

State-of-the-art input for simulation of nuclear processes at high energy: where heavy-ion and nuclear structure meets

Benjamin Bally

QCD challenges from pp to AA collisions

Münster - 04/09/2024



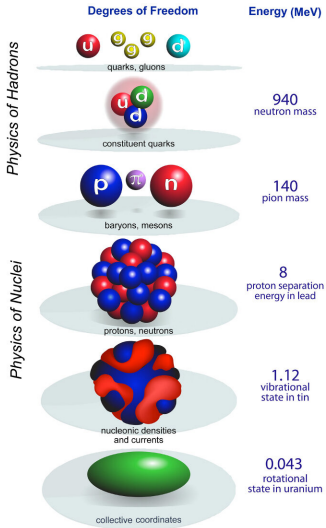
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 - ◇ Exploit versatility of LHCb/SMOG2 Giacalone *et al.*, arXiv:2405.20210 (2024)

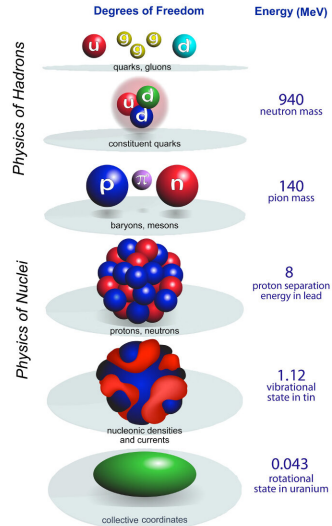
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- Potential new method to probe low-energy observables

- Nuclear matter made of quarks and gluons



Courtesy of ORNL

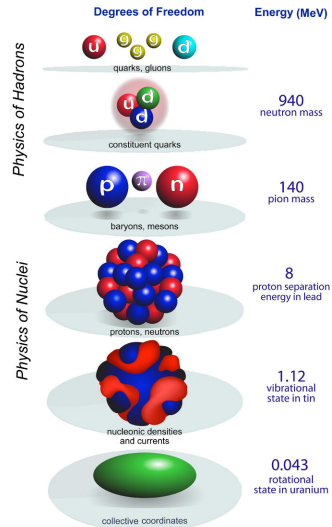
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 - ◊ Possible only for lightest nuclei ($A \leq 4$)
 - ◊ Even if possible, would be very inefficient
 - ◊ What would we learn?



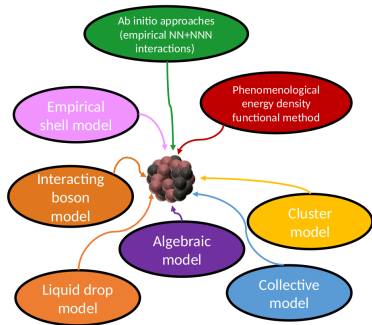
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- Define appropriate degrees of freedom for the scale
- Connect different scales

