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MMM

Office of Science

Overview of recent UPC results from ALICE with a focus on incoherent measurements Simone Ragoni

Creighton University, USA

OUR UPC PHYSICS PROGRAM IS THE FORERUNNER OF EIC





- Introduction to ultra-peripheral collisions (UPC)
- The ALICE detector
- \bullet Results on exclusive and dissociative J/ ψ in p-Pb UPC
- \bullet Results on coherent and incoherent J/ ψ in UPC Pb-Pb
- Measurements of the energy dependence of the photonuclear cross sections





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Introduction to ultraperipheral collisions (UPC) Creighton



- Large impact parameter (beyond the reach of the strong interaction)
- Vector meson production
- e.g. ρ^0 , J/ψ , $\psi(2S)$

Only QED involved at this vertex! $\frac{d\sigma^{T}(\gamma p \rightarrow J/\Psi + p)}{dt} = \frac{|M|^{2}}{16\pi s^{2}} LO$ $= [F_{N}^{2G}(t)]^{2} \frac{\alpha_{e}^{2} \Gamma_{ee}^{J} m_{J}^{3}}{3\alpha_{e.m.}} \pi^{3} \left[\bar{x} G(\bar{x}, \bar{q}^{2}) \frac{2\bar{q}^{2} - |q_{t}^{J}|^{2}}{(2\bar{q}^{2})^{3}} \right]^{2}$ Ryskin: Z. Phys. C 57, 89-92 (1993)

Hard scale assured by high mass states i.e. J/ψ , $\psi(2S)$

- Coherent photoproduction: photon couples with the entire nucleus
- Incoherent photoproduction: photon couples with a single nucleon only
- Different average p_T of the vector mesons for the two processes

ALICE

Introduction to ultraperipheral collisions (UPC) $\frac{Creighton}{UNIVERSITY}$



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|t| the square of the momentum transferred between the incoming and outgoing target nucleus

W the centre-of-mass energy of the photon-target system

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Physics of ultra-peripheral collisions (UPC)









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Exclusive J/ψ in p-Pb

 $x = e^{\pm|y|} \underline{M}_{\mathrm{J}/\psi}$ • Probing Bjorken- $x \sim 10^{-5}$ with ALICE

- power-law growth of cross-sections \rightarrow power-law growth of gluon distributions down to $x \sim 10^{-6} \rightarrow 10^{-6}$ clear signs of gluon saturation
- ALICE points: forward, semiforward and midrapidity configurations
 - Forward: two muons in the spectrometer
 - Semiforward: one in the spectrometer, one in the central barre
 - Midrapidity: two muons/electrons in the central barrel

Bjorken-*x* 10^{-3} 10^{-1} 10^{-2} 10^{-5} 10^{-4} (dh) (q+ $\psi \rightarrow J/\psi$ +p) (nb) ALICE p-Pb $\sqrt{s_{NN}}$ = 8.16 TeV 10^{3} ALICE p-Pb $\sqrt{s_{NN}}$ = 5.02 TeV LHCb pp \sqrt{s} = 7 TeV and 13 TeV (W+ solutions) LHCb pp \sqrt{s} = 7 TeV and 13 TeV (W- solutions) Fixed target (E401, E516, E687) H1 **ZEUS** 10² **JMRT NLO** CCT Power-law fit to ALICE data

ALICE, Phys.Rev.D 108 (2023) 11, 112004

Eur. Phys. J. C (2019) 79: 402 (ALICE midrapidity and semiforward), Phys. Rev. Lett. 113 no. 23, (2014) 232504 (ALICE forward)

10²

30 40

20

10

ALI-PUB-568478

 2×10^{2}

 2×10^{3}

10³

 $W_{\gamma p}$ (GeV)

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Exclusive and dissociative J/ψ in p-Pb







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Coherent vs incoherent J/ψ



- Coherent (dimuon $p_{\rm T} < 0.2 \text{ GeV}/c$) photon couples to entire nucleus coherently
- Incoherent J/ ψ features a much wider $p_{\rm T}$ distribution photon interacts with a single nucleon of the target nucleus

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1. Angular distributions



4.
$$W_{\gamma Pb,n}^2 = m \sqrt{s_{NN}} e^{-y}$$



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1. Angular distributions

2. **y**

4. $W_{\gamma Pb,n}^2 = m \sqrt{s_{NN}} e^{-y}$

Coherent J/ ψ cross section

- ALICE data exhibit moderate nuclear shadowing
- Nuclear suppression factor

$$S_{\rm Pb}(y \sim 0) = \sqrt{\frac{d\sigma}{dy}} \frac{d\sigma}{dy}_{IA} = 0.64 \pm 0.04$$

- IA = impulse approximation (no nuclear effects)
- $S(W_{\gamma p})$ nuclear suppression factor provides a way to test the consistency of the data with the available nuclear and nucleon PDFs and to measure the nuclear shadowing factor







1. Angular distributions



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Complete view of J/ψ t-dependence



Fig. 1. Coherent (thick) and incoherent (thin lines) diffractive J/Ψ production cross section as a function of *t*, with (solid lines) and without (dashed lines) subnucleonic fluctuations. The band shows statistical uncertainty of the calculation.

