

Event properties and hydro in small and large systems

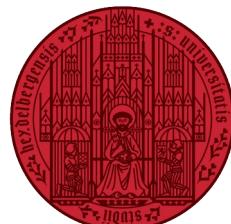
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Institute for Theoretical Physics, Heidelberg University

September 2, 2024

QCD Challenges from pp to AA collisions



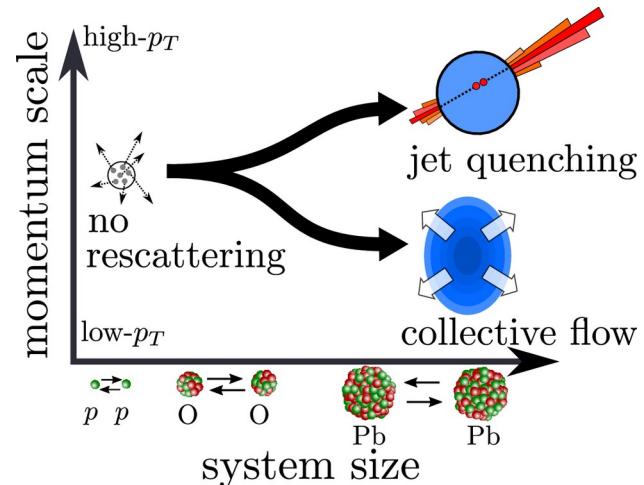
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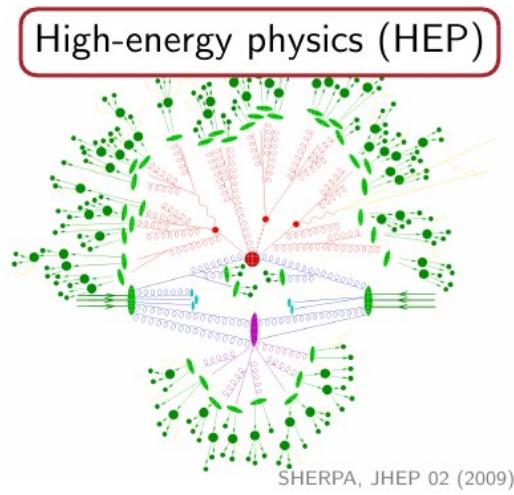


Outline

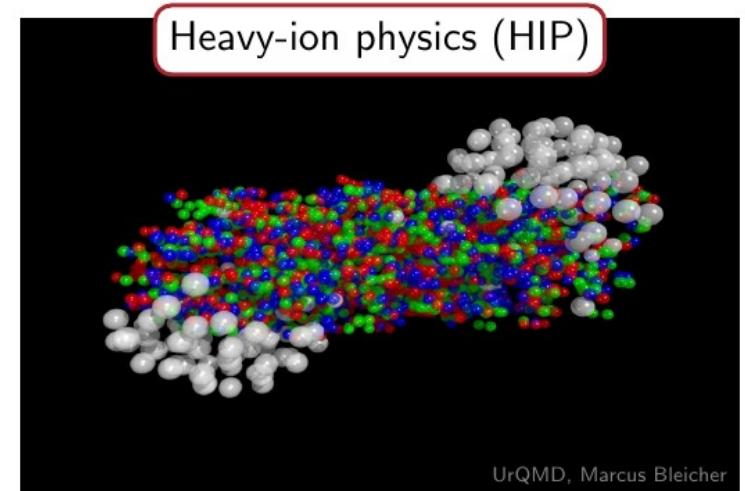
- Emergent phenomena in many-body systems.
- Causality and stability of relativistic hydrodynamics
- Stochastic viscous hydrodynamics
- Collectivity with few ultracold fermionic atoms

Emergent phenomena in complex systems

- HEP: concentrate higher energy in smaller and smaller volume.
- HIP: distribute high energy or high nucleon density over a relatively large volume. – T.D. Lee, 1974, Bear Mountain workshop



more scatterings →



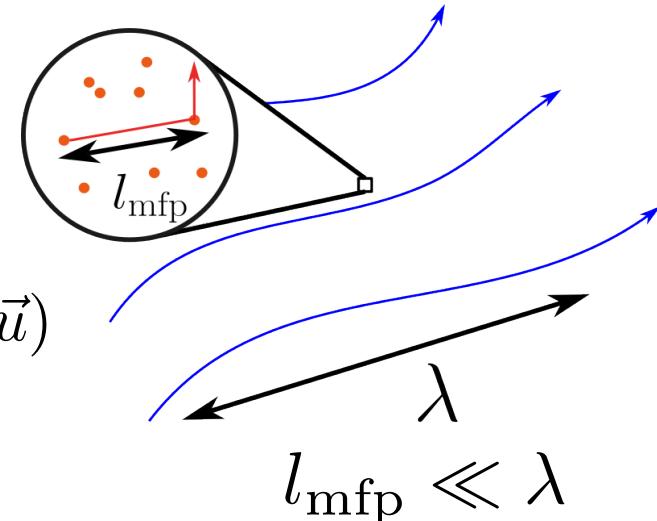
“More is different” – P.W. Anderson (1972)

Hydrodynamics

Navier-Stokes equation for fluid velocity \vec{u}

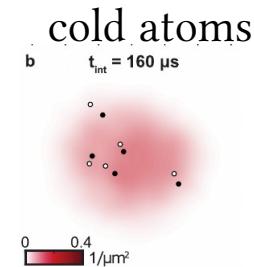
$$\rho \left(\partial_t + \vec{u} \cdot \vec{\nabla} \right) \vec{u} = -\vec{\nabla} p + \eta \nabla^2 \vec{u} + \left(\frac{1}{3} \eta + \zeta \right) \vec{\nabla} (\vec{\nabla} \cdot \vec{u})$$

density pressure shear and bulk viscosity

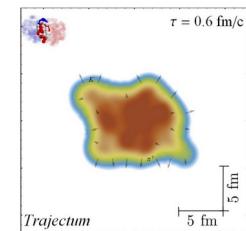


Universal effective theory

Even in small systems!