→ **Tower of Effective Field Theories**



Courtesy of ORNL



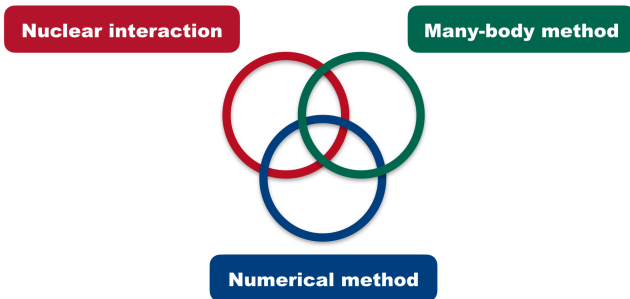
Courtesy of J.-P. Ebran

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Courtesy of P. Arthuis

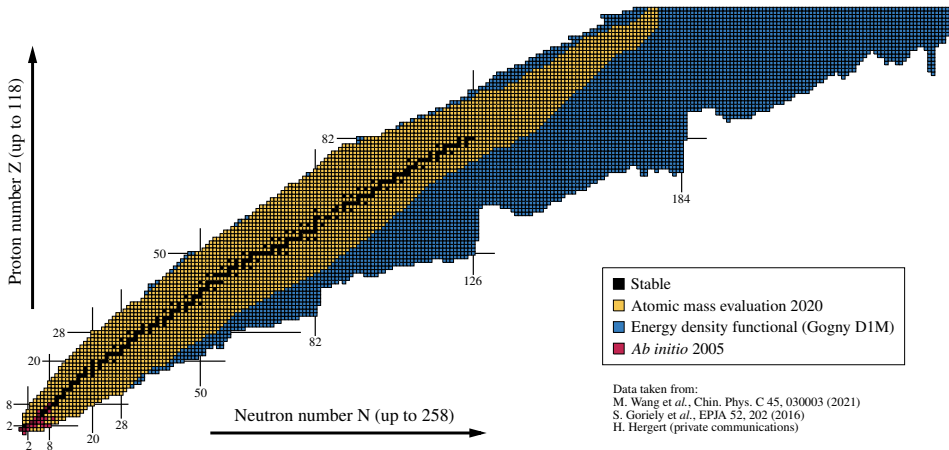
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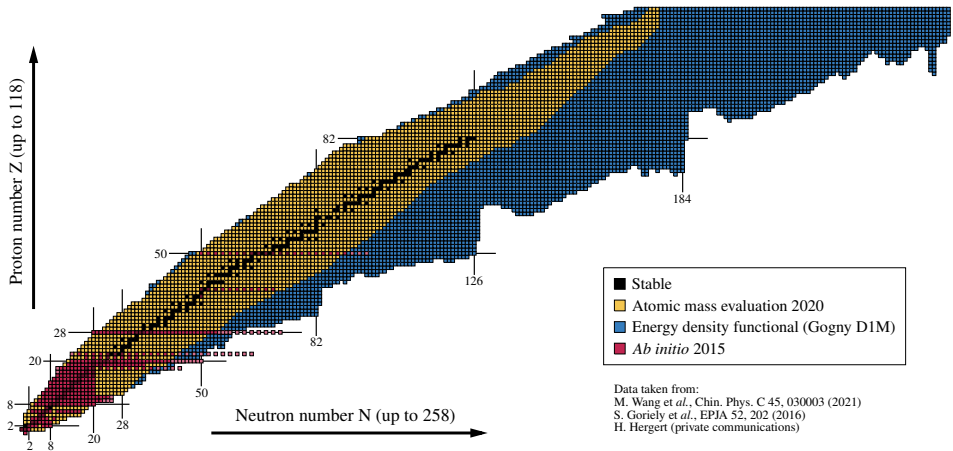
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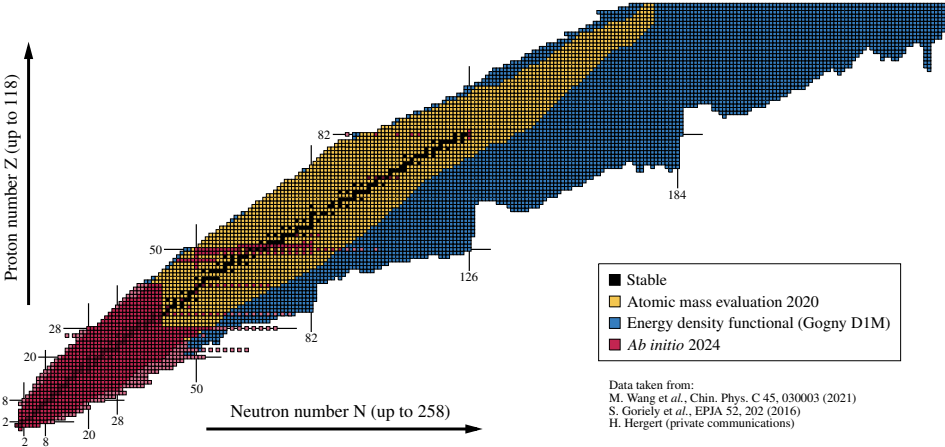
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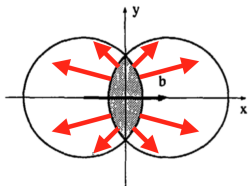
	Two-nucleon force	Three-nucleon force	Four-nucleon force
LO (Q^0)		—	—
NLO (Q^2)		—	—
N ² LO (Q^3)			—
N ³ LO (Q^4)			
N ⁴ LO (Q^5)			

