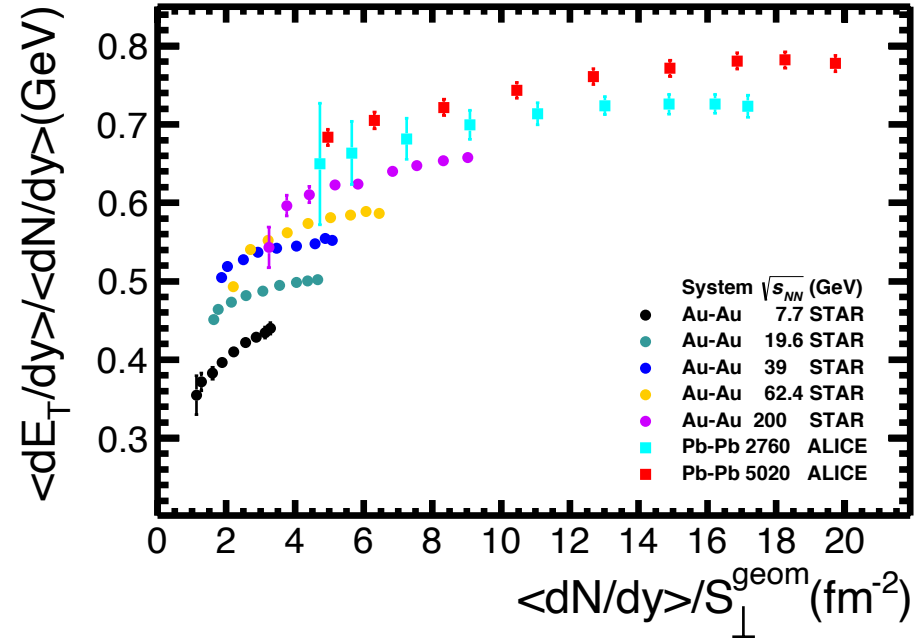
- BES  
• J. Adam et al. (STAR Collaboration), Phys. Rev. C 102, 034909 (2020).

- RHIC 62.4 GeV and 200 GeV  
• M. M. Aggarwal et al. (STAR Collaboration), Phys. Rev. C 83, 024901 (2011).  
• J. Adams et al. (STAR Collaboration), Phys. Rev. Lett. 98, 062301 (2007).  
• G. Agakishiev et al. (STAR Collaboration), Phys. Rev. Lett. 108, 072301 (2012).  
• L. Adamczyk et al. (STAR Collaboration), Phys. Rev. C 96, 044904 (2017). - RHIC, 62.4 si 200 GeV  
• B. I. Abelev et al. (STAR Collaboration), Phys. Rev. C 79, 034909 (2009). - ALICE 2.76 TeV

- ALICE 2.76 TeV  
• B. Abelev et al. (ALICE Collaboration), Phys. Rev. Lett. 111, 222301 (2013).  
• B. Abelev et al. (ALICE Collaboration), Phys. Lett. B 728, 216 (2014); 734, 409 (2014).  
• B. I. Abelev et al. (STAR Collaboration), Phys. Rev. C 79, 034909 (2009). - ALICE 2.76 TeV

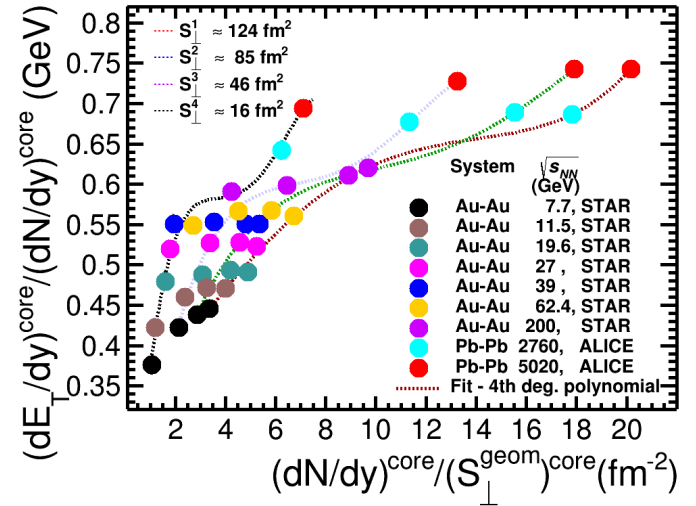
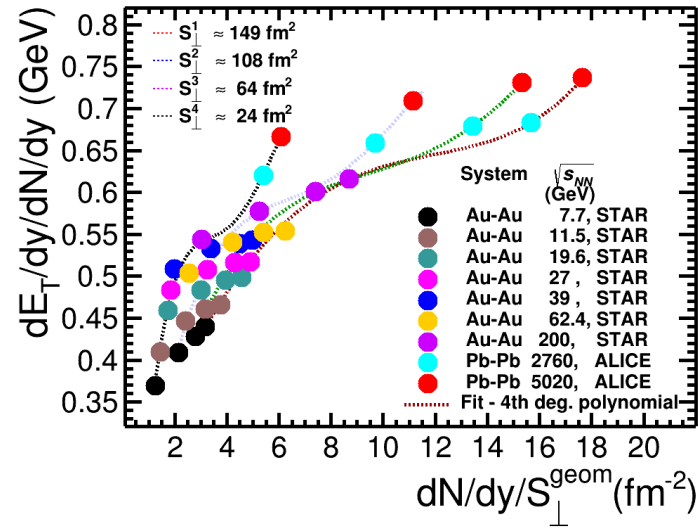
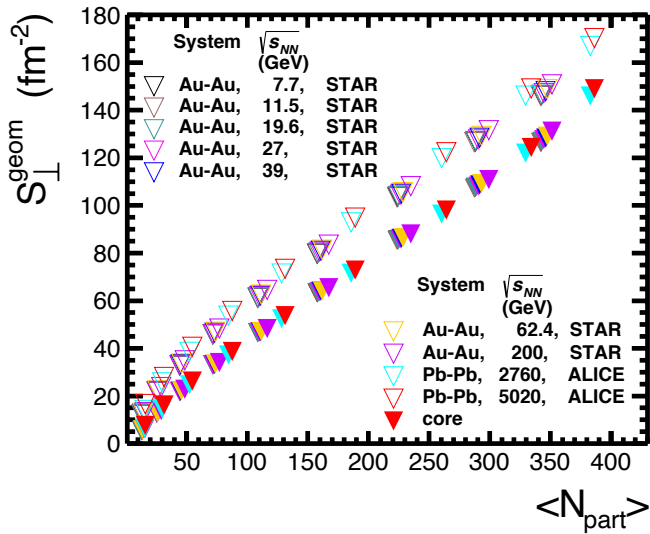
- ALICE 5.02 TeV  
• D. S. de Albuquerque, Ph.D. thesis (2019), CERN-THESIS-2019-135.  
• P. Kalinak for the ALICE Collaboration, European Physical Society Conference on High Energy Physics, 5-12 July 2017, Venice, Italy, PoS(EPS-HEP2017)168 (2017),  
• <https://pos.sissa.it/314/168/pdf>.  
• D. S. de Albuquerque for the ALICE Collaboration, Nucl. Phys. A 982, 823 (2019), XXVIIIth International Conference on Ultrarelativistic Nucleus-Nucleus Collisions (Quark Matter 2018).  
• L. Adamczyk et al. (STAR Collaboration), Phys. Rev. C 96, 044904 (2017). - RHIC, 62.4 si 200 GeV  
• B. Abelev et al. (ALICE Collaboration), Phys. Rev. C 88, 044910 (2013). - ALICE 5.02 TeV  
• S. Acharya et al. (ALICE Collaboration), Phys. Rev. C 101, 044907 (2020).

# $(dE_T/dy)/(dN/dy) - (dN/dy)/S_{\perp}$ correlation

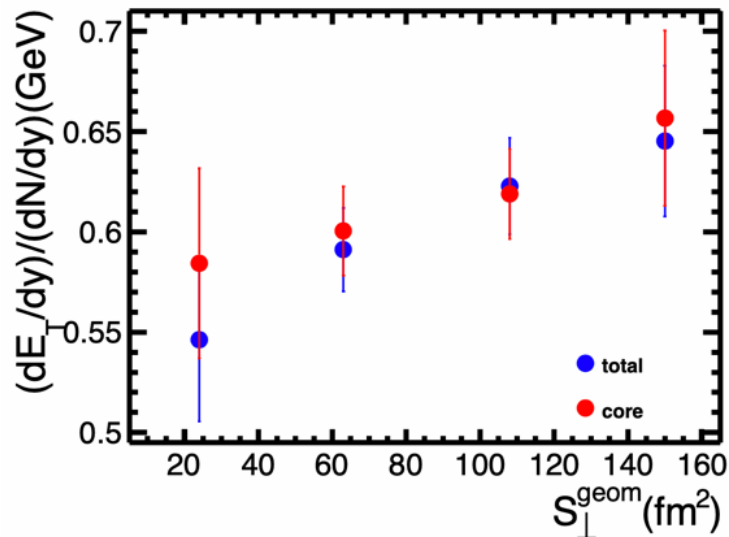


M. Petrovici and A. Pop, Phys.Rev. C107(2023)034913

# $(dE_T/dy)/(dN/dy) - (dN/dy)/S_{\perp}$ correlation - core-corona $\pi^{\pm}, K^{\pm}, p, pbar$ and their neutrals

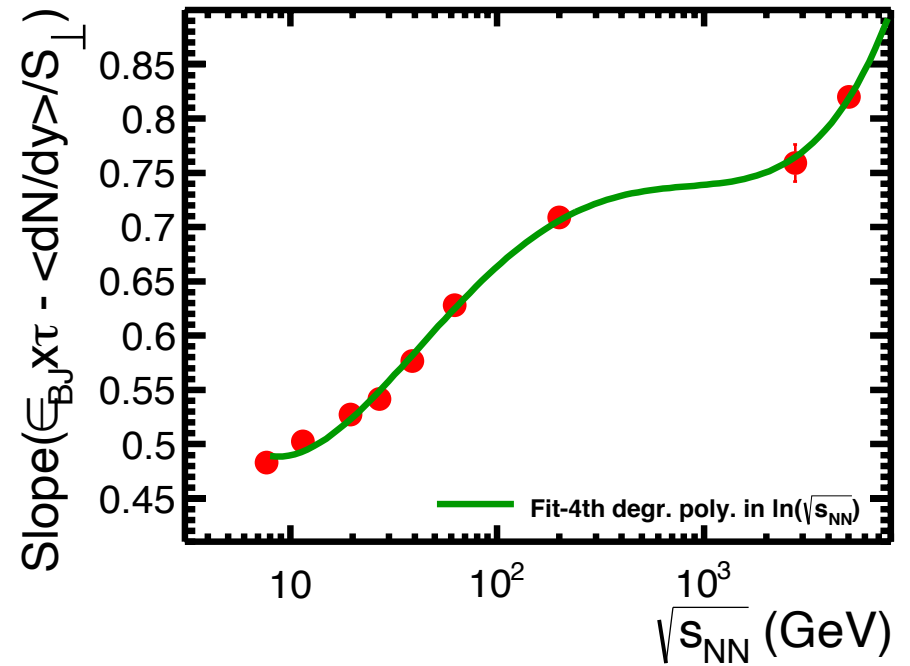
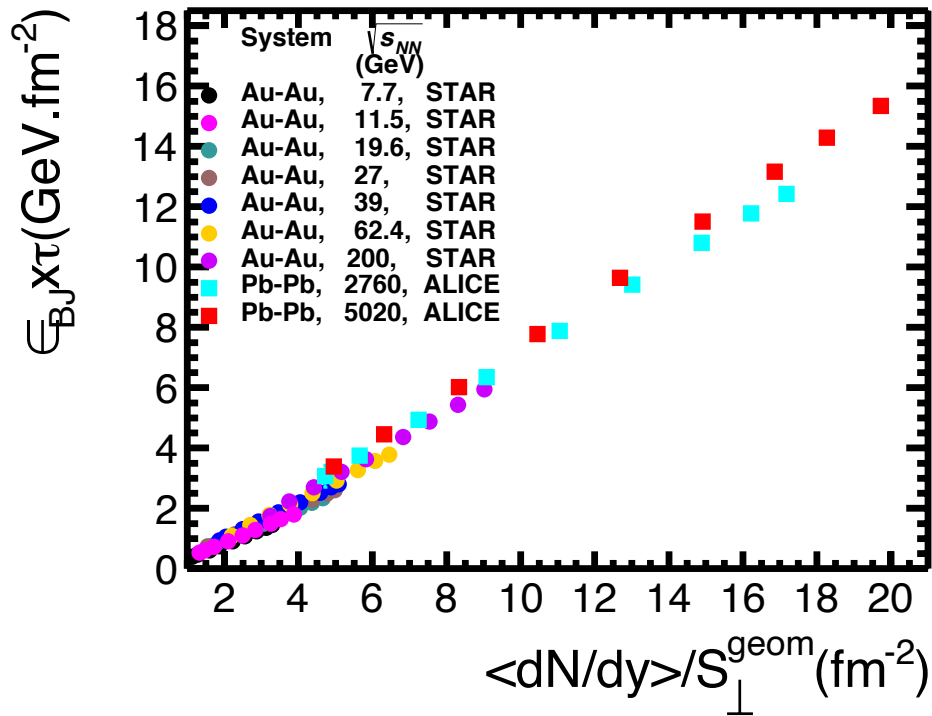