quark gluon plasma



Shear viscosity

$$\eta \sim n \times \langle p \rangle \times l_{\text{mfp}}$$

Large mean free path \rightarrow large viscosity

Small mean free path \rightarrow small viscosity

$$\Delta x \Delta p \geq \hbar/2$$

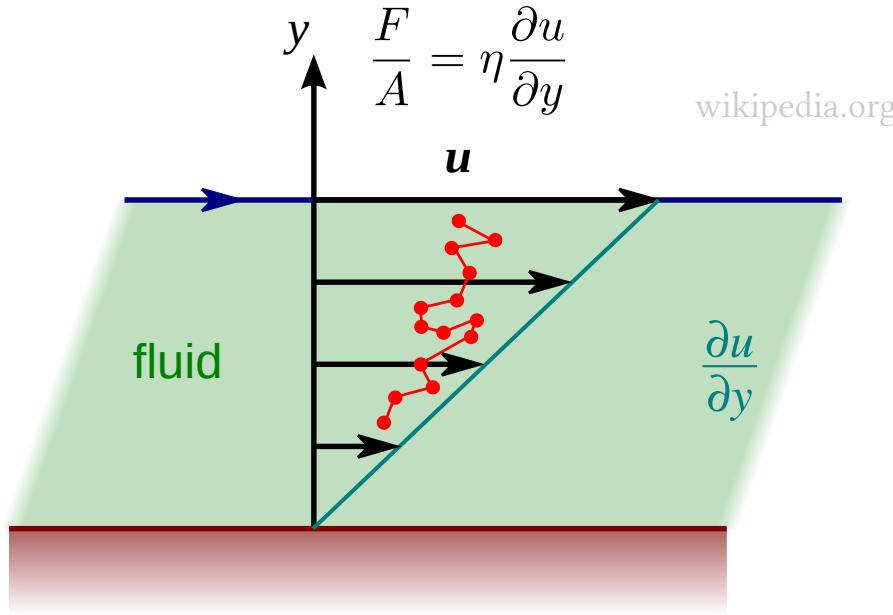
Viscosity in strongly interacting QFT

Kovtun, Son, Starinets, PRL, 2005

$$\frac{\eta}{s} = \frac{\hbar}{4\pi k_B} \approx 0.08 \frac{\hbar}{k_B}$$

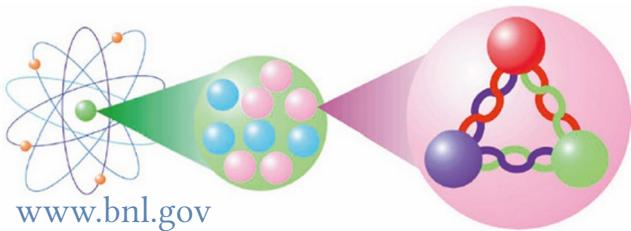
“perfect” fluid

entropy $s \sim nk_B$

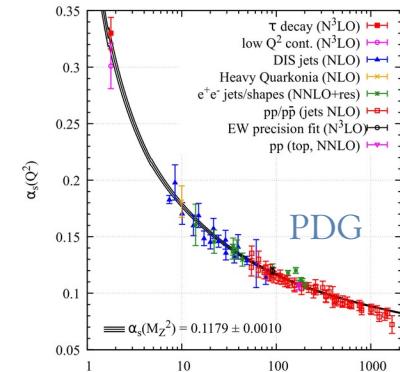


Strongly interacting quantum systems

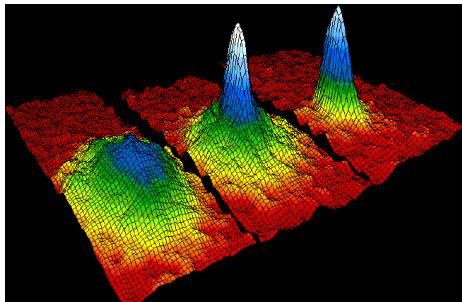
Quark Gluon Plasma



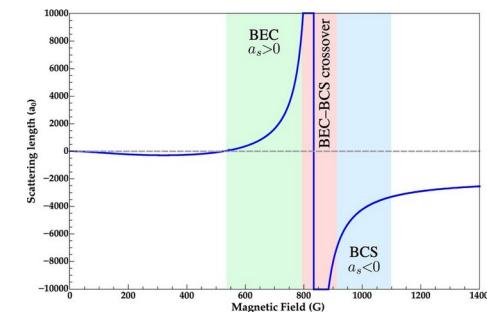
QCD coupling α_s runs
with energy scale



Ultracold quantum gases



Interaction strength diverges
near Feschbach resonance



Hernández-Rajkovic et al., 2008.05046

How many particles make a fluid?

wikipedia.org

Sorites ('heap') paradox:

If n grains of sand is a heap, then $n-1$ is also a heap
⇒ a single grain of sand is a heap?



Duck test:

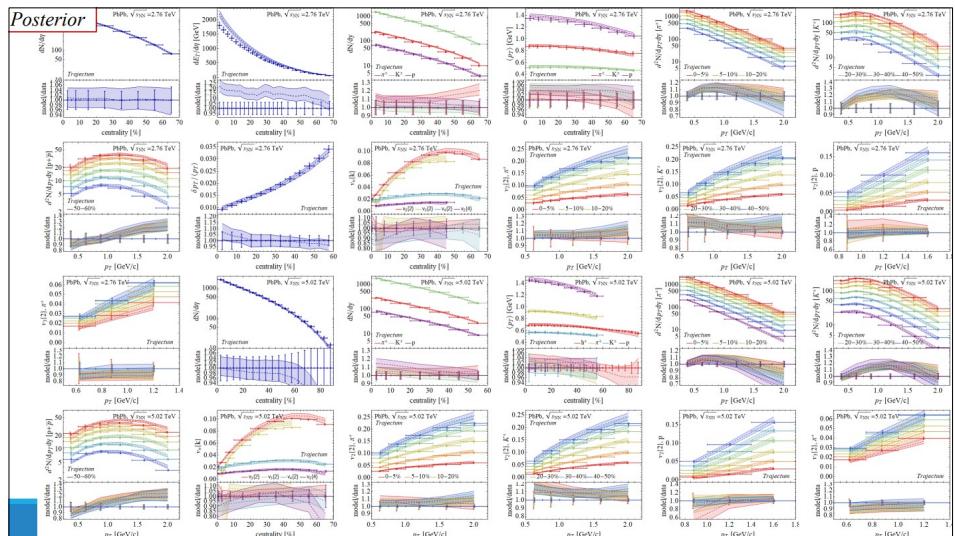
If it looks like a duck, swims like a duck,
and quacks like a duck, then it probably is a duck.



Can a few-particle system **behave** like a fluid?

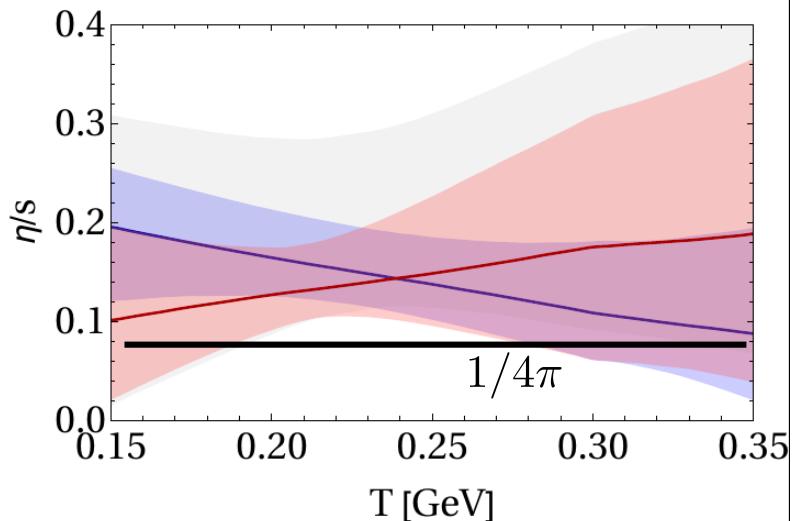
Shear viscosity of Quark Gluon Plasma

Multi-observable Bayesian fits



©Wilke van der Schee

Fitted viscosity to entropy ratio



Nijs, van der Schee, PRL 2022

Quark Gluon Plasma behaves like a nearly “perfect” fluid!