Epelbaum et al., Front. Phys. 8, 98 (2020)



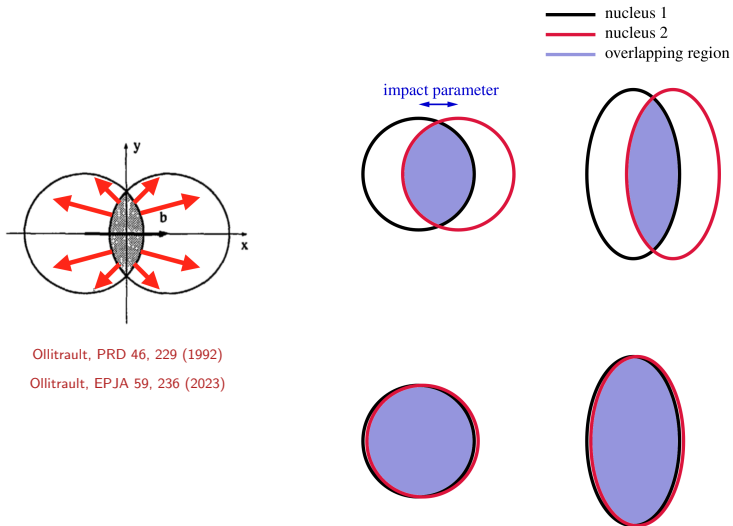




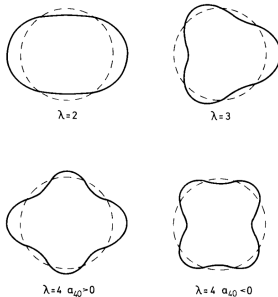


Ollitrault, PRD 46, 229 (1992)

Ollitrault, EPJA 59, 236 (2023)

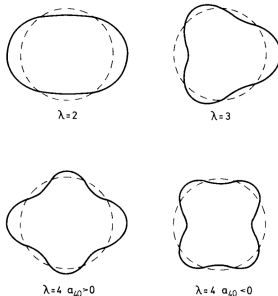


- Nuclear phenomenology often relies on the picture of intrinsic shapes
- Multipole expansion where small values of λ are the most important



Ring and Schuck, *The Nuclear Many-Body Problem* (1980)

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- Multipole expansion where small values of λ are the most important



Ring and Schuck, *The Nuclear Many-Body Problem* (1980)

- Warning: these shapes are not observables!
→ Convenient way to include many-body correlations in a one-body framework

- Woods-Saxon (WS) density profile

$$\rho_{ws}(r, \theta, \phi) = \frac{\rho_0}{1 + \exp\left[\frac{1}{a}\left(r - R(\theta, \phi)\right)\right]}$$

with nuclear radius (**surface** deformation)

$$R(\theta, \phi) = R_0 \left\{ 1 + \sum_{\lambda} \sum_{\mu=-\lambda}^{\lambda} a_{\lambda\mu} Y_{\lambda\mu}(\theta, \phi) \right\}$$