## Inflection points



# $\epsilon_{Bj} - (dN/dy)/S_{\perp}$ correlation for A-A - centrality dependence

$$\epsilon_{Bj} \cdot \tau = (dE_T/dy)/S_{\perp}$$

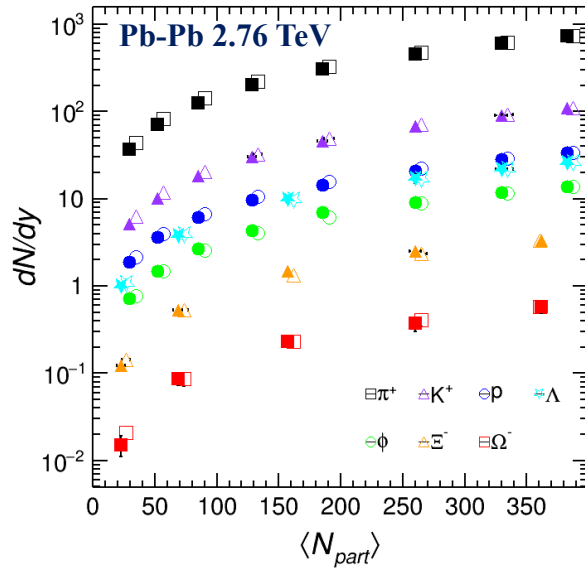


M.Petrovici and A.Pop, Phys.Rev. C107(2023)034913

# Strangeness production - smoking gun of deconfinement

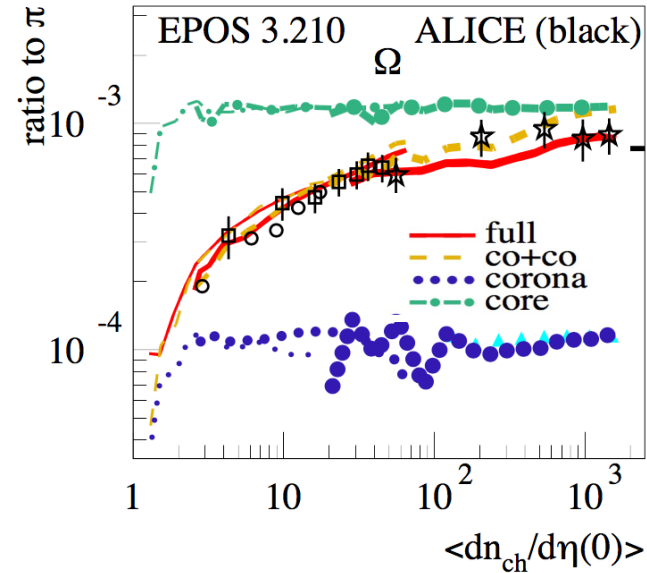
J.Rafelski and B.Muller, Phys.Rev.Lett. 48(1982)1066

$$\left(\frac{dN}{dy}\right)_i^{cen} = N_{part} [(1 - f_{core}) M_i^{ppMB} + f_{core} M_i^{core}] \quad (1)$$

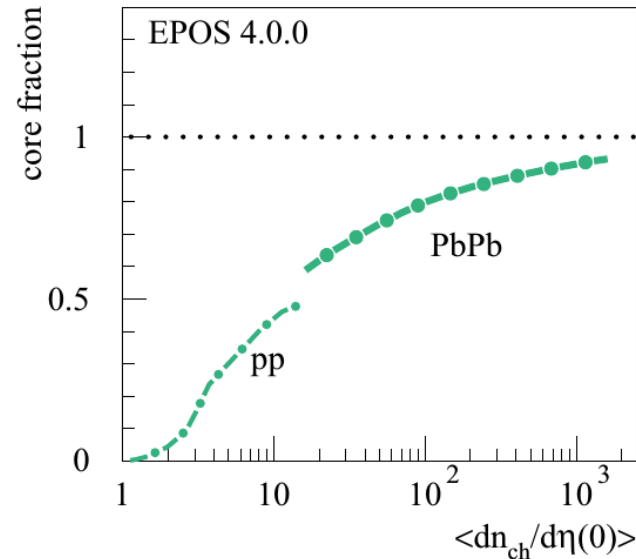


open symbols - Eq.1  
full symbols - exp. points

M. Petrovici et al., Phys.Rev. C96(2017)014908



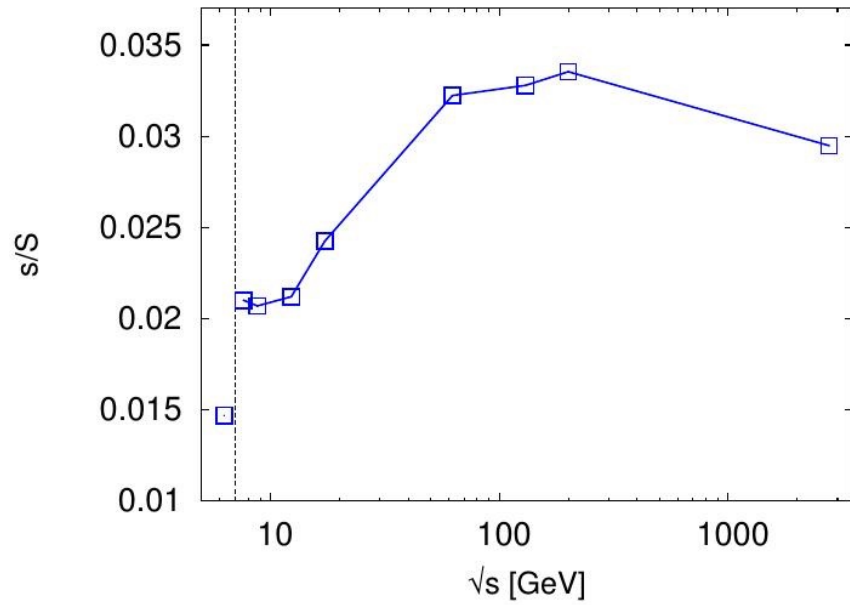
K. Werner, SQM 2017, July 10-15 2017, Utrecht



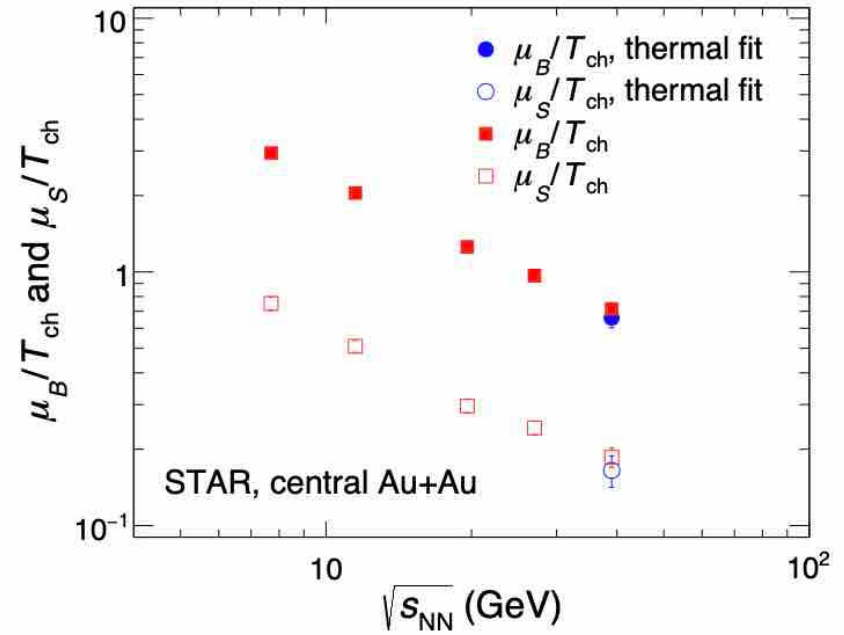
K. Werner, Phys.Rev. C109(2024)014910



# Strangeness production - smoking gun of deconfinement



J.Rafelski and M.Petran, *arXiv[nucl-th]1403.4036*

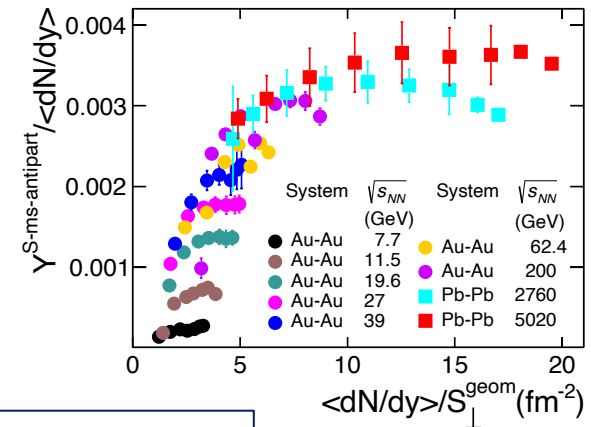
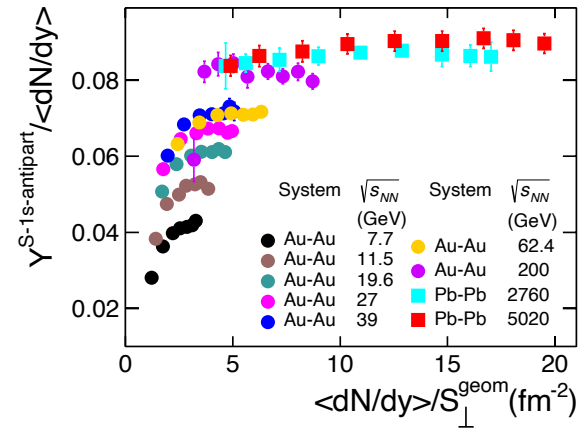
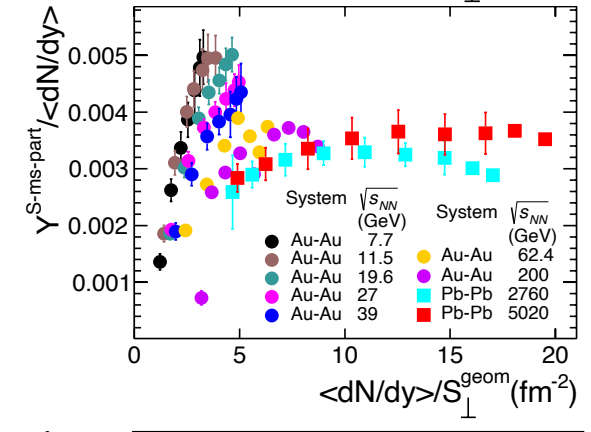
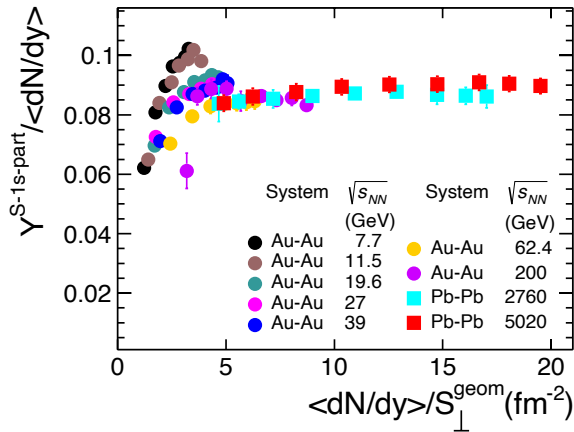
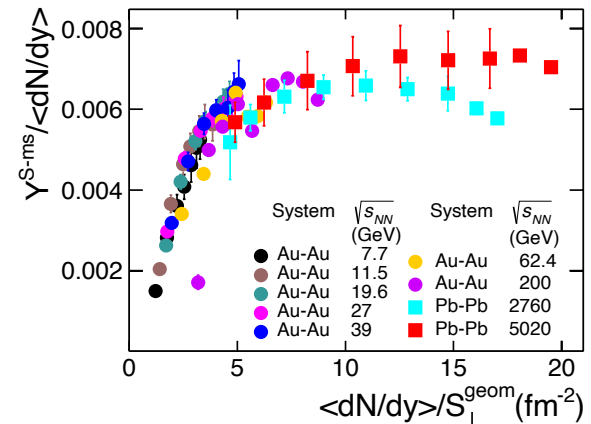
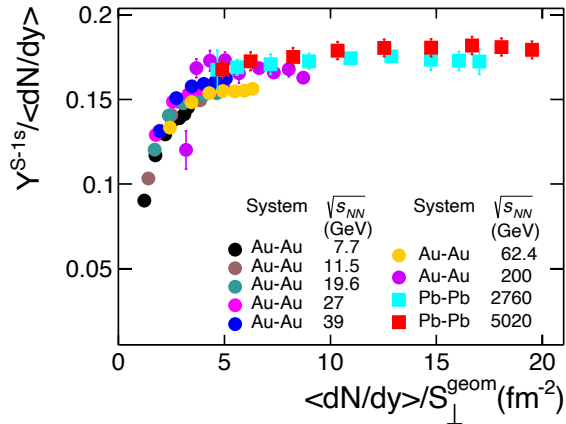


J.Adam et al, STAR Collaboration, *Phys.Rev. C102(2020)034909*

# $(dN/dy)_{\text{strange and multi strange}} / (dN/dy) - (dN/dy) / S_{\perp}$ correlation

$$Y^{1s} = \frac{dN^{1s}}{dy} = \frac{dN^{(K^+ + K^-)}}{dy} + 2 \frac{dN^{K_s^0}}{dy} + \frac{dN^{(\Lambda + \bar{\Lambda})}}{dy} + 2 \frac{dN^{(\Sigma^- + \bar{\Sigma}^+)}}{dy}$$