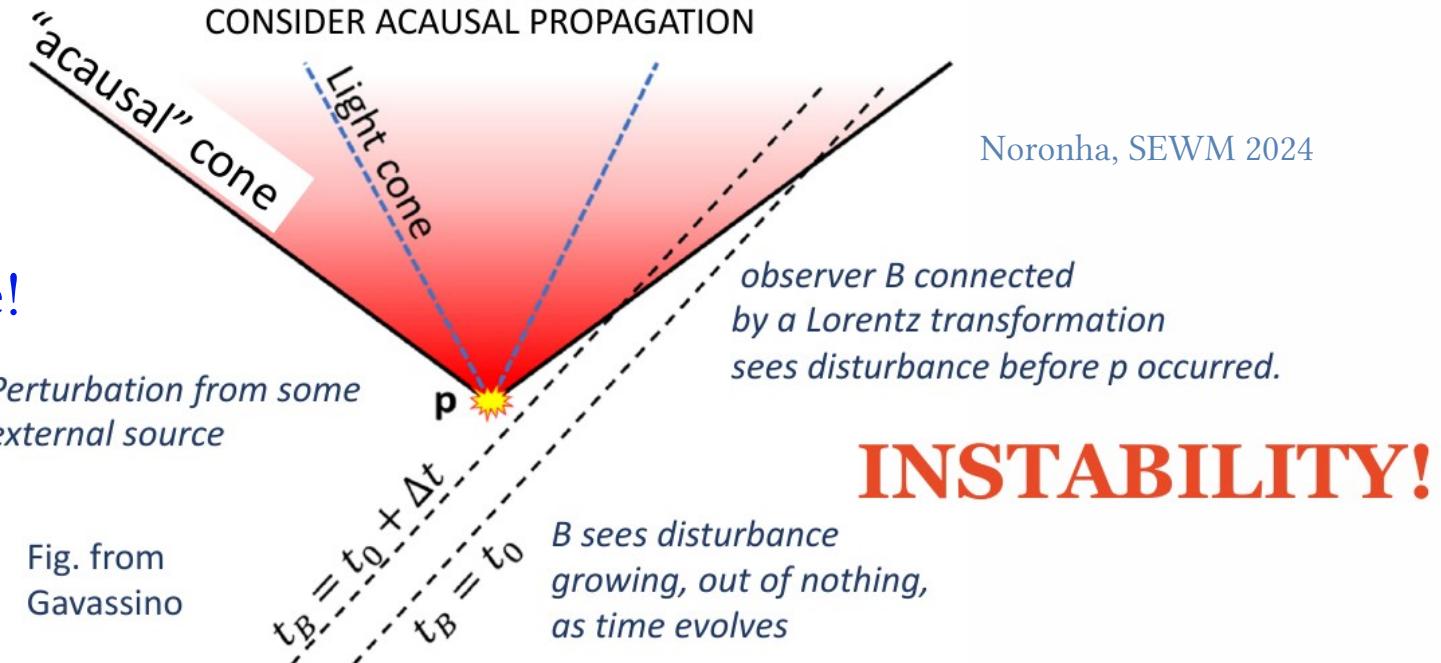
Challenge I: causality and stability

Relativistic Navier-Stokes equation

$$\rho \left(\partial_t + \vec{u} \cdot \vec{\nabla} \right) \vec{u} = -\vec{\nabla} p + \eta \nabla^2 \vec{u} + \left(\frac{1}{3} \eta + \zeta \right) \vec{\nabla} (\vec{\nabla} \cdot \vec{u})$$

Second-order time-derivatives: $(g^{\alpha\beta} + u^\alpha u^\beta) \partial_\alpha \partial_\beta u^\mu$

Naive formulation
is acausal and unstable!



Muller-Israel-Stewart (2nd order hydro)

from TRAJECTUM, 2010.15134

Relaxation equations for shear
and bulk pressures

$$D\Pi = -\frac{1}{\tau_\Pi} [\Pi + \zeta \nabla \cdot u + \delta_{\Pi\Pi} \nabla \cdot u \Pi - \lambda_{\Pi\pi} \pi^{\mu\nu} \sigma_{\mu\nu}], \quad (6)$$

New transport coefficients

$$\Delta_\alpha^\mu \Delta_\beta^\nu D\pi^{\alpha\beta} = -\frac{1}{\tau_\pi} [\pi^{\mu\nu} - 2\eta\sigma^{\mu\nu} + \delta_{\pi\pi} \pi^{\mu\nu} \nabla \cdot u - \phi_7 \pi_\alpha^{\langle\mu} \pi^{\nu\rangle\alpha} + \tau_{\pi\pi} \pi_\alpha^{\langle\mu} \sigma^{\nu\rangle\alpha} - \lambda_{\pi\Pi} \Pi \sigma^{\mu\nu}]. \quad (7)$$

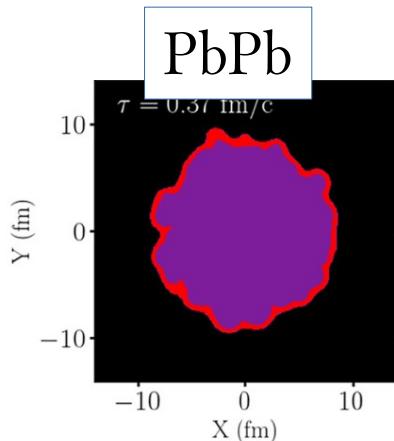
Nonlinear constraints for dissipative fluxes

$$c_s^2 + \frac{4}{3} \frac{\eta}{\tau_\pi(e+p)} + \frac{\zeta}{\tau_\Pi(e+p)} \leq 1$$

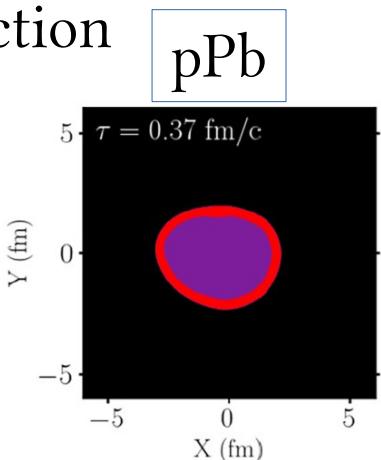
Bemfica, Disconzi, Noronha, PRL (2019), Bemfica, Disconzi, Hoang, Noronha, Radosz, PRL (2021)

Causality violation in small and large systems

Only small percentage of total energy in acausal cells

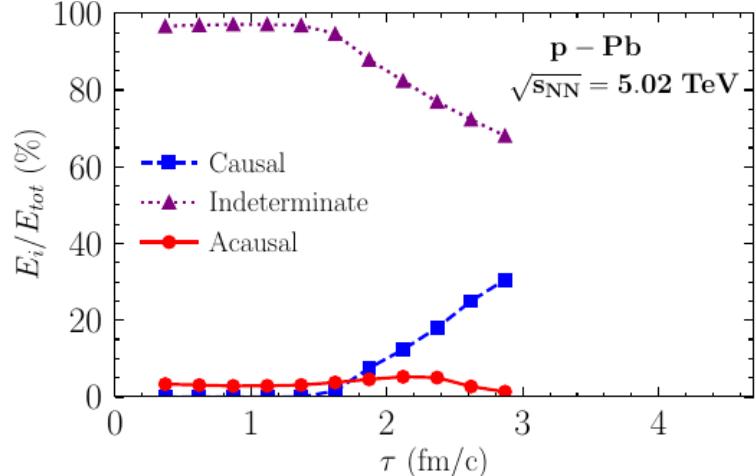
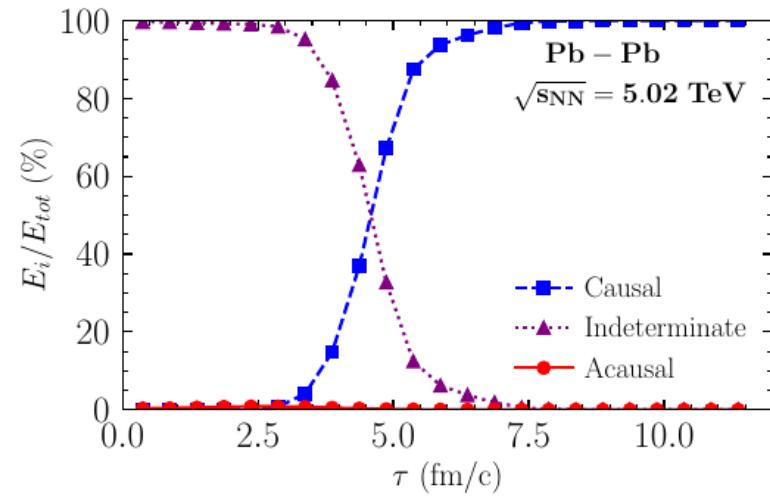