ALI-DER-559110



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Techniques for the photon direction ambiguity



Neutron emission:

•
$$x = \frac{M_{\rm VM}}{\sqrt{s_{\rm NN}}} \cdot e^{\pm y}$$

• Ambiguity due to sign in the rapidity of the photon emitter $\rightarrow 10^{-2}$, 10^{-5}



- Using the neutron ZDCs on the A and C side to detect the neutrons!
- E.g. 0N0N: no neutrons on either ZDCs
- E.g. 0NXN: neutrons only on one side

 $\frac{d\sigma_{PbPb}^{0N0N}}{dy} = n_{0N0N}(\gamma, +y) \cdot \sigma_{\gamma Pb}(+y) + n_{0N0N}(\gamma, -y) \cdot \sigma_{\gamma Pb}(-y)$ $\frac{d\sigma_{PbPb}^{0NXN}}{dy} = n_{0NXN}(\gamma, +y) \cdot \sigma_{\gamma Pb}(+y) + n_{0NXN}(\gamma, -y) \cdot \sigma_{\gamma Pb}(-y)$

Guzey et al., Eur.Phys.J.C 74 (2014) 7, 2942 • Effectively leveraging on the impact parameter

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 $2R_{\rm Pb}$

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Coherent J/ ψ with neutron emission



ALICE, JHEP 10 (2023) 119

- First measurement of the energy dependence of the photonuclear cross section down to Bjorken- $x \sim 10^{-5}!$
- At low-*x* data favours both saturation and shadowing models
- New Run 2 results probe unprecedented Bjorken-*x* region like no other LHC experiment!



Neutron emission extends the range in energy being explored by about 300 GeV!

Coherent J/ ψ with neutron emission



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"ALICE DOES THE DOUBLE-SLIT" HOME.CERN, AUG 12TH, 2024





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Coherent ρ^0 with neutron emission



ALICE, arXiv:2405.14525 [nucl-ex], submitted to PLB

- Modulation increases as the impact parameter lowers
- ALICE results compatible with both theory and STAR results
- Modulation: linearly polarised photons + quantum interference at the fermi scale

ALICE Pb-Pb UPC $\sqrt{s_{NN}} = 5.02 \text{ TeV}$

 $Pb + Pb \rightarrow Pb + Pb + \rho^{0}$

z 0.35

0.3

0.2

0.1

0.05



 $p_{-} < 0.1 \text{ GeV}/c, |y| < 0.8$ 0.25 - ALICE ***** STAR Au–Au UPC $\sqrt{s_{NN}}$ = 200 GeV, p_{τ} < 0.06 GeV/c STAR U–U UPC $\sqrt{s_{NN}}$ = 193 GeV, p_{τ} < 0.06 GeV/c ALI-PUB-571042 H. Xing et al. W. Zhao et al. 0.15 Xn0n + 0nXn XnXn 0n0n ALI-PUB-571047

https://home.cern/news/news/physics/alice-does-double-slit Same technique: neutron emission classes \rightarrow impact parameter range

ALICE IN THE FUTURE UPGRADES IN RUN 3 AND 4

ALICE in Run 3 and 4

- Significant increase in integrated lumi from 1 nb⁻¹ for Run 2 to 13 nb⁻¹ for Run 3 and Run 4 together
- Great increase in statistics with continuous readout
- Uncertainties for nuclear suppression factor expected to be at the level of 4%
 - Nuclear shadowing studied as a function of x and Q^2
- New measurements e.g. bottomonium states
- MFT in Run 3, FoCal in Run 4!

		1010	
	σ	Central 1	Forward 1
Meson		Total	Total 1
$\rho \to \pi^+ \pi^-$	5.2b	5.5 B	4.9 B
$\rho' \to \pi^+ \pi^- \pi^+ \pi^-$	730 mb	210 M	190 M
$\phi \to \mathrm{K}^+\mathrm{K}^-$	0.22b	82 M	15 M
${ m J}/\psi o \mu^+\mu^-$	1.0 mb	1.1 M	600 K
$\psi(\mathrm{2S}) o \mu^+ \mu^-$	30µb	35 K	19 K
$ m Y(1S) ightarrow \mu^+ \mu^-$	2.0 µb	2.8 K	880

PhPh

CERN Yellow Rep. Monogr. 7 |y| < 0.9 2.5 < |y| < 4(2019) 1159-1410, arXiv 1812.06772 ALICE Simulation, $Pb + Pb \rightarrow Pb + Pb + V$ $\sqrt{s_{NN}} = 5.5 \text{ TeV}, L = 13 \text{ nb}^{-1}$ 0.8 (х) о.6 На 0.4 CMS Y(1S) pseudodata EPS09LO, $Q = m_{V(1S)}/2$ ALICE Y(1S) pseudodata EPS09LO, $Q = m_{10}(2S)/2$ 0.2 ALICE $\psi(2S)$ pseudodata EPS09LO, $Q = m_{1/1}/2$ ALICE J/ψ pseudodata 10^{-4} 10^{-3} 10^{-2} 10 Х

ALI-SIMUL-313259

40

Dissociative J/ ψ in Run 3 and 4 (with FOCAL)

- Already in Run 3, but with FOCAL much easier to measure higher energies (using the channel J/ψ $\rightarrow ee$)
- CCT model features gluonic hotspots
- Significant reduction at higher energies would imply the onset of the saturation regime



Exclusive J/ ψ and ψ' in Run 3 and 4 (with FOCAL)

- The ratio is also a very sensitive observable
- Both Run 3 and 4 are needed to find the onset of the saturation regime



A. Bylinkin, J. Nystrand, D. Tapia Takaki, 2023 J. Phys. G: Nucl. Part. Phys. **50** 055105





Backup slides