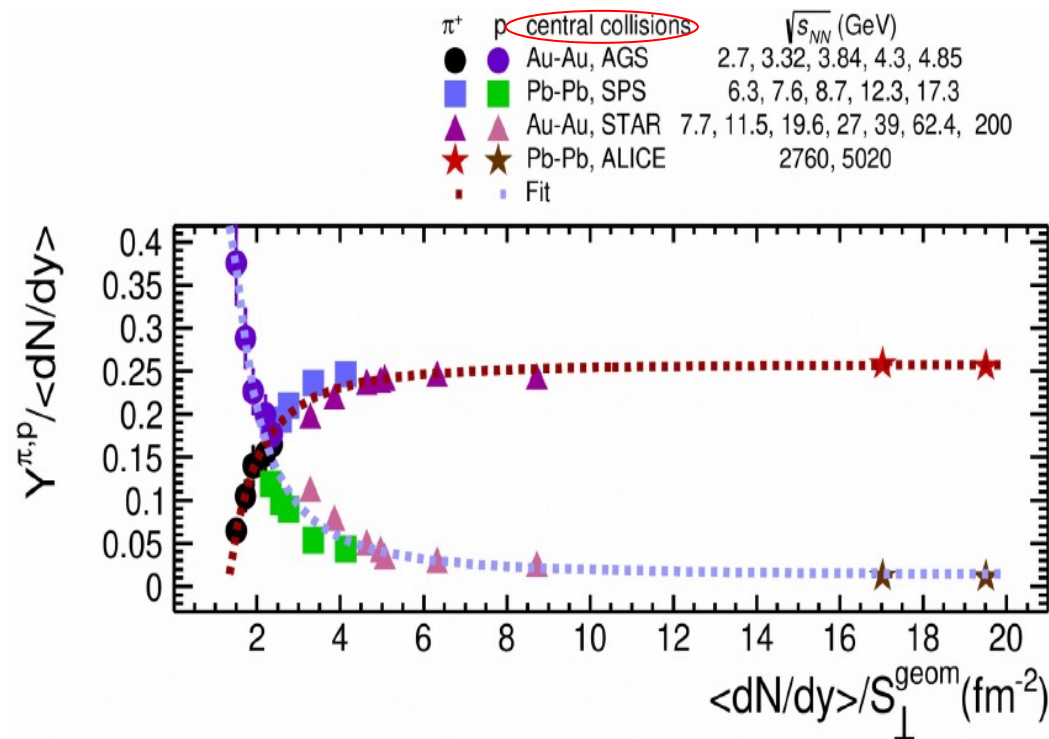
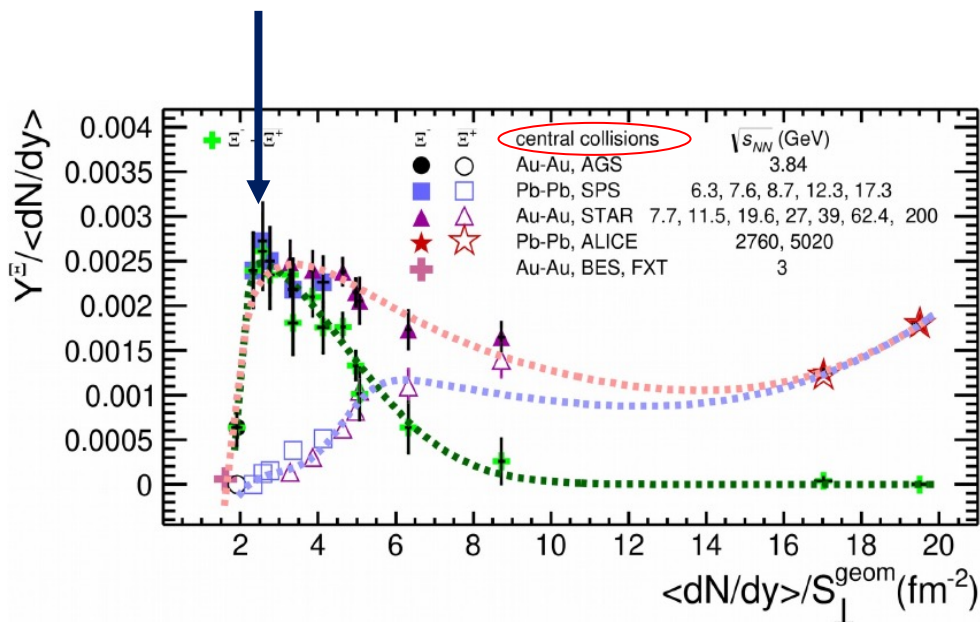
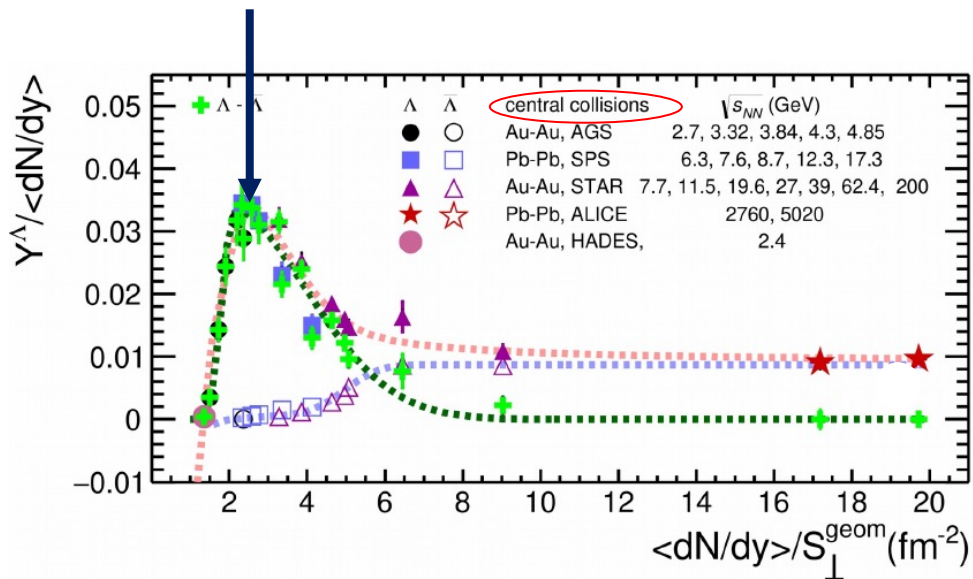
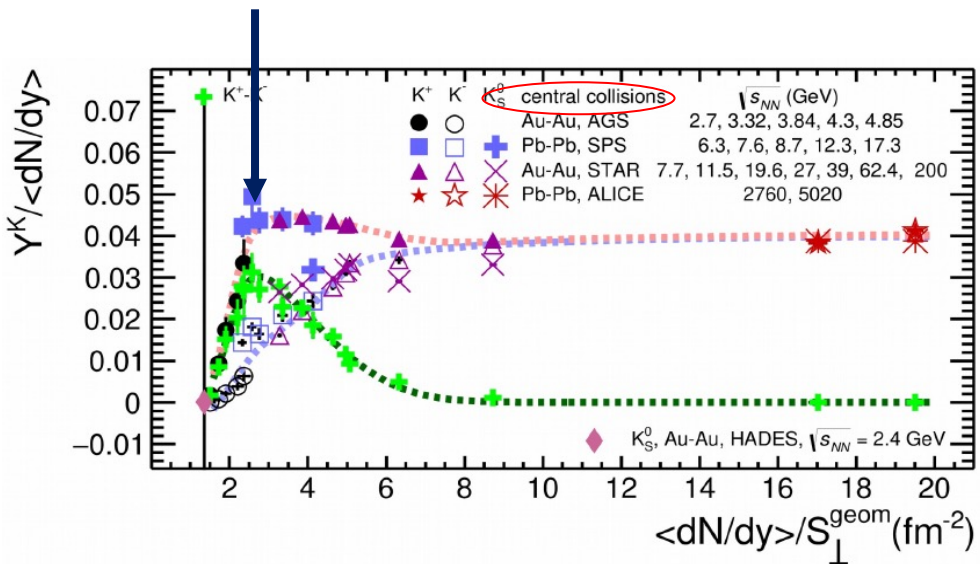
$$Y^{ms} = \frac{dN^{ms}}{dv} = \frac{dN^{(\Omega^- + \bar{\Omega}^+)}}{dv} + 2 \frac{dN^{(\Xi^- + \bar{\Xi}^+)}}{dv}$$



M. Petrovici and A. Pop, EuNPC 2022  
A. Pop and M. Petrovici, arXiv:2402.19115[hep-ph]

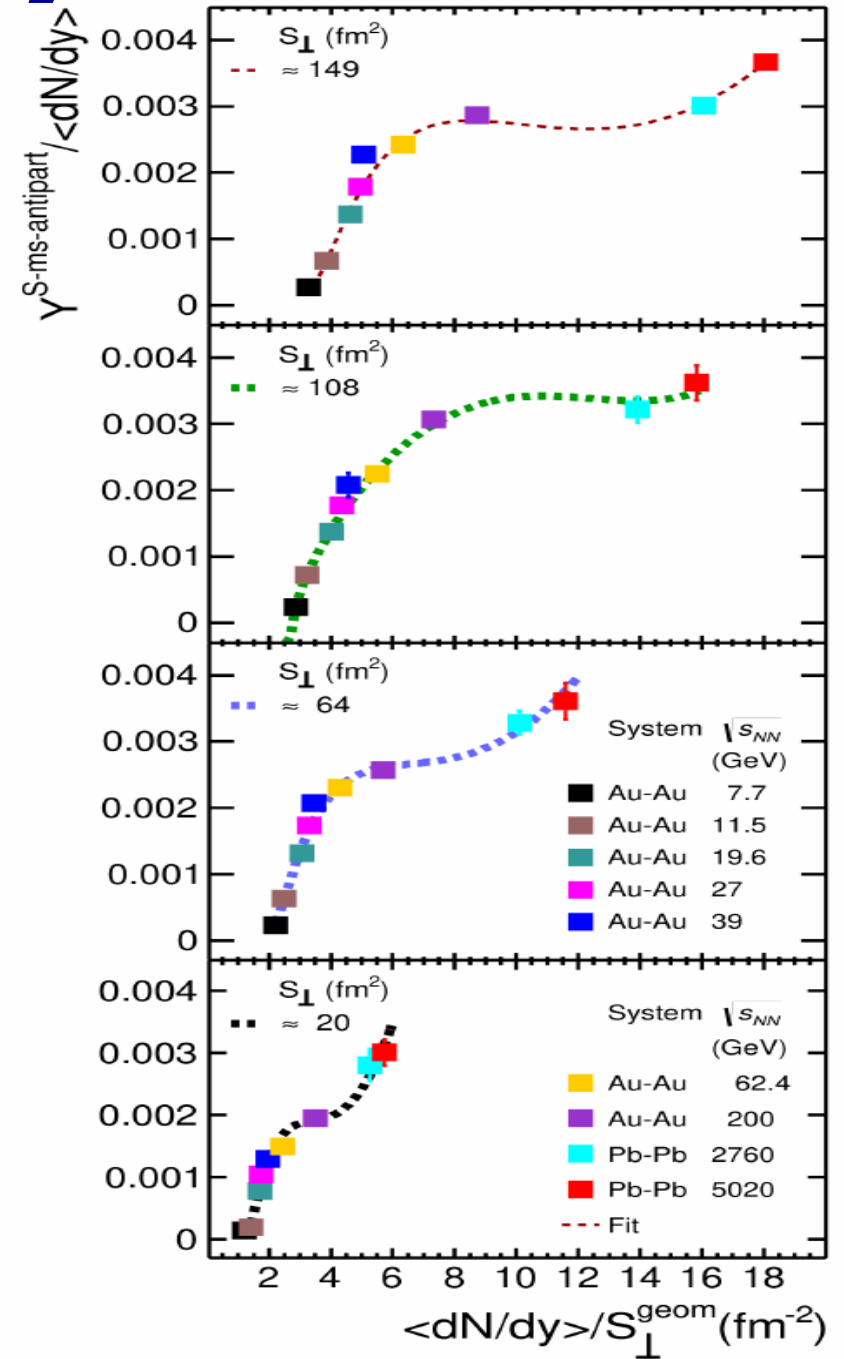
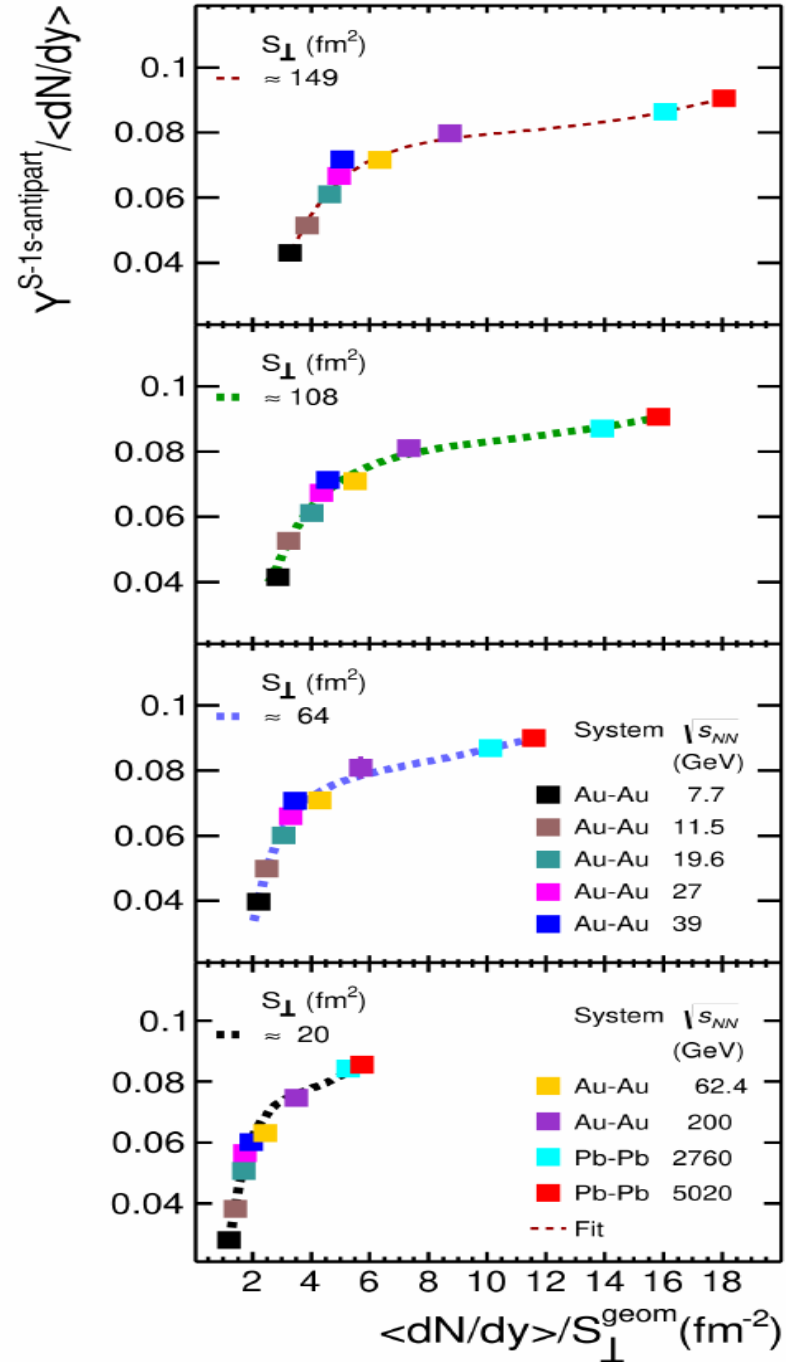
# $(dN/dy)_{(strange\ and\ multi\ strange)} / (dN/dy) - (dN/dy) / S_{\perp}$ correlation