Might still affect viscosity extraction
Especially bulk viscosity



see talk by Renata

Krupczak et al. (ExTrEMe) PRC 2024



BDNK theory (1st order hydro)

Bemfica, Disconzi, noronha PRD (2018), Kovtun JHEP (2019)

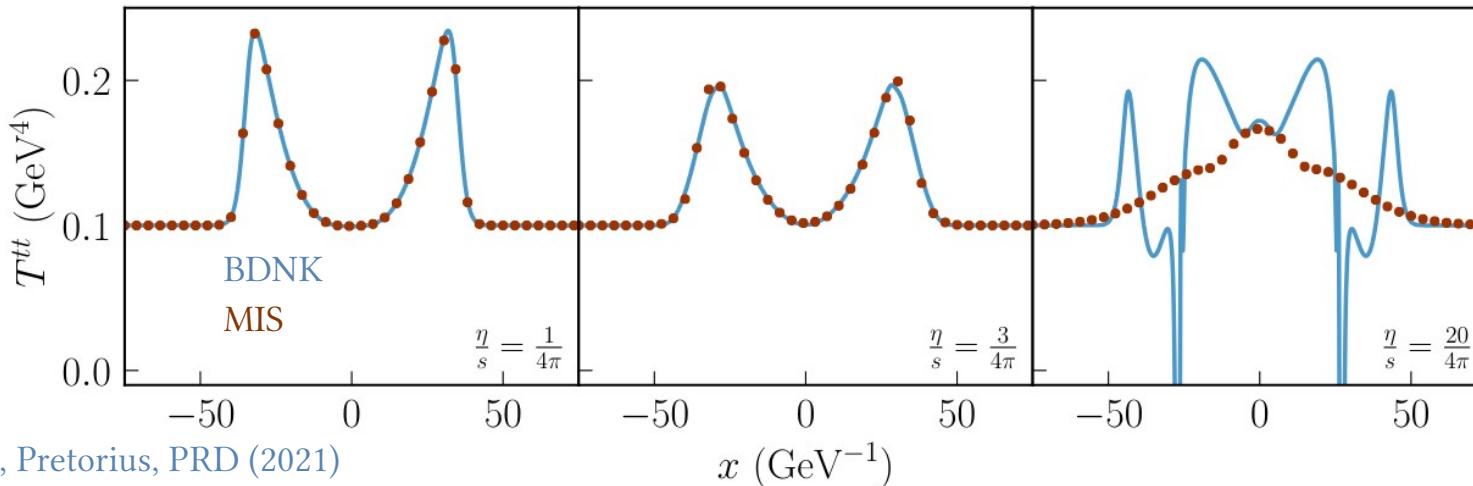
Use freedom of hydrodynamic frame to obtain causal theory

→ additional transport coefficients

First proof of causality and stability

First (1+1D) numerical simulations

$$\mathcal{E} = \epsilon + \tau_\epsilon [u^c \nabla_c \epsilon + \rho \nabla_c u^c]$$
$$\mathcal{P} = P - \zeta \nabla_c u^c + \tau_P [u^c \nabla_c \epsilon + \rho \nabla_c u^c]$$
$$\mathcal{Q}^a = \tau_Q \rho u^c \nabla_c u^a + \beta_\epsilon \Delta^{ac} \nabla_c \epsilon + \beta_n \Delta^{ac} \nabla_c n$$
$$\mathcal{T}^{ab} = -2\eta \sigma^{ab} \equiv -2\eta \nabla^{< a} u^{b >}$$



Density frame

Armas, Jain, Scipost (2011), Başar, Bhambure, Singh, Teaney, 2403.04185

Consider relativistic diffusion equation

$$\partial_\mu J^\mu = 0, \quad J^\mu = \cancel{n} u^\mu + j_D^\mu, \quad j_D^\mu = -D \Delta^{\mu\nu} \partial_\nu n$$

density in fluid rest frame

$$\text{Density frame: } N = J^0 \quad \partial_t N + \vec{\nabla} \cdot \vec{J} = 0$$

$$\text{Using ideal equations: } \vec{J} = N \vec{v} + (D/\gamma) (\vec{\nabla} N - \vec{v} (\vec{v} \cdot \vec{\nabla}) N)$$

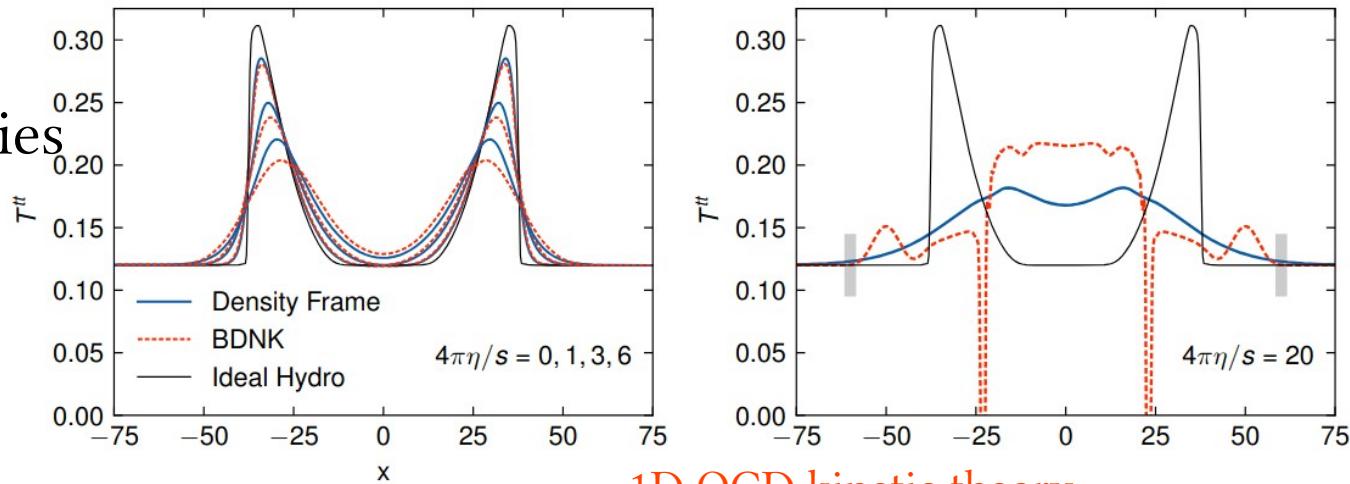
density in a lab frame fluid 3-velocity

Strictly first order, numerically stable and no new coefficients,
...but does not obey Lorentz causality