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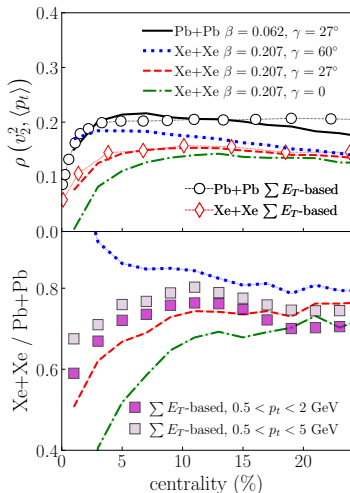
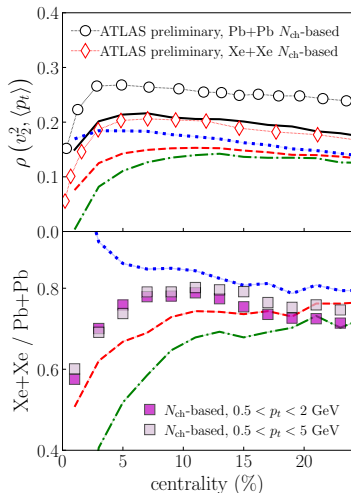
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- In HIC, people usually and incorrectly
 - ◇ assume a spherical WS
 - ◇ fit the WS parameters to the charge density/radius
 - ◇ consider only nucleon density (not proton and neutron densities)
 - ◇ take the WS parameters from NS calculations based on different definitions
→ in general NS use **volume** deformations related to $\langle \Phi | r^{\lambda} Y_{\lambda\mu} | \Phi \rangle$

Ryssens *et al.*, PRL 130, 212302 (2023)

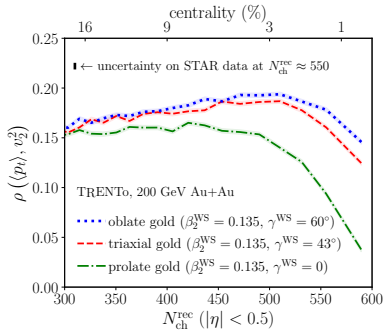
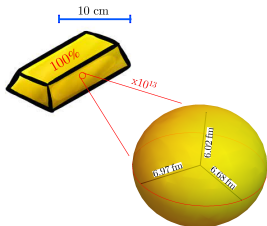
- A better, and still very simple, approach would be
 - ◇ perform NS calculations to determine $\rho_{\text{mic},n}(r, \theta, \phi)$ and $\rho_{\text{mic},p}(r, \theta, \phi)$
 - ◇ fit ρ_{ws} to ρ_{mic} for protons and neutrons separately
 - ◇ sample nucleons using $\rho_{\text{ws},n}$ and $\rho_{\text{ws},p}$

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 - ◇ sample nucleons using $\rho_{\text{ws},n}$ and $\rho_{\text{ws},p}$
- Parameters from *phenomenological* calculations available for
 - ◇ ^{129}Xe → Bally *et al.*, PRL 128, 082301 (2022)
 - ◇ ^{197}Au → Bally *et al.*, EPJA 59, 58 (2023); Ryssens *et al.*, PRL 130, 212302 (2023)
 - ◇ ^{208}Pb → Bally *et al.*, PRL 128, 082301 (2022)
 - ◇ ^{238}U → Ryssens *et al.*, PRL 130, 212302 (2023)



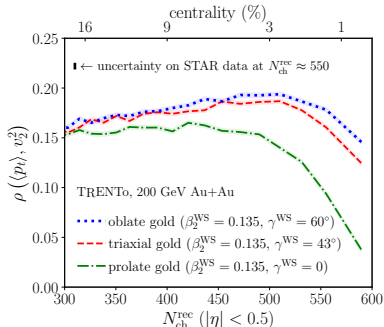
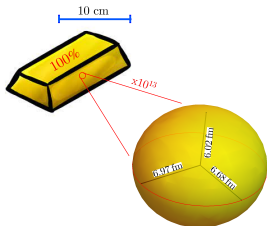
- Only triaxiality explains LHC results

Example: $^{197}\text{Au} + ^{197}\text{Au}$ at RHIC



Bally *et al.*, EPJA 59, 58 (2023)

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Bally *et al.*, EPJA 59, 58 (2023)