## central collisions

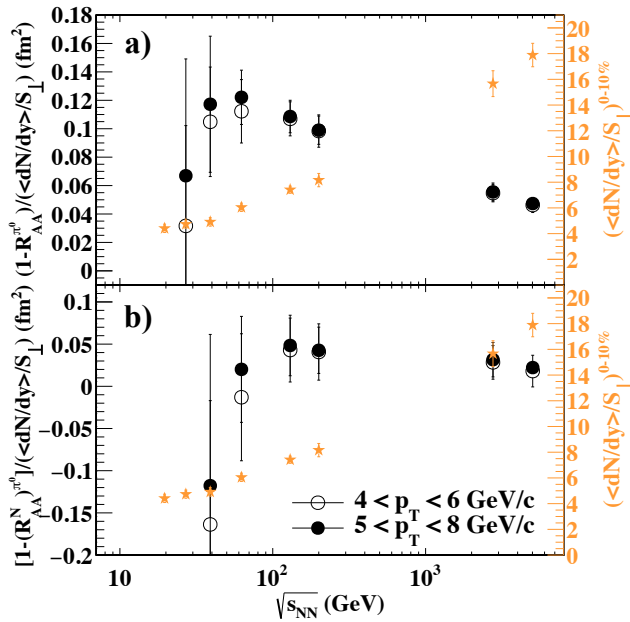


# $(dN/dy)_{\text{strange and multi strange antihadron}} / (dN/dy) - (dN/dy) / S_{\perp}$ correlation

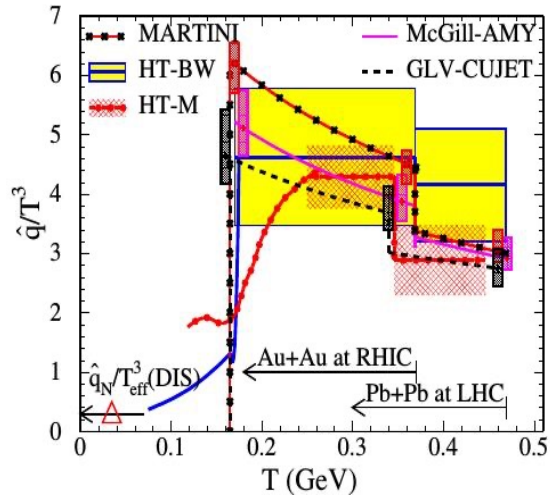
(different  $S_{\perp}$ )



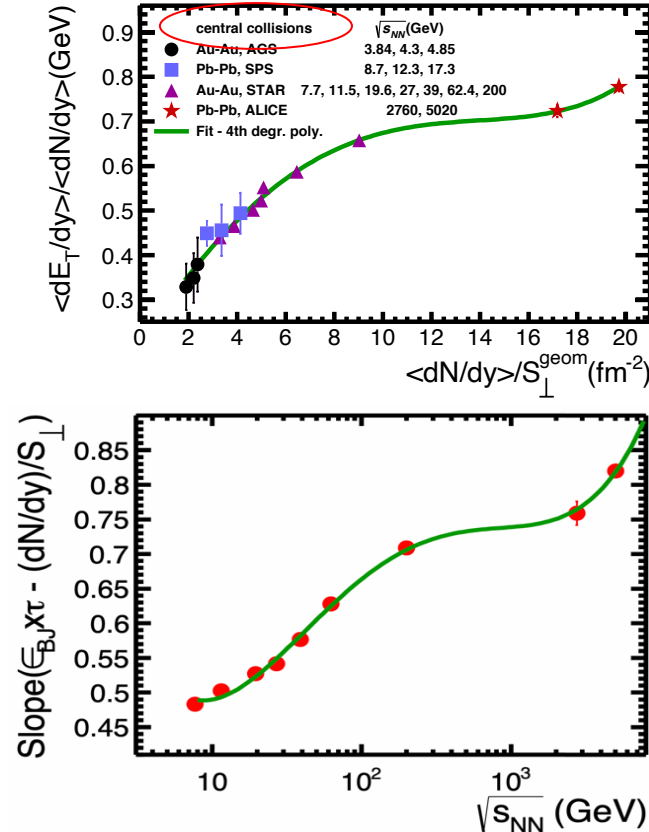
# Do we see a new state of deconfined matter at LHC energies?



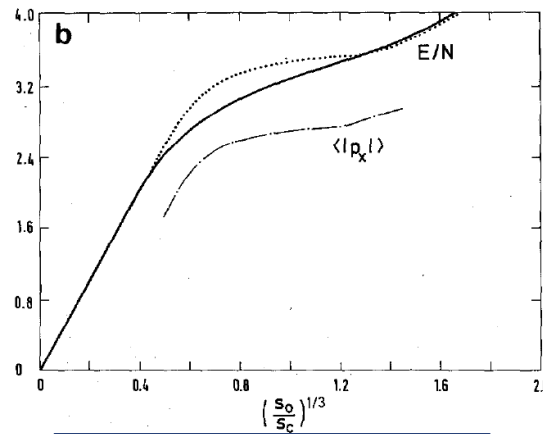
M.Petrovici et al., Phys. Rev. C103(2021)034903



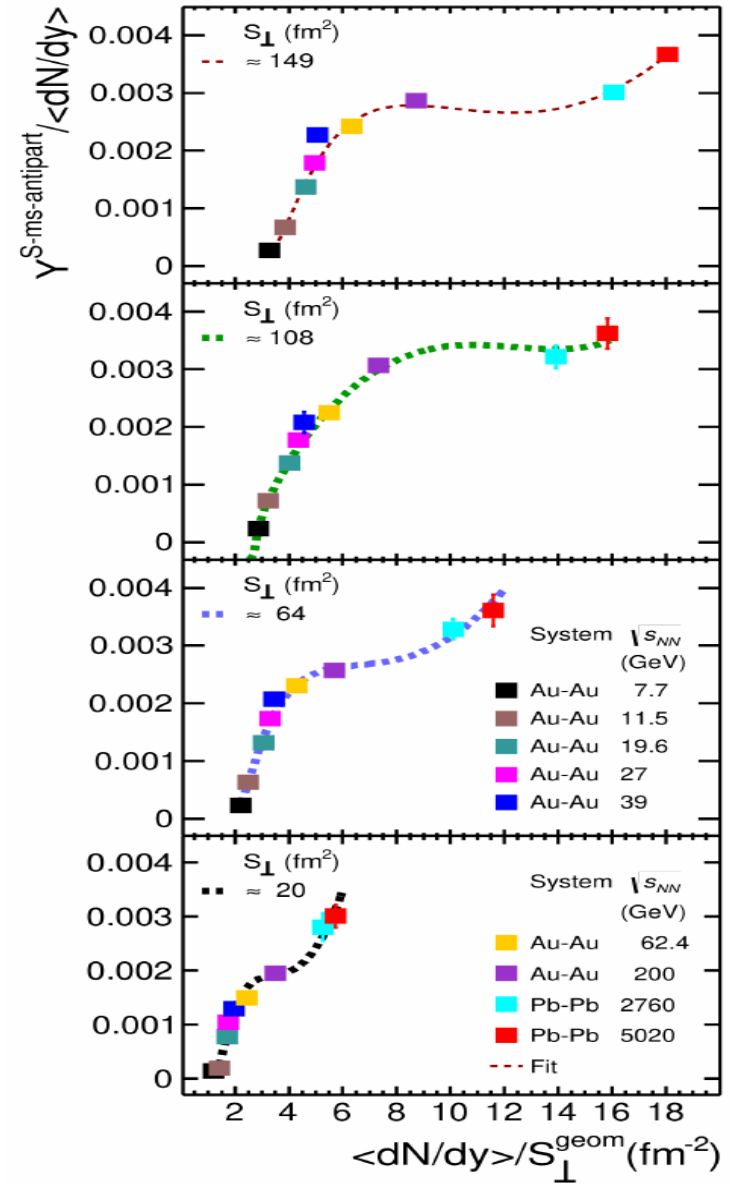
K.M. Burke et al., JET Collaboration, Phys. Rev. C90(2014)014909



M.Petrovici and A.Pop, Phys.Rev. C107(2023)0



J.-P. Blaizot and J.-Y. Ollitrault, Phys.Lett 191B(1987)21



Pop and M. Petrovici, arXiv:2402.19115[hep-ph]

# Short review $pp$ vs $A-A$ @ LHC

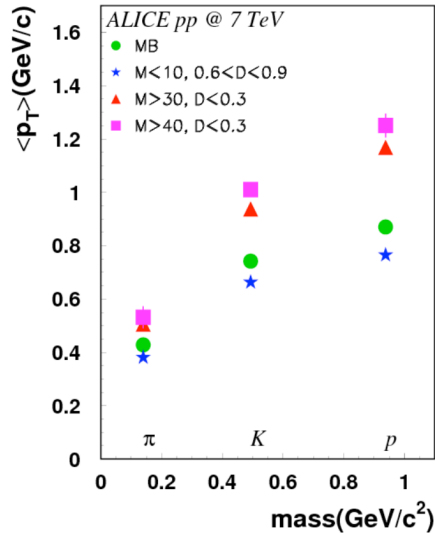


Fig.30

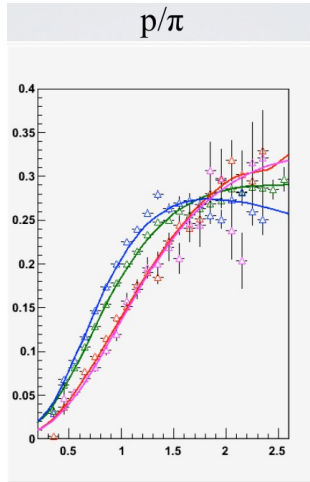


Fig.31

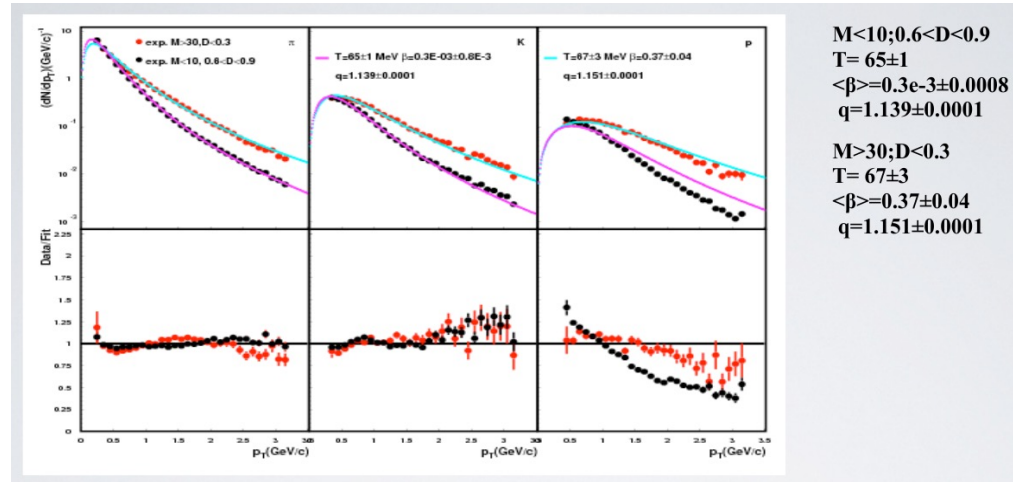


Fig.32

$M < 10; 0.6 < D < 0.9$   
 $T = 65 \pm 1$   
 $\langle \beta \rangle = 0.3e-3 \pm 0.0008$   
 $q = 1.139 \pm 0.0001$

$M > 30; D < 0.3$   
 $T = 67 \pm 3$   
 $\langle \beta \rangle = 0.37 \pm 0.04$   
 $q = 1.151 \pm 0.0001$

Eq. 1

$$D = \frac{|\sum_i p_t^i|}{\sum_i |p_t^i|} \Big|_{\eta > 0},$$

Eq. 2

$$f(p_t) = m_t \int_{-Y}^Y \cosh(y) dy \int_{-\pi}^{\pi} d\phi \int_0^R r dr \left( 1 + \frac{q-1}{T} (m_t \cosh(y) \cosh(\rho) - p_t \sinh(\rho) \cos(\phi)) \right)^{-1/(q-1)}$$

C.Andrei et al., ALICE Week, PWG2-Soft Physics, 9.11.2010  
 C.Andrei et al., Paper draft, 14.03.2011

# Short review pp vs A-A @ LHC

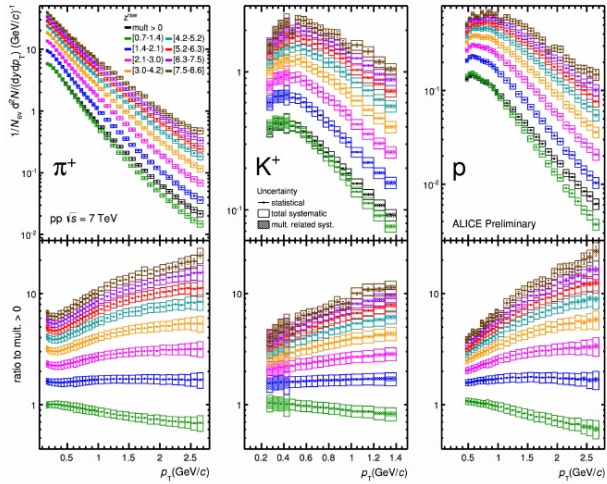


Fig.33

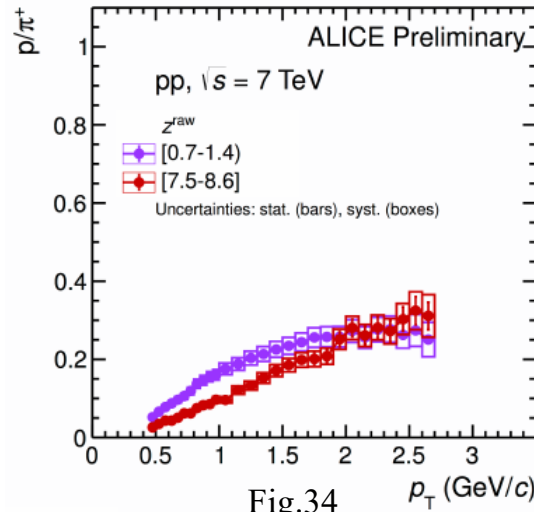


Fig.34

$$z^{raw} = \frac{(N_{ch}^{raw})_{limit}}{\langle N_{ch}^{raw} \rangle_{mult>0}}$$

$$\langle N_{ch}^{raw} \rangle_{mult>0} = 9.6, |\eta| < 0.8$$

$N_{ch}^{raw}$	$z^{raw}$
7 - 12	0.7 - 1.3
13 - 19	1.4 - 2.0
20 - 28	2.1 - 2.9
29 - 39	3.0 - 4.1
40 - 49	4.2 - 5.1
50 - 59	5.2 - 6.2
60 - 71	6.3 - 7.4
72 - 82	7.5 - 8.6