Comparison with BDNK and kinetic theory

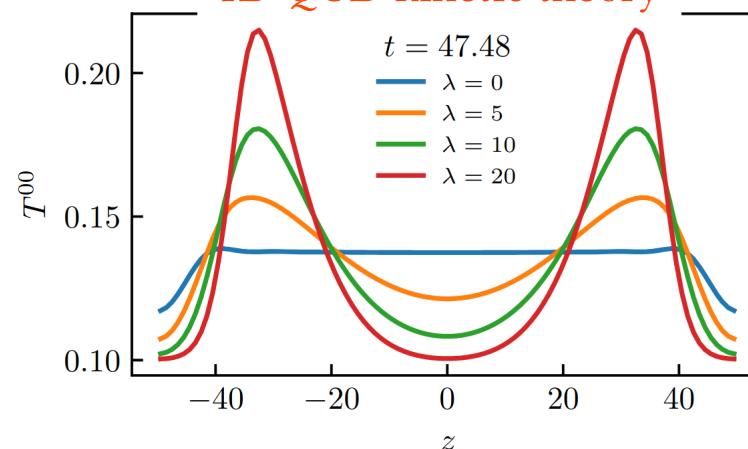
For $\eta/s \rightarrow 0$ agreement
with BDNK, but discrepancies
for large η/s

Paquet, Bhambure, AM, Singh, Teaney, Zhou (in preparation)



Kinetic theory smoothly interpolates
to free streaming at $\eta/s \rightarrow \infty$.
→ use to test density-frame results

1D QCD kinetic theory

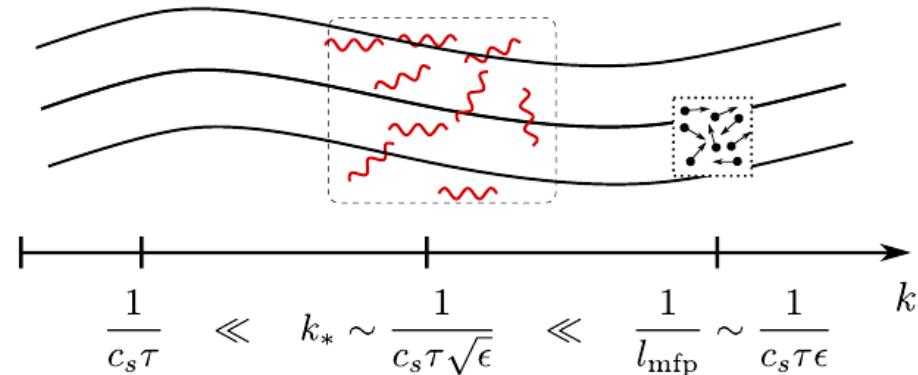


Challenge II: stochastic noise

Hydrodynamics with noise

No dissipation without fluctuations!

$$T^{\mu\nu} = \underbrace{eu^\mu u^\nu + \mathcal{P}(e)\Delta^{\mu\nu}}_{\text{ideal}} - \underbrace{\eta\sigma^{\mu\nu}}_{\text{visc}} + \underbrace{\xi^{\mu\nu}}_{\text{noise}}$$



$$\langle \xi^{\mu\nu}(x)\xi^{\rho\sigma}(x') \rangle \equiv 2T\eta\delta_{xx'}^4 \left[\Delta^{\mu\rho}\Delta^{\nu\sigma} + \Delta^{\nu\rho}\Delta^{\mu\sigma} - \frac{2}{3}\Delta^{\mu\nu}\Delta^{\rho\sigma} \right]$$

Hydrodynamic noise should be important in small systems

- Thermal fluctuations renormalize EOS and transport coefficients
- Fluctuations grown near critical points

Stickiness of sound and long-time tails

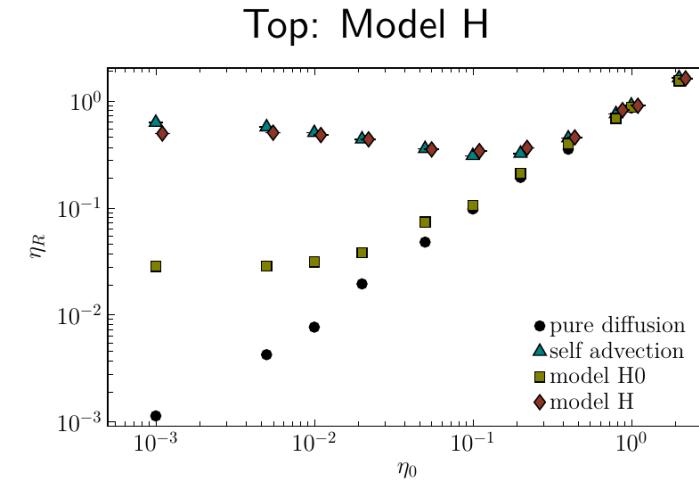
Kovtun, Moore and Romatschke PRD (2011)

Thermal noise contributes to viscosity

$$\eta_R = \eta + \frac{7}{60\pi^2} \frac{\rho T \Lambda}{\eta}$$

Schaefer SEWM (2024)

Lower limit of viscosity!



Chattopadhyay, Ott, Schaefer, Skokov, PRL (2024)

Out-of-equilibrium noise creates corrections between 1st and 2nd orders

$$\frac{\langle\langle \tau^2 T^{\eta\eta}(\tau) \rangle\rangle}{e+p} = \frac{p}{e+p} - \frac{4\gamma_\eta}{3\tau} + \frac{1.08318}{s (4\pi\gamma_\eta\tau)^{3/2}},$$

Akamatsu, AM, Teaney, PRC (2017,2018)

Stochastic hydrodynamics

Singh, Shen, McDonald, Jeon, Gale NPA (2018)

Adding noise to an update, e.g. Brownian particle

$$\partial_t p = -\gamma p + \xi$$

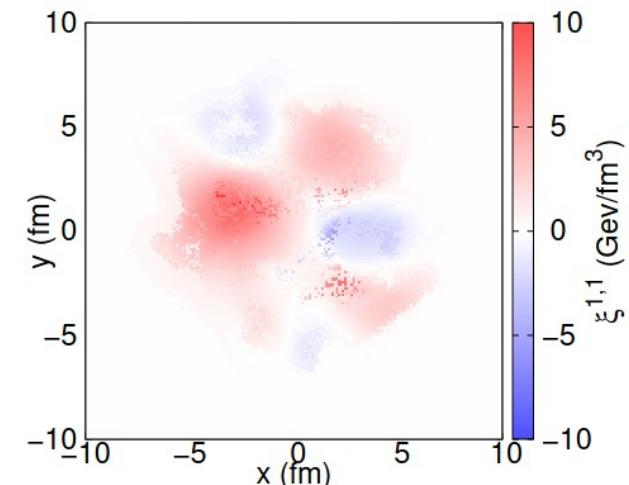
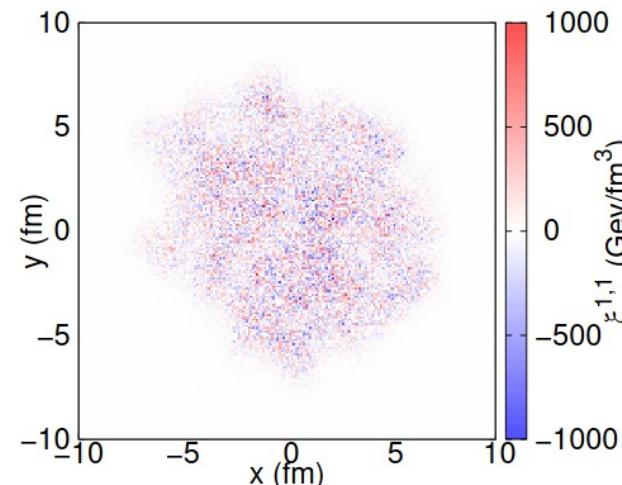
$$\langle \xi \xi \rangle = 2TM\delta(t - t')$$

$$p(t + \Delta t) = p(t) + \Delta t(-\gamma p(t) + \sqrt{2TM\gamma/\Delta t}\theta)$$