- Remark: excellent agreement for neutron skin (although WS not perfect)

$$\Delta r_{np}[\text{STAR}] = 0.17 \pm 0.03 \text{ (stat.)} \pm 0.08 \text{ (syst.) fm} \quad \text{STAR, Science Adv. 9, eabq390 (2023)}$$

$$\Delta r_{np}[\text{MREDF}] = 0.17 \text{ fm}$$

$$\Delta r_{np}[\text{WS fit}] = 0.19 \text{ fm}$$

- Collisions scheduled or already performed at LHC/RHIC
 - ◊ $^{16}\text{O} + ^{16}\text{O}$
 - ◊ $^{208}\text{Pb} + \text{Ne}$

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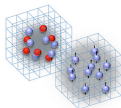
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- New ideas related to
 - ◇ collectivity in small systems
 - ◇ physics program/opportunities at LHCb/SMOG2

EFT \rightarrow Nuclear structure \rightarrow Relativistic hydrodynamic \rightarrow Hadron transport \rightarrow Nucl-th

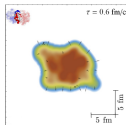
χ EFT
N3LO



PGCM



NLEFT



Trajectum



arXiv:2402.05995

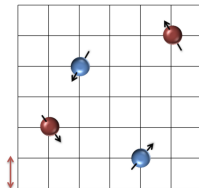
- TAURUS: <https://github.com/project-taurus>
- Trajectum: <https://sites.google.com/view/govertnijs/trajectum>
- SMASH: <https://github.com/smash-transport/smash>

- Nuclear Lattice Effective field Theory (NLEFT)

Lee, *Front. in Phys.* 8, 174 (2020)

Lähde and Meißner, *Lectures Notes in Phys.*, Springer (2019)

- Mesh with 8 sites and spacing $a = 1.315$ fm



- Minimal pionless EFT Hamiltonian with $SU(4)$ symmetry

- Pin-hole algorithm \rightarrow sample nucleon positions from A -body density (include all correlations)

Elhatisari *et al.*, *PRL* 119, 222505 (2017)

- Projected Generator Coordinate Method (PGCM)

Hill and Wheeler, Phys. Rev. 89, 1102 (1953); Griffin and Wheeler, Phys. Rev. 108, 311 (1957)

Bally *et al.*, EPJA 60, 62 (2024)

- Approximate wave function of the form: $|\Psi\rangle = \sum_q f(q)|\Phi(q)\rangle$
 $q \equiv$ collective degrees of freedom

- Variational principle: $\delta \frac{\langle \Phi | H | \Phi \rangle}{\langle \Phi | \Phi \rangle} = 0 \Leftrightarrow \text{diag} [\text{span}(\{|\Phi(q)\rangle\})]$

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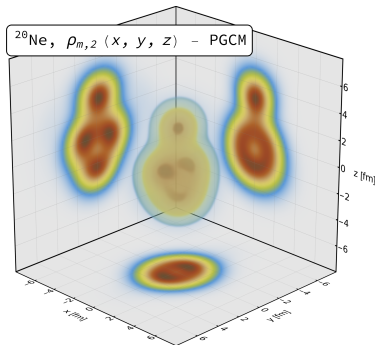
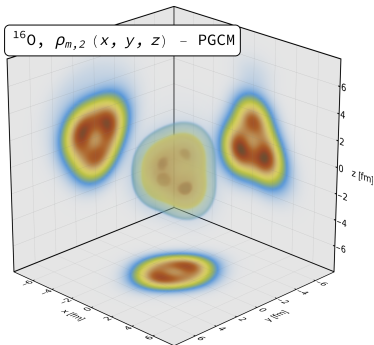
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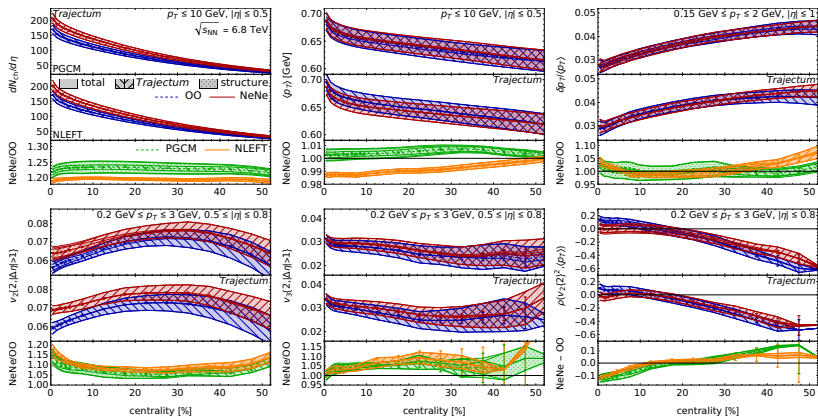
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- Large-scale computations using numerical suite TAURUS on Topaze supercomputer (CEA/CCRT)