Eq. 5 
$$E \frac{d^3N}{dp^3} \sim f(p_T) = \int_0^R m_T K_1(m_T \cosh \rho / T_{kin}) I_0(p_T \sinh \rho / T_{kin}) r dr$$

were

$$m_T = \sqrt{m^2 + p_T^2}; \beta_r(r) = \beta_s \left(\frac{r}{R}\right)^n$$

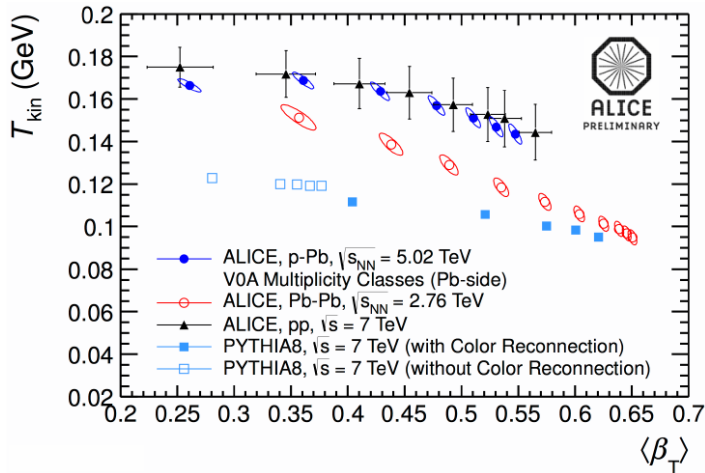


Fig.35

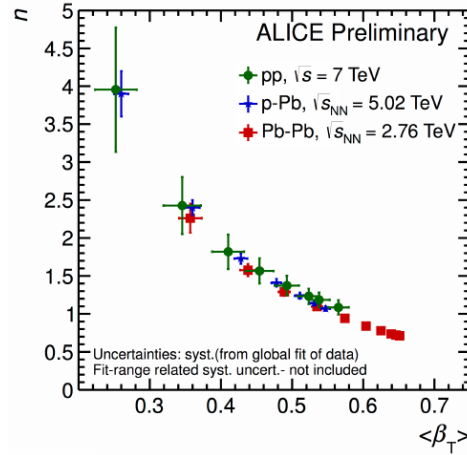
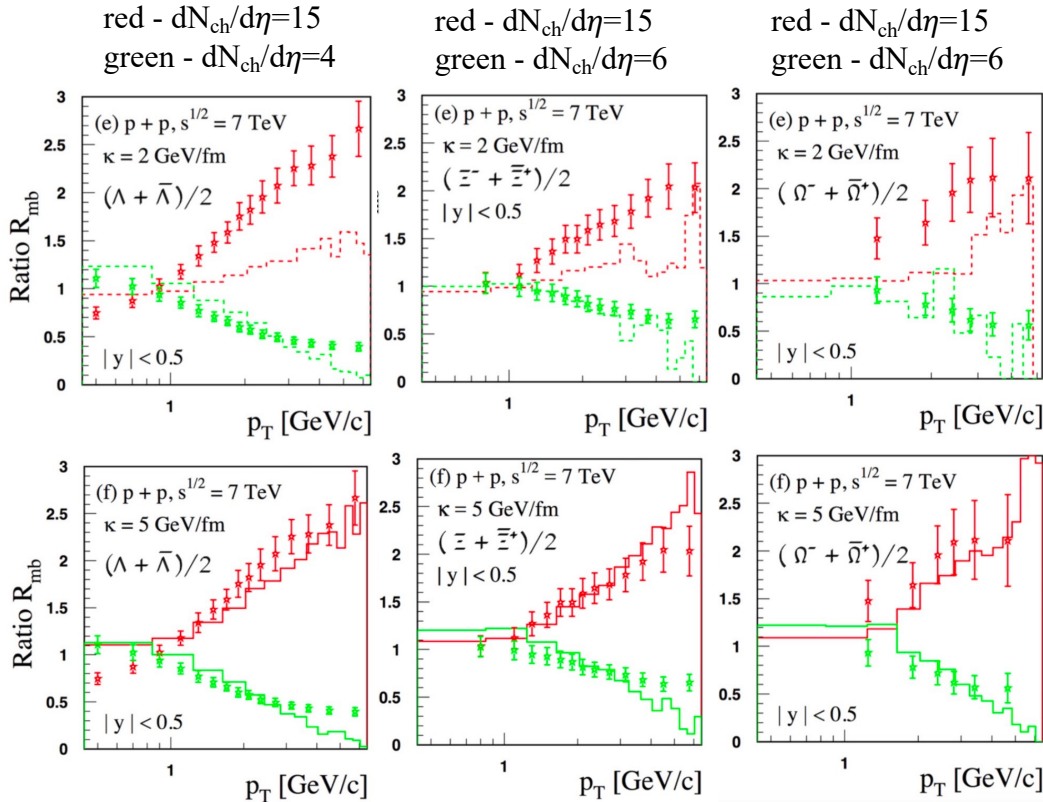
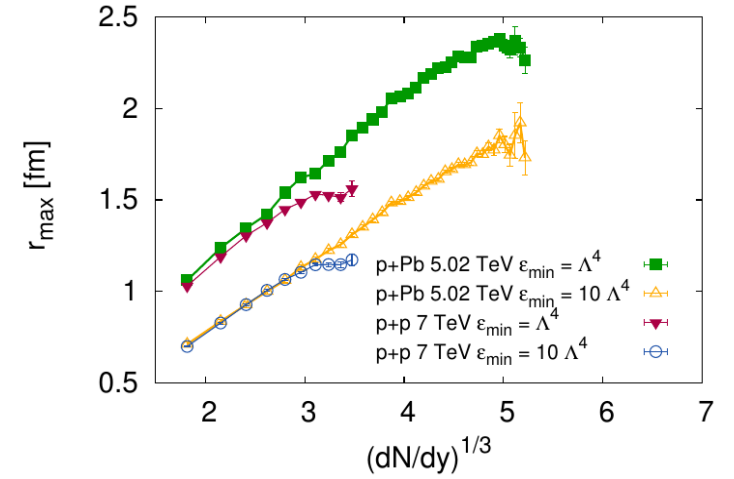


Fig.36

# Short review $pp$ vs $A-A$ @ LHC



*V. Topor Pop and M. Petrovici, Phys. Rev. C 98, 064903 (2018).*



$R_{pp} = l_{fm} \cdot f_{pp}$  - maximal radius for which the energy density of the Yang-Mill fields is larger than  $\varepsilon = \alpha \Lambda_{QCD}^4$  ( $\alpha \in [1, 10]$ )

$$S_{\perp}^{pp} = \pi R_{pp}^2$$

$$\alpha=1 \quad f_{pp} = \begin{cases} 0.387 + 0.0335x + 0.274x^2 - 0.0542x^3 & \text{if } x < 3.4 \\ 1.538 & \text{if } x \geq 3.4 \end{cases}$$

$$x = (dN_g/dy)^{1/3}$$

$$dN_g/dy \approx dN/dy$$

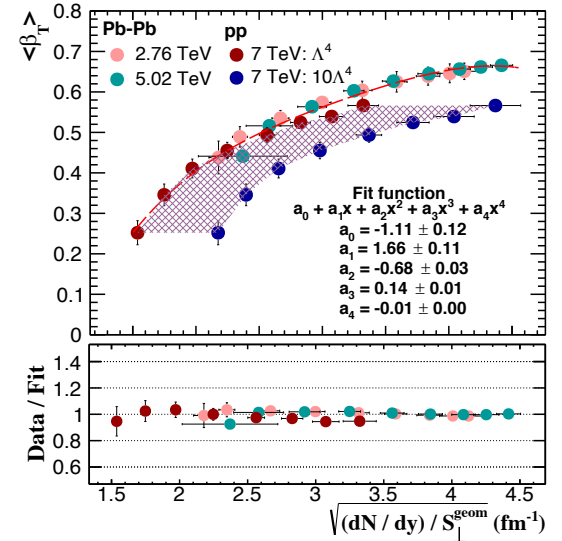
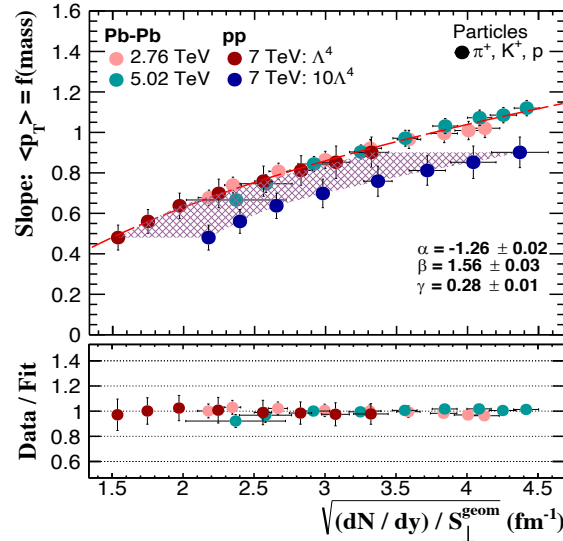
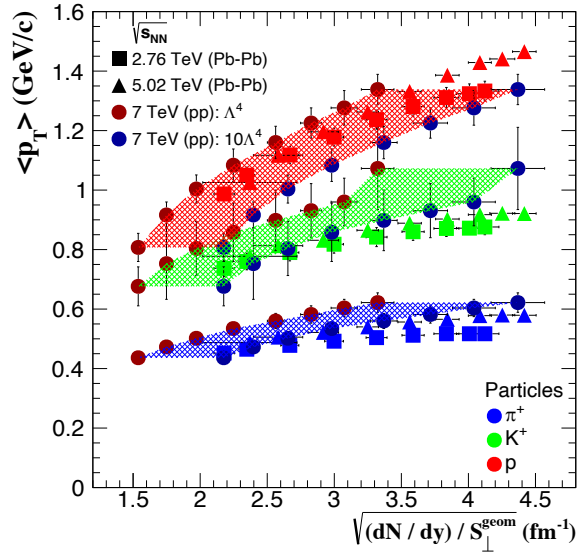
*A. Bzdak et al., Phys.Rev. C87(2013)064906*

*McLarren, M. Praszalowicz and B. Schenke, Phys.Rev. C87(2013)064906*



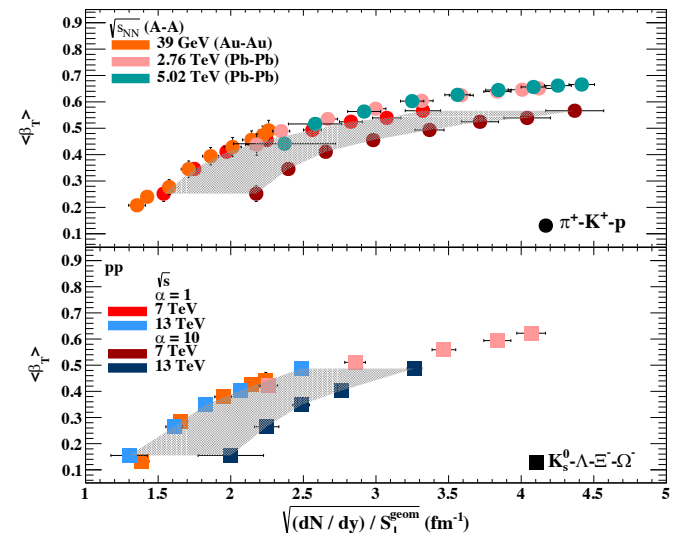
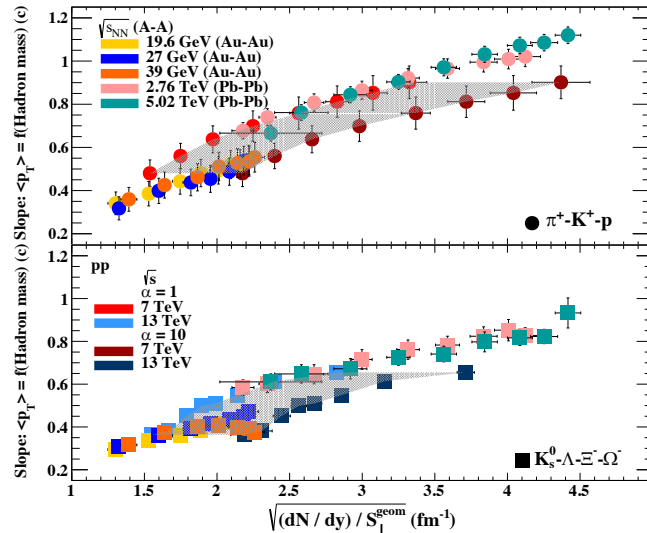
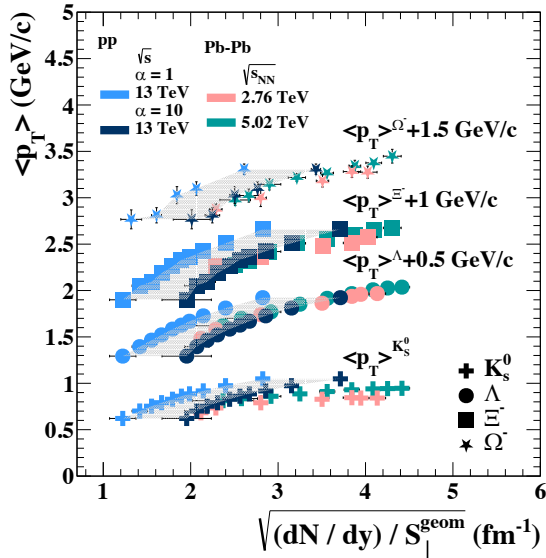
# A-A vs. pp @ LHC