Noise contribution increases as $\Delta t \rightarrow 0$

Singh, Shen, McDonald, Jeon, Gale NPA (2018)

Need to filter small wavenumbers to do simulations



See also Sakai, Murase, Hirano, PRC 2020

Hydro kinetics

Deterministic equations for 2-point correlation functions

$$\partial_t \langle pp \rangle = -2\gamma(\langle pp \rangle - MT)$$

Akamatsu, AM, Teaney, PRC (2017,2018)

Heroic generalization to n-point functions
and general frames

An, Basar, Stephanov, Yee, PRL (2021)

So far solutions for simple flows

Akamatsu, AM, Teaney, PRC (2017,2018)

$$\partial_t (\bullet) = \Delta \circ \bullet$$

$$\partial_t (\bullet) = \Delta \circ \bullet + \Delta \circ \bullet$$

$$\partial_t (\bullet) = \Delta \circ \bullet + \Delta \circ \bullet$$

$$+ \Delta \circ \bullet + \Delta \circ \bullet$$

Metropolis algorithm

Başar, Bhambure, Singh, Teaney, 2403.04185

Implement noise with accept-reject algorithm

$$p(t + \Delta t) = p(t) + \sqrt{2TM\gamma\Delta t}\theta \quad P = \min(1, e^{-\Delta H})$$

dissipation

In one step implement both (on average) dissipation and noise

$$\langle \Delta p \rangle = -\gamma p \Delta t, \quad \langle \Delta p \Delta p \rangle = MT\Delta t^2$$

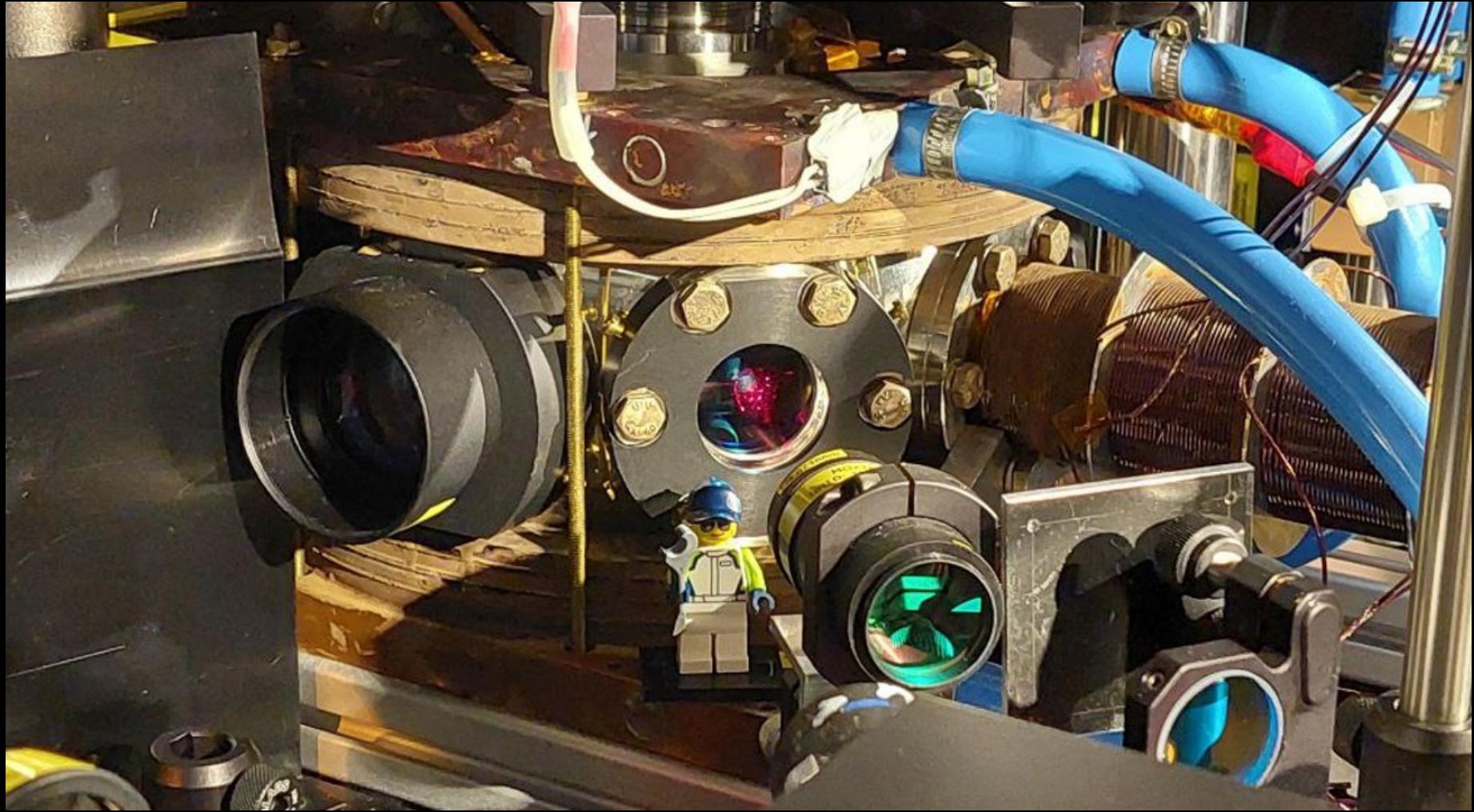
Generalizable to full hydrodynamic with random momentum updates.

Paquet, Bhambure, AM, Singh, Teaney, Zhou (in preparation)

Simulation of fluctuations near critical point (Model G and H)

Florio, Grossi, Soloviev, Teaney, PRD (2022) Chattopadhyay, Ott, Schaefer, Skokov, PRL (2024)