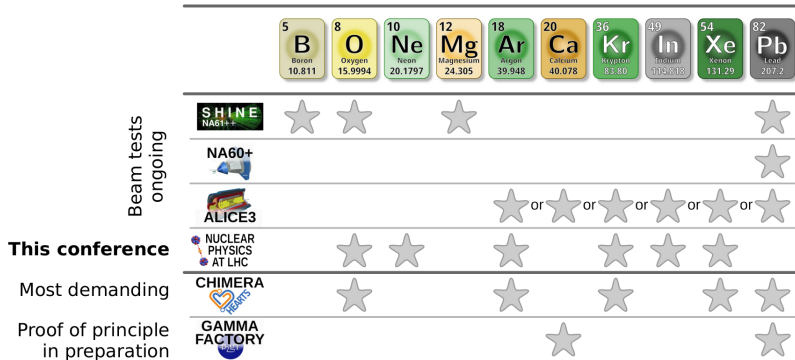


- Determine average deformation of PGCM ground state: \bar{q}
- One-body density: $\rho_m(x, y, z) = \sum_{st} \frac{\langle \Phi(\bar{q}) | a_{xyzst}^+ a_{xyzst} P^Z P^N | \Phi(\bar{q}) \rangle}{\langle \Phi(\bar{q}) | P^Z P^N | \Phi(\bar{q}) \rangle}$
- Sample ρ_m



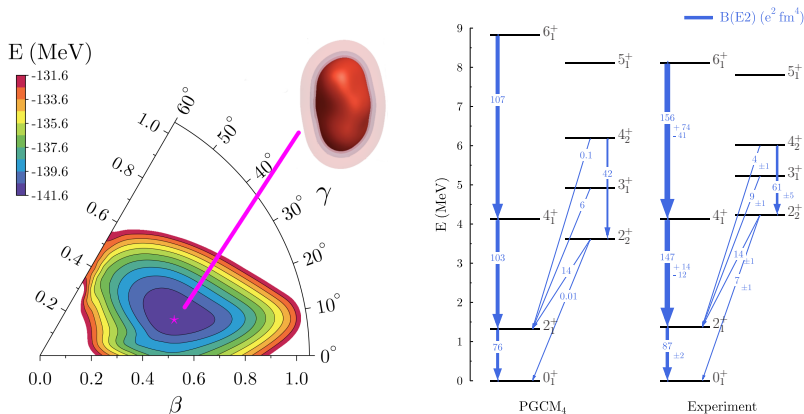
- Impact of ^{20}Ne and ^{16}O structure
- NLEFT and PGCM in good agreement

Experimental landscape at CERN



Presentation by Maciej Slupecki (CERN), Beijing, 2024

⇒ All these nuclei are within the reach of ab initio nuclear structure



Bally *et al.*, EPJA 60, 62 (2024)

- Ground state exhibits large intrinsic triaxial deformation
- Excellent description using χ EFT Hamiltonian

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- Fully exploit the versatility of LHCb/SMOG2