$\pi, K, p$



M.Petrovici, A.Lindner and A.Pop, Phys. Rev. C 98(2018)024904

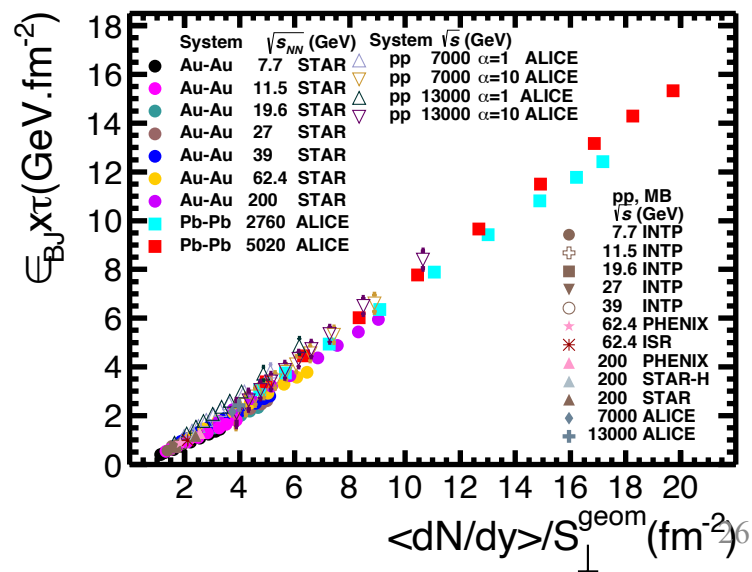
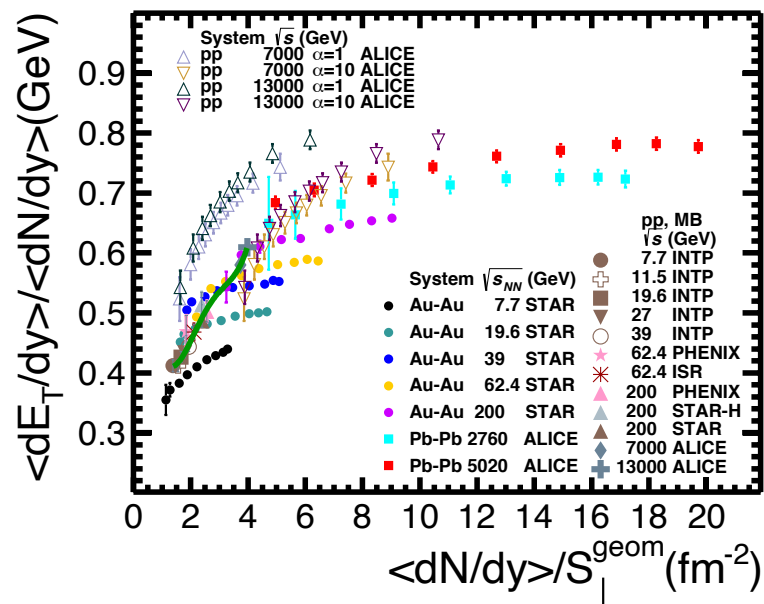
$K^0, \Lambda, \Xi, \Omega$



A. Lindner et al., Proceedings of Science (PoS) 380(2021)197 (PANIC2021), <https://pos.sissa.it/380/197/>.

# A-A vs. pp @ LHC

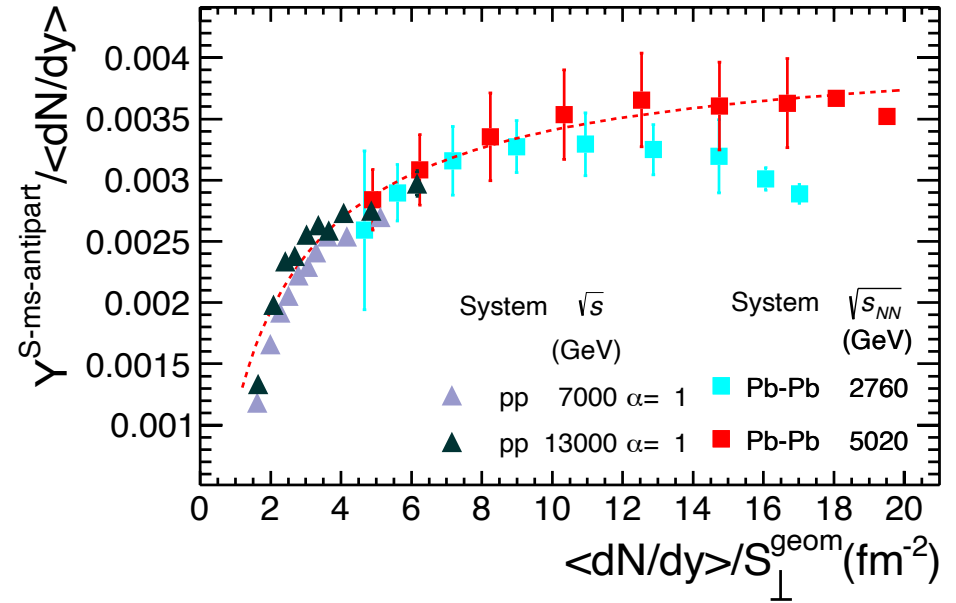
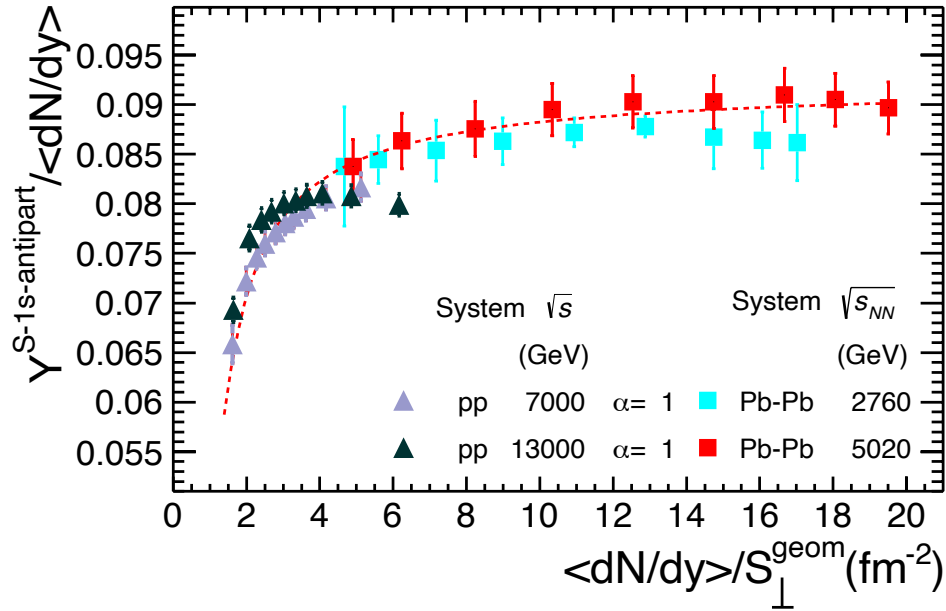
$$(dE_T/dy)/(dN/dy) - (dN/dy)/S_{\perp} \text{ and } \epsilon_{Bj} - (dN/dy)/S_{\perp}$$



M. Petrovici and A. Pop, Phys.Rev. C107(2023)034913

# *A-A vs. pp @ LHC*

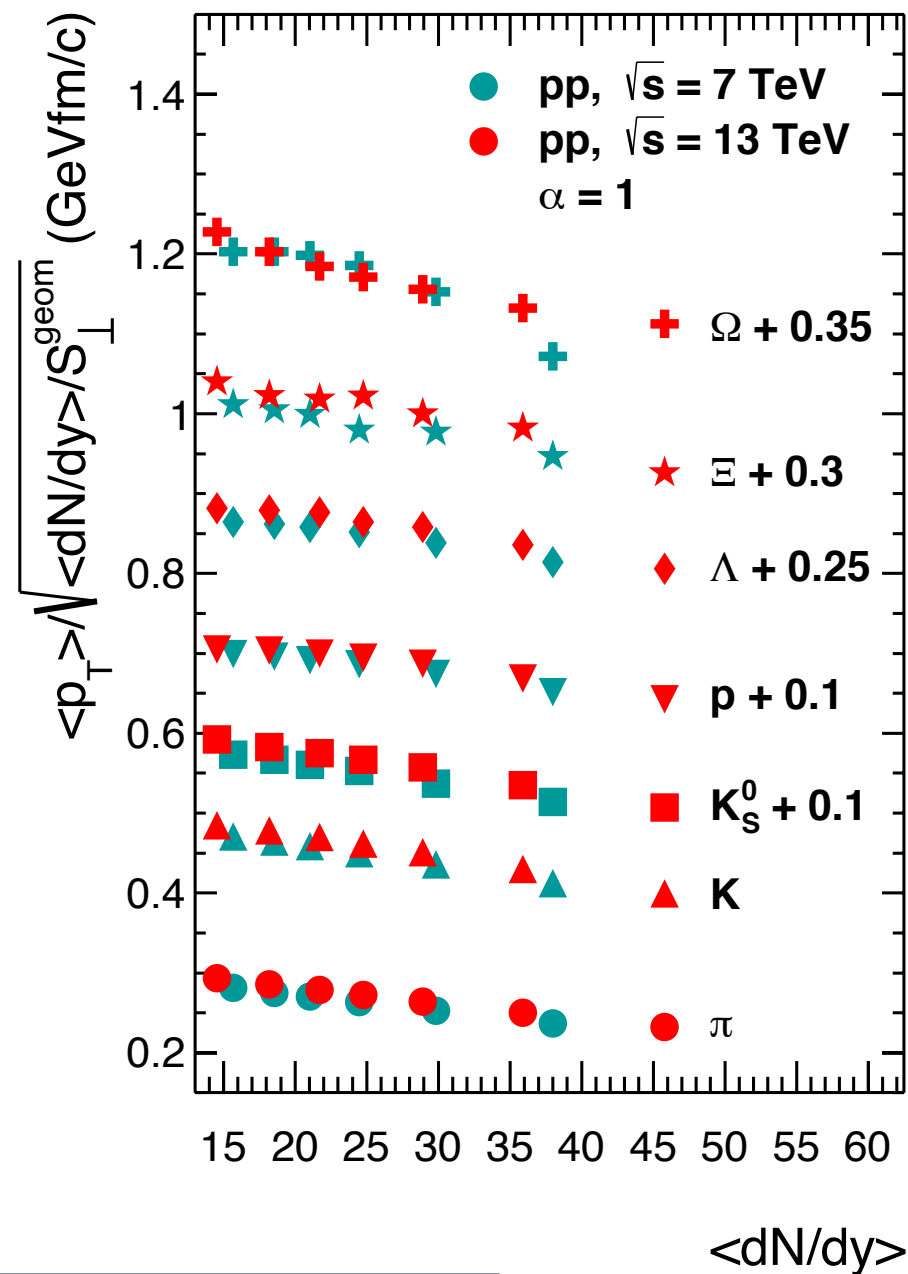
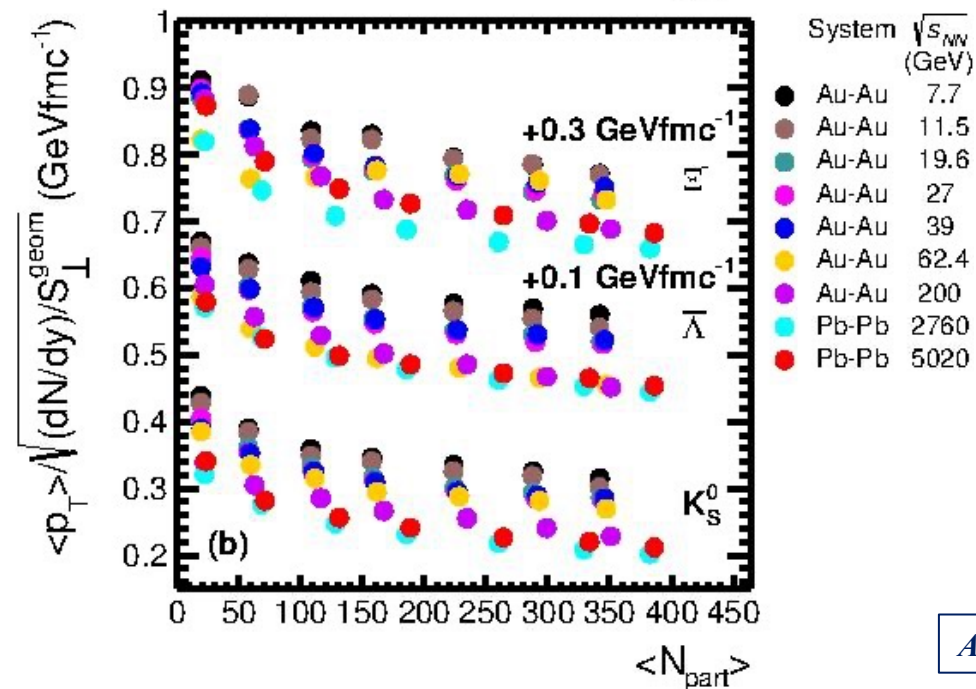
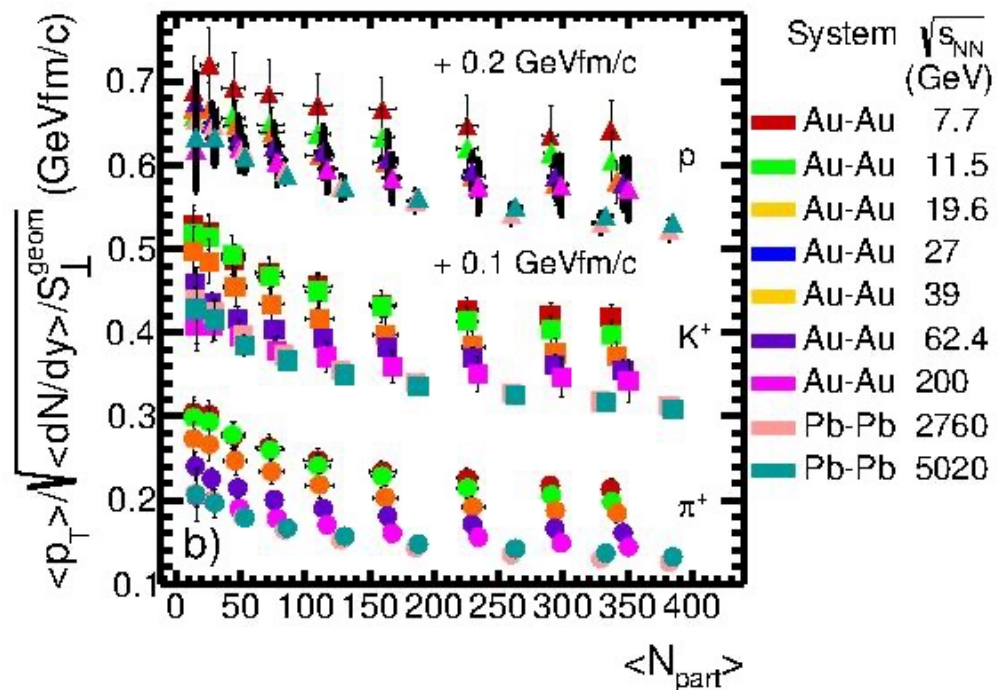
$$(dN/dy)^{\text{(strange and multi strange)}} / (dN/dy) - (dN/dy) / S_{\perp}$$