Challenge III: mesoscopic systems

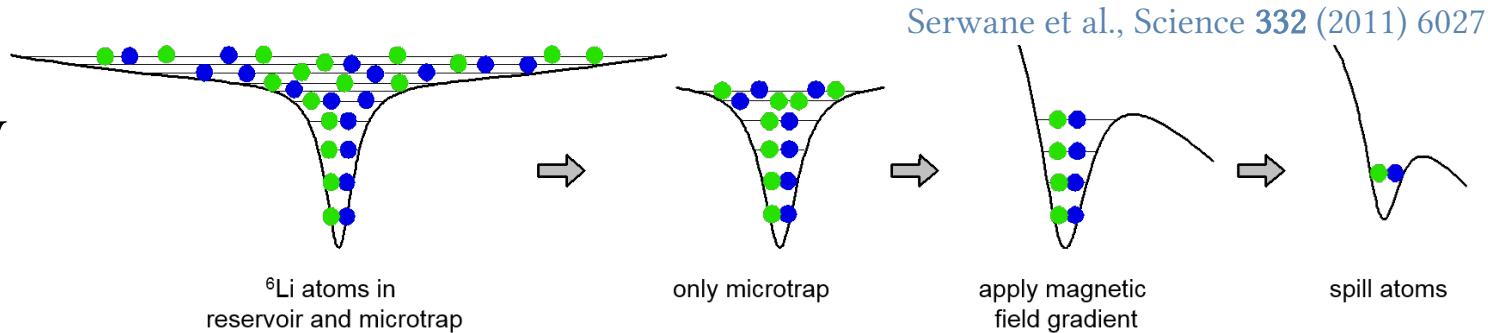


Brandstetter @ Jochim's lab, Heidelberg University

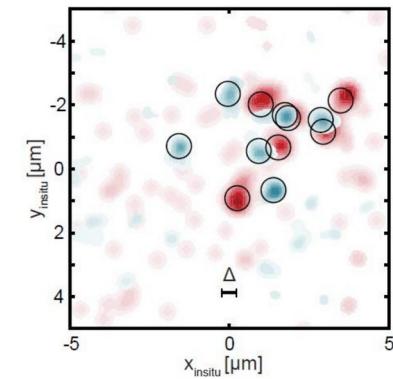
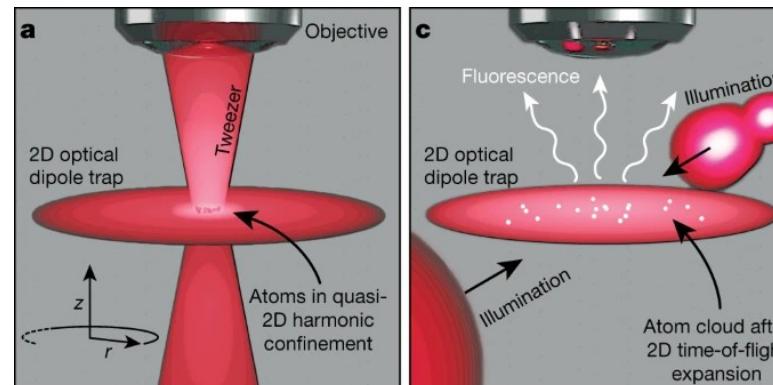
Aleksas Mazeliauskas, aleksas.eu

Lithium atoms in 2D harmonic trap

Preparation of few atoms in a trap



- time-resolved imaging
- spatial or momentum
- individual atoms
- controllable interactions

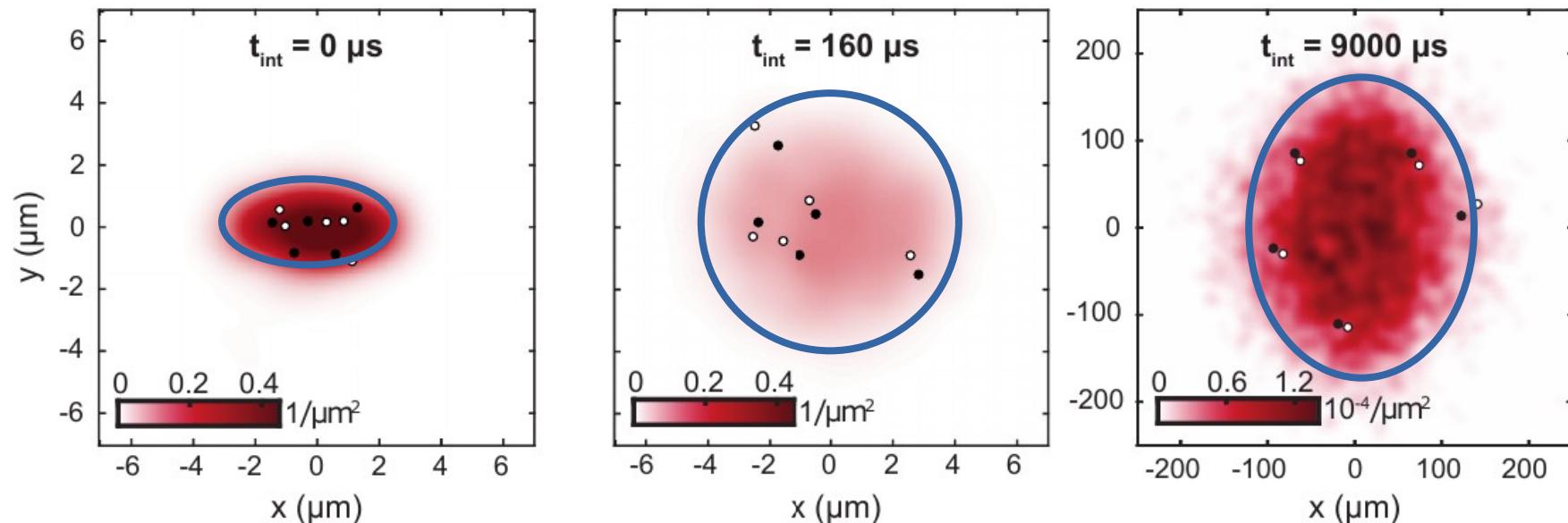


Geometry inversion of 10 atoms



Theory proposal: Flörchinger, Giacalone, Heyen, Tharwat, PRC 105 (2022) 4, 044908

Experimenta: Brandstetter, Lunt, Heintze, Giacalone, Heyen, Galka, Subramanian, Holten, Preiss, Floerchinger, Jochim arXiv:2308.09699

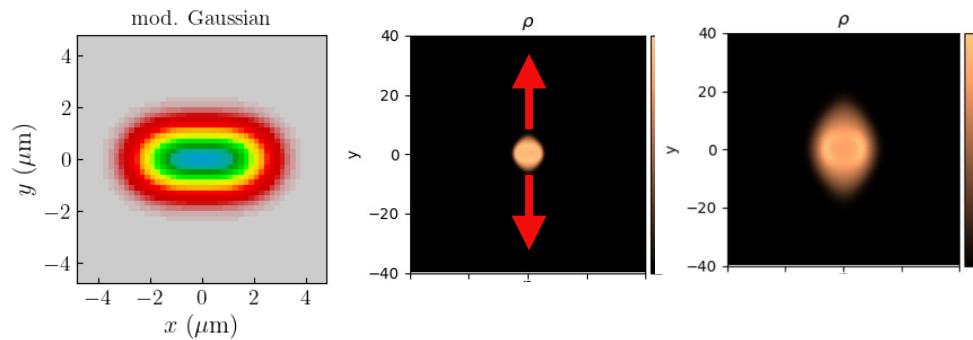


Elliptic flow of 10 Lithium atoms

Fluid dynamic description

Brandstetter, Lunt, et al. arXiv:2308.09699

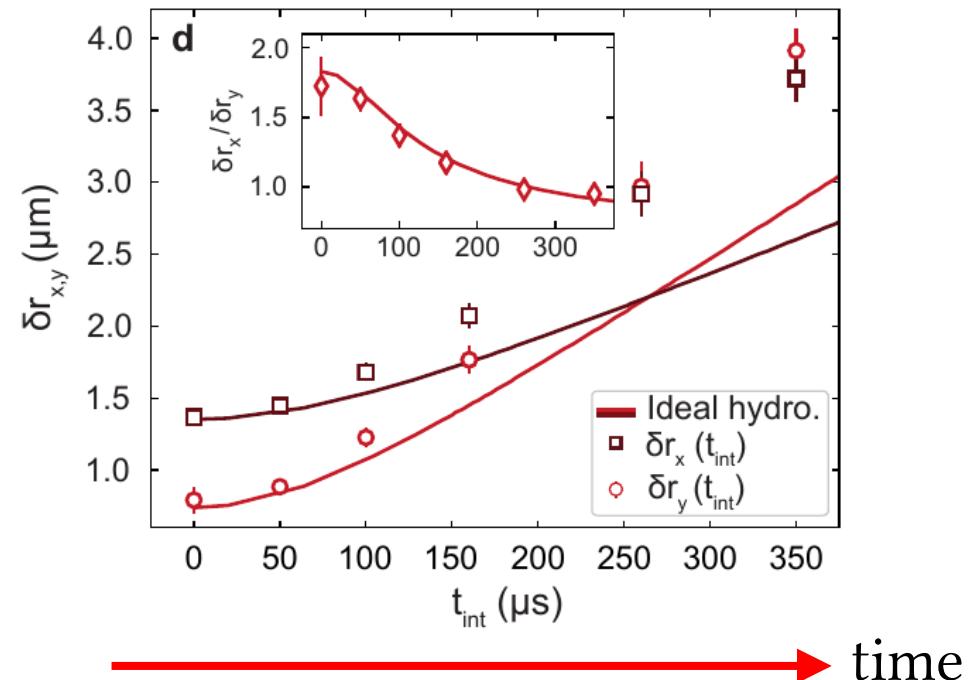
average geometry → expansion → inversion