*M. Petrovici and A. Pop, EuNPC 2022*  
*A. Pop and M. Petrovici, arXiv:2402.19115[hep-ph]*

**Highest charged particle multiplicity in pp at midrapidity selected by "V0M" by ALICE Collaboration !!!**

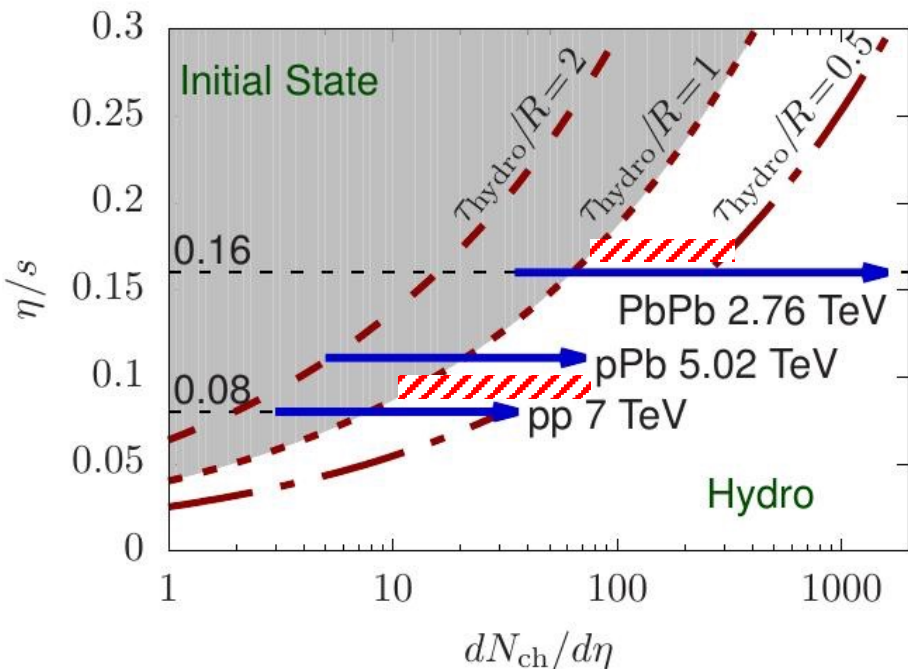
$$\langle p_T \rangle / [ (dN/dy) / S_{\perp}^{geom} ]$$



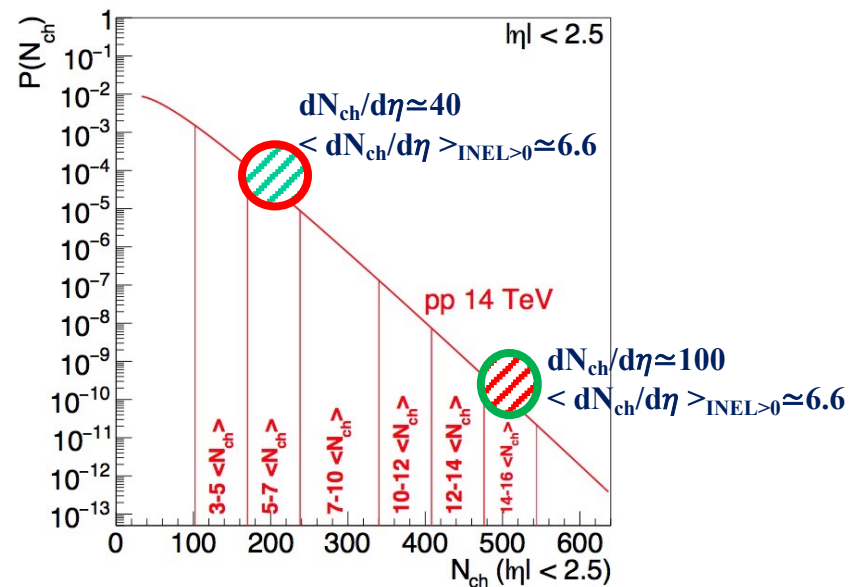
A. Pop and M. Petrovici, will be published

# What's next ?

A.Kurkela et al., PoS(Confinement 2018)152

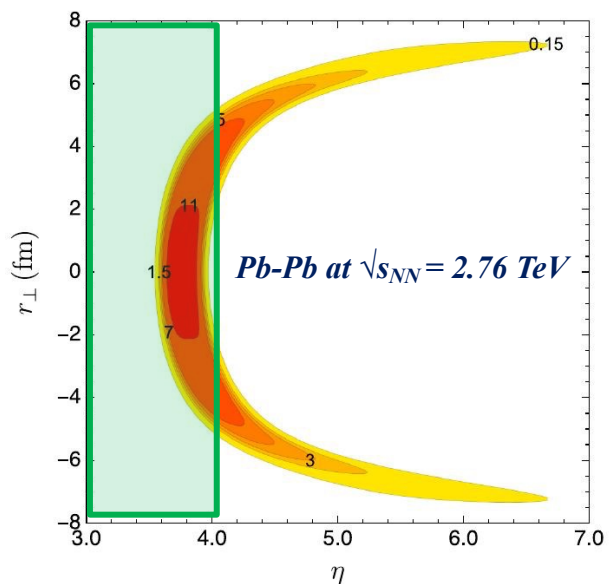
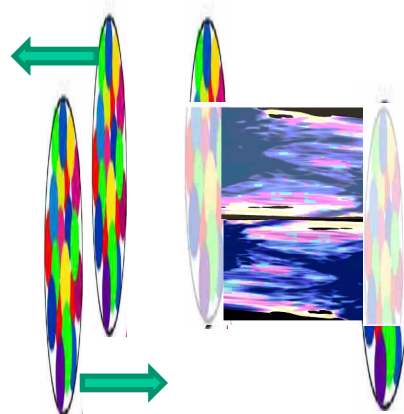
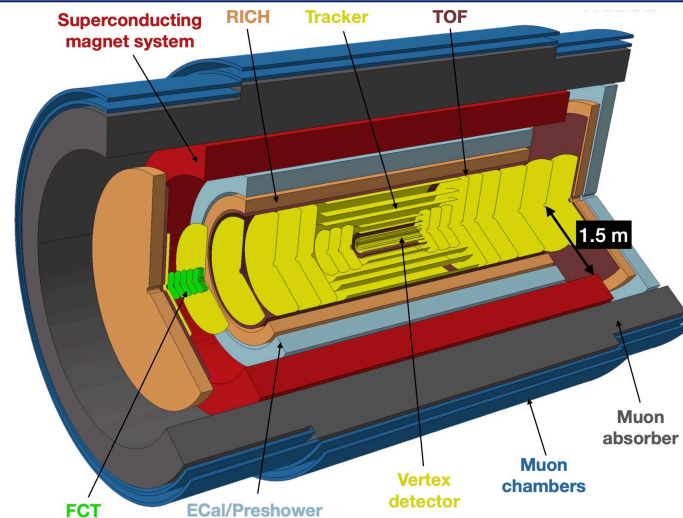


ALICE Coll., arXiv:1812.06772



## ALICE3

ALICE Collaboration, arXiv:2211.02491v1 [physics.ins-det] 4 Nov 2022



M. Li and J.I. Kapusta, Phys.Rev. C99(2019)014906

## *Concluding remark*



*“We have found it of paramount importance that in order to progress we must recognize the ignorance and leave room for doubt. Scientific knowledge is a body of statements of varying degrees of certainty some most unsure, some nearly sure, none absolutely certain.”*

*Richard Feynman*

*Backup slides*

# Expectations based on QCD

## QCD – non-Abelian gauge theory & asymptotic freedom

D.J.Gross, H.D.Politzer and F.Wilczek - Nobel Prize 2004

QCD - running coupling constant 
$$\alpha_s(Q^2) = \frac{\alpha_s(\mu^2)}{1 + \frac{\alpha_s(\mu^2)}{12\pi}(33 - 2n_f)\log(Q^2/\mu^2)}$$

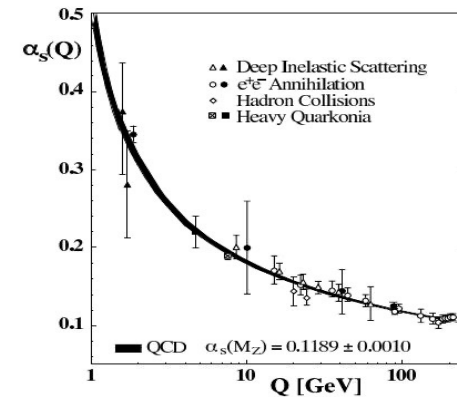
QCD – intrinsic scale 
$$\Lambda^2 = \mu^2 \exp\left[\frac{12\pi}{(33 - 2n_f)\alpha_s(\mu^2)}\right]$$

$$\Rightarrow \alpha_s(Q^2) = \frac{12\pi}{(33 - 2n_f)\log(Q^2/\Lambda^2)}$$

for  $Q^2 \gg \Lambda^2$   $\alpha_s$  is small  $\Rightarrow$  a perturbative description in terms of Quarks and Gluons interacting weakly

$\Rightarrow$

for  $Q^2 \sim \Lambda^2$  Quarks and Gluons arrange themselves in Strongly Bound Clusters - Hadrons

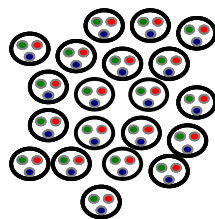


Since  $\Lambda_{QCD} \sim 200$  MeV a phase transition is expected at:

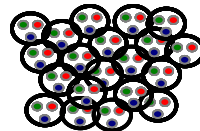
$$T \sim \Lambda_{QCD} \sim O(10^{12} \text{ K})$$

or

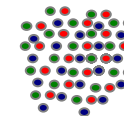
$$\rho_B \sim \Lambda_{QCD}^3 \sim 1 \text{ fm}^{-3}$$



Strongly Bound Clusters  
Hadrons

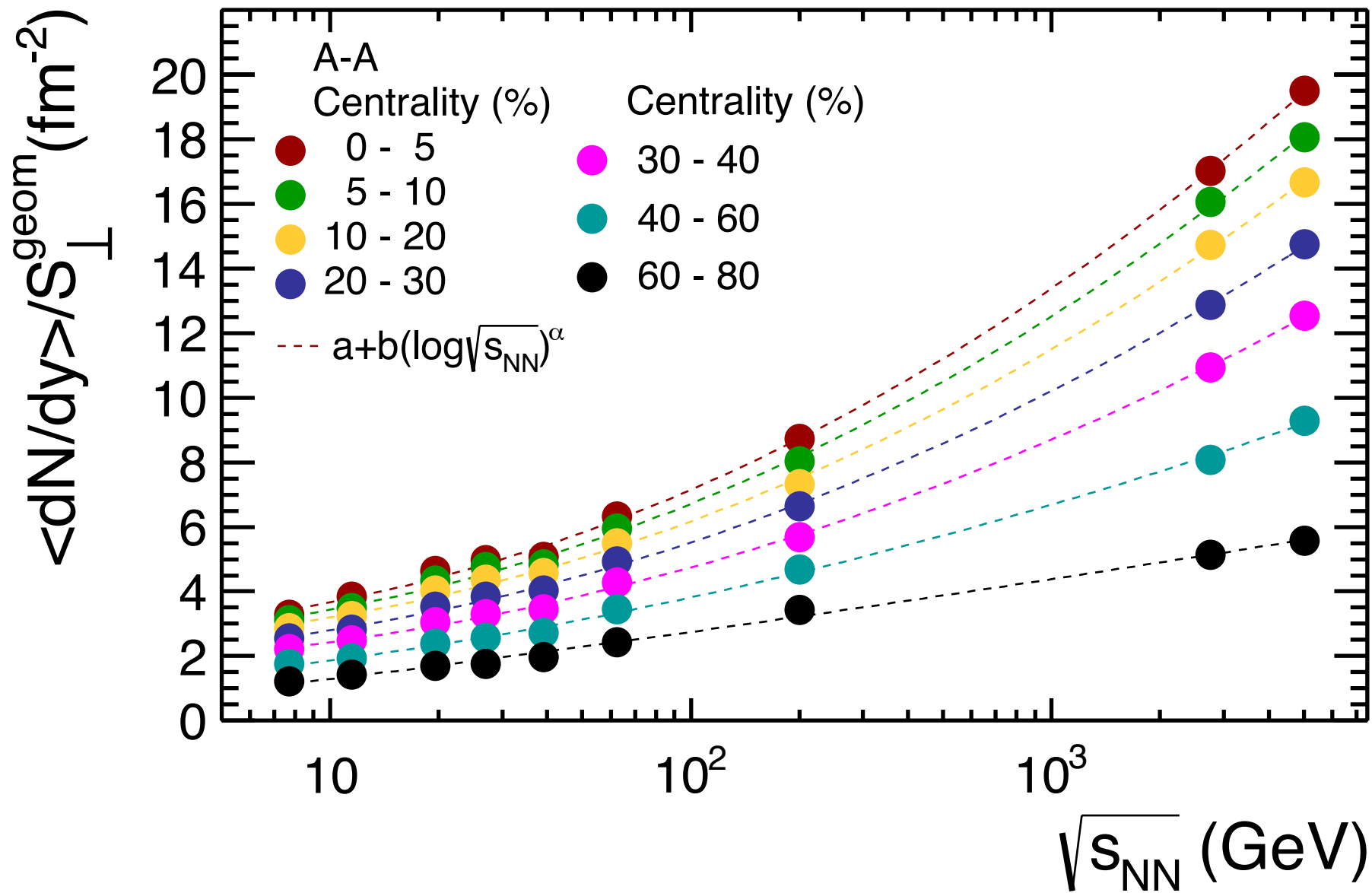


Phase transition



Weakly interacting  
Quarks and Gluons

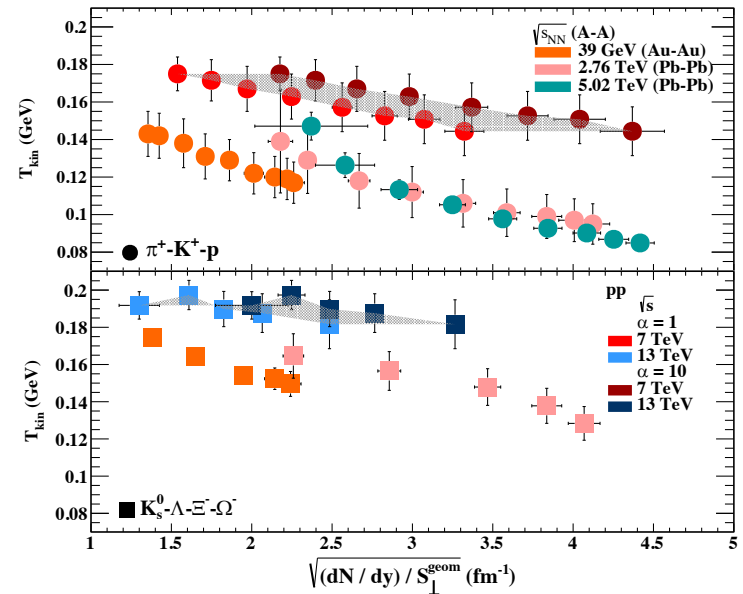
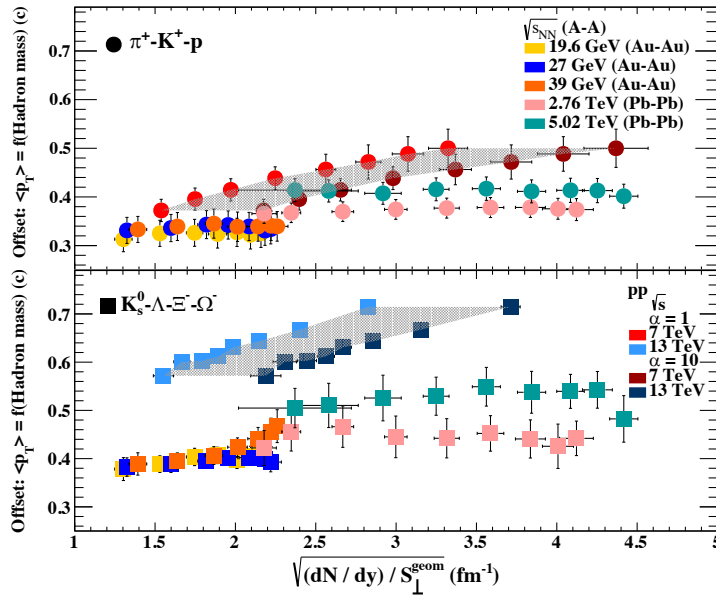




# *pp vs. Pb-Pb @ LHC - $(dN/dy)/S_{\perp}$ scaling*

Observable	$\alpha$	species
$\langle p_T \rangle = f([(dN/dy)/S_{\perp}]^{1/2})$	10	$\pi, K^-, K_s^0, \Lambda, \Xi, \Omega$
	1 (low mult. $\rightarrow$ 10 (high mult.))	p
$\langle dE_T/dy \rangle / \langle dN/dy \rangle = f([(dN/dy)/S_{\perp}]^{1/2})$	10	$\pi, K^-, K_s^0, p, \Lambda, \Xi, \Omega$
Slope $p_T = f(\text{mass})$	1	$\pi, K^-, K_s^0, p, \Lambda, \Xi, \Omega$
$\langle \beta_T \rangle$	1	$\pi, K^-, K_s^0, p, \Lambda, \Xi, \Omega$
$Y_{1s(\text{ms})} / \langle dN/dy \rangle$	1	$K, \Lambda, \Xi, \Omega$

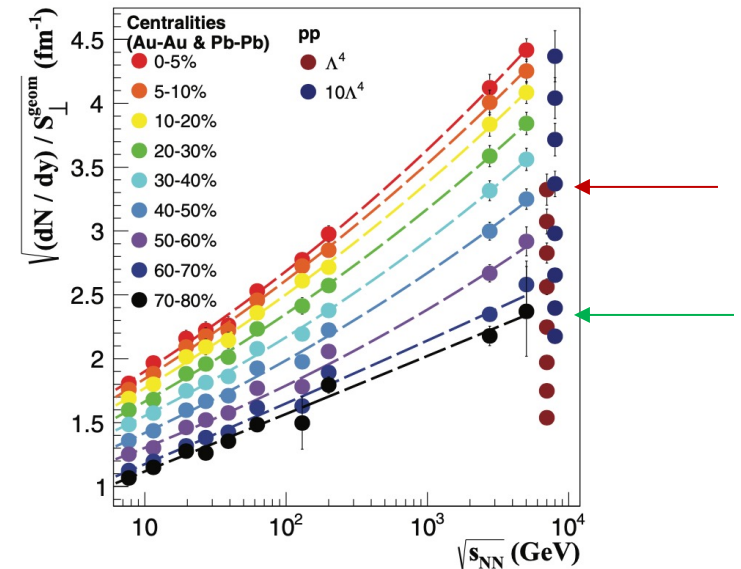
*Why the offset of  $p_T = f(\text{mass})$  and  $T_{kin}^0$  from BGBW fits do not scale ?*



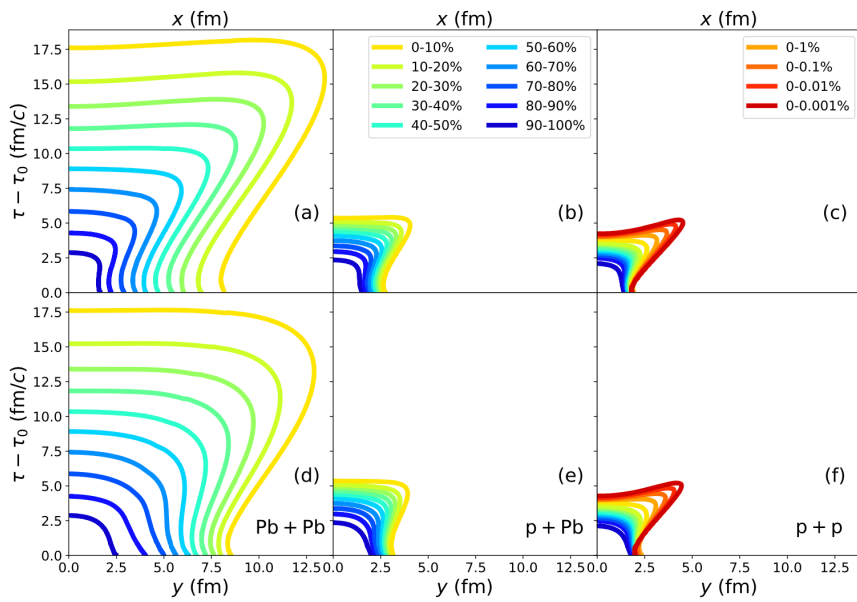
# *pp vs. Pb-Pb @ LHC*

System	$\sqrt{s_{NN}}$ (GeV)	Cen. (%)	$\langle N_{part} \rangle$	$S_{\perp}^{geom}$ (fm <sup>2</sup> )	$S_{\perp}^{var}$ (fm <sup>2</sup> )	$f_{core}$	$(S_{\perp}^{geom})_{core}$ (fm <sup>2</sup> )	$(S_{\perp}^{var})_{core}$ (fm <sup>2</sup> )	$dN/dy$
Pb-Pb	2760	0–5	382.5 ± 3.1	166.9 ± 0.7	170.7 ± 0.7	0.94 ± 0.00	146.0 ± 0.7	148.0 ± 0.6	2837.0 ± 144.0
		5–10	329.4 ± 4.9	146.1 ± 0.7	154.7 ± 0.6	0.90 ± 0.00	121.9 ± 0.7	126.5 ± 0.5	2345.5 ± 112.4
		10–20	259.9 ± 2.9	119.8 ± 0.8	132.4 ± 0.6	0.86 ± 0.00	96.3 ± 0.7	102.7 ± 0.4	1763.2 ± 84.8
		20–30	185.4 ± 3.9	92.9 ± 0.8	107.5 ± 0.5	0.81 ± 0.00	71.5 ± 0.8	78.4 ± 0.3	1195.8 ± 54.2
		30–40	128.1 ± 3.3	71.4 ± 0.8	87.2 ± 0.4	0.76 ± 0.00	52.4 ± 0.8	59.7 ± 0.2	784.8 ± 35.9
		40–50	84.2 ± 2.6	53.7 ± 0.8	70.3 ± 0.3	0.70 ± 0.00	37.2 ± 0.8	44.8 ± 0.2	482.7 ± 21.4
		50–60	52.1 ± 2.0	38.6 ± 0.8	56.1 ± 0.3	0.63 ± 0.00	24.7 ± 0.9	33.1 ± 0.1	274.8 ± 12.5
		60–70	29.5 ± 1.3	25.7 ± 0.8	43.6 ± 0.2	0.54 ± 0.00	14.6 ± 0.9	23.8 ± 0.1	141.8 ± 5.4
		70–80	14.9 ± 0.6	14.2 ± 0.8	30.8 ± 0.2	0.43 ± 0.00	6.4 ± 0.7	15.1 ± 0.1	67.2 ± 3.0
Pb-Pb	5020	0–5	385 ± 2	170.2 ± 0.7	174.2 ± 0.7	0.94 ± 0.00	149.0 ± 0.7	151.5 ± 0.6	3320.6 ± 131.4
		5–10	333 ± 4	149.2 ± 0.7	158.5 ± 0.6	0.90 ± 0.00	124.4 ± 0.7	129.9 ± 0.5	2698.7 ± 117.2
		10–20	263 ± 4	122.4 ± 0.8	135.8 ± 0.6	0.86 ± 0.00	98.1 ± 0.7	105.6 ± 0.4	2042.5 ± 84.7
		20–30	188 ± 3	94.9 ± 0.8	110.5 ± 0.5	0.82 ± 0.00	72.9 ± 0.7	80.8 ± 0.3	1401.4 ± 62.9
		30–40	131 ± 2	73.4 ± 0.8	90.0 ± 0.4	0.77 ± 0.00	53.8 ± 0.8	61.8 ± 0.3	931.0 ± 44.5
		→ 40–50	86.3 ± 1.7	55.7 ± 0.8	73.1 ± 0.3	0.71 ± 0.00	38.6 ± 0.8	46.9 ± 0.2	588.6 ± 27.8
		50–60	53.6 ± 1.2	40.7 ± 0.8	58.7 ± 0.3	0.63 ± 0.00	26.3 ± 0.8	34.9 ± 0.2	346.9 ± 26.1
		60–70	30.0 ± 0.8	27.9 ± 0.8	45.9 ± 0.2	0.54 ± 0.01	16.2 ± 0.8	25.5 ± 0.1	186.1 ± 26.0
		→ 70–80	15.6 ± 0.5	16.6 ± 0.7	33.0 ± 0.2	0.43 ± 0.01	7.7 ± 0.7	17.0 ± 0.1	93.5 ± 27.4

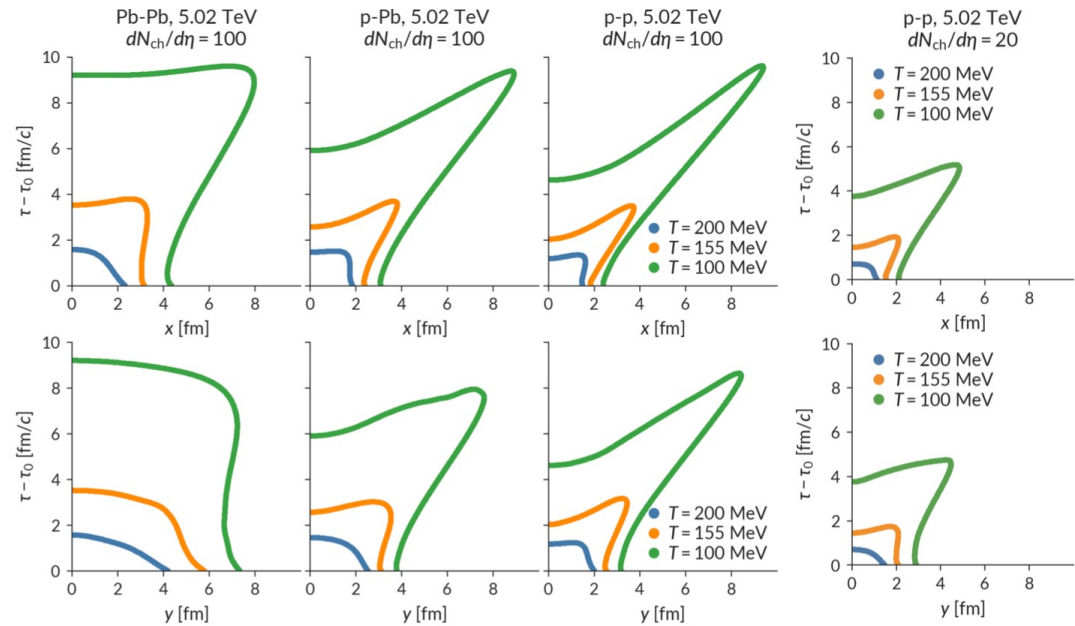
$\sqrt{s}$ (TeV)	$dN/dy$	$S_{\perp}$ (fm <sup>2</sup> )	
		$\alpha = 1$	$\alpha = 10$
(pp)	→ 82.1 ± 2.8	7.43 ± 0.48	4.30 ± 0.36
	70.2 ± 2.2	7.43 ± 0.41	4.30 ± 0.31
	59.4 ± 1.7	7.43 ± 0.35	4.30 ± 0.27
	48.8 ± 1.3	7.43 ± 0.30	4.30 ± 0.23
	→ 37.3 ± 0.9	7.39 ± 0.02	4.20 ± 0.02
	26.8 ± 0.6	6.89 ± 0.05	3.80 ± 0.03
	18.2 ± 0.4	5.94 ± 0.06	3.16 ± 0.04
	10.8 ± 0.2	4.58 ± 0.06	2.29 ± 0.04



# *pp vs. Pb-Pb @ LHC - hydro models*



*C. Plumberg, Phys.Rev. C102(2020)054908*



*U. Heinz et al., Journal of Physics: Conf. Series 1271(2019)012018*