ideal hydro simulation

initial conditions + equation of state

shape vs time

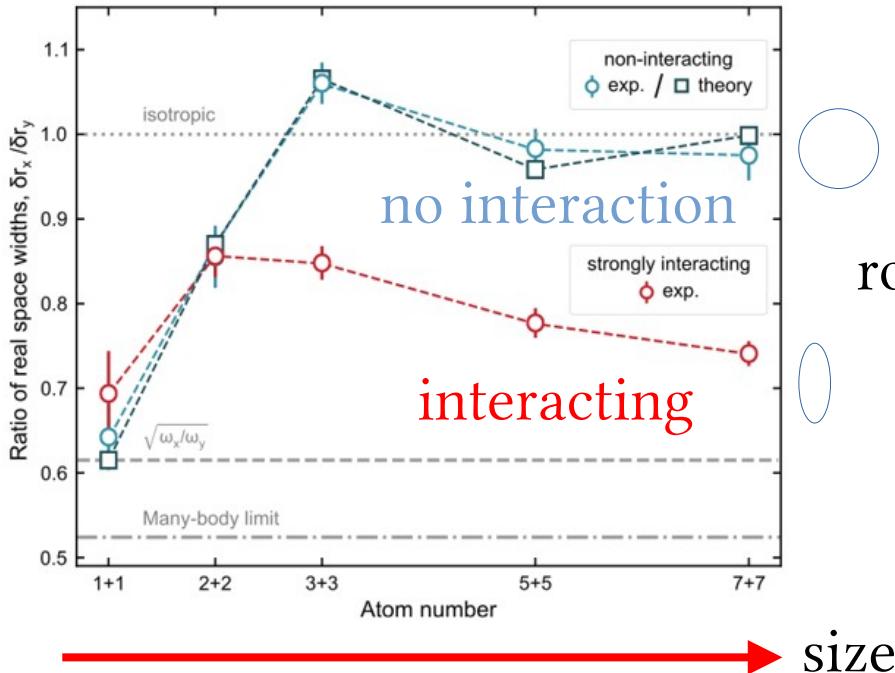
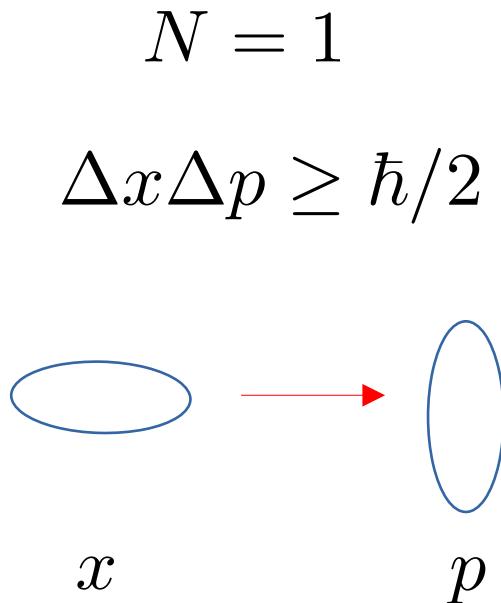


Good hydrodynamic description of 10 atoms!

System size dependence

Brandstetter, Lunt, et al. arXiv:2308.09699

final shape vs system size



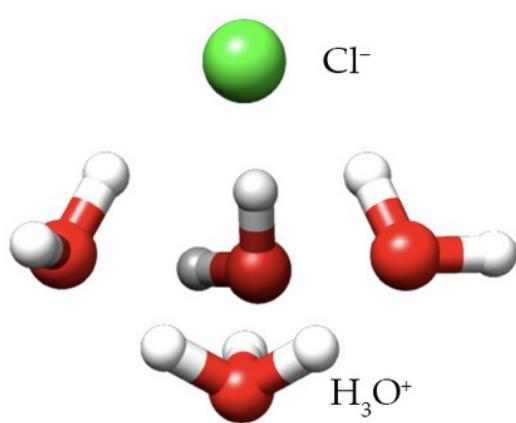
$N \gg 1$
round Fermi surface

Elliptic flow with > 5 atoms!

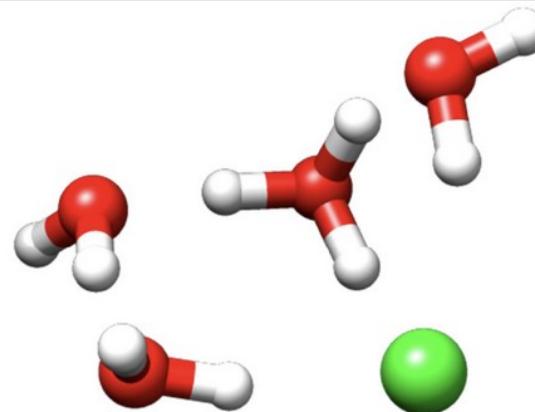
How many H₂O molecules make a droplet?

Miller, Physics Today 77 (8), 10–12 (2024)

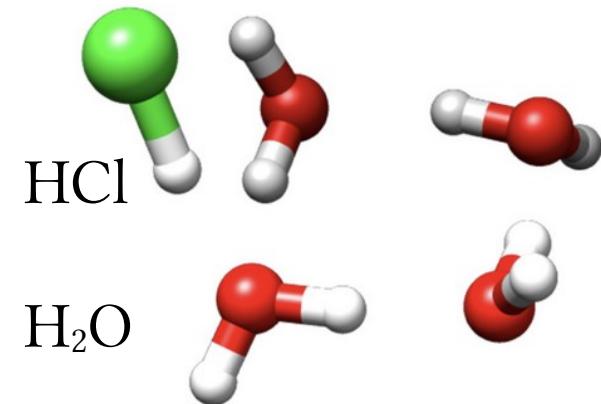
Xie, Tikhonov, Schnell, Science (2024)



Separated ion pair



Contact ion pair



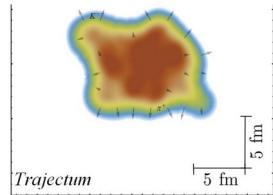
H_2O

Undissociated

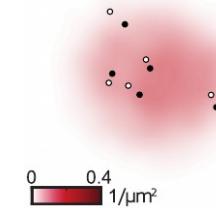
$n=5$ molecules are enough to dissolve HCl!

Similarities and differences between systems

quark gluon plasma



cold atoms



- Initial conditions: sampling of nucleons/atoms from wavefunction,
but no particle production in cold atoms
- Equation of state: many-body equilibrium EOS are used,
but cold atoms are at $T=0$.
- Hydrodynamic evolution: show geometry inversion,
but no dissipation measured in few-body cold atom systems
- Observables: event-by-event momentum observables in heavy-ions vs
ensemble averaged momentum/coordinate observables in cold atoms

Summary and outlook

Summary

QGP passes the “Duck test” as a near perfect fluid in large systems, but there are many challenges to interpretation in small systems

- Matching to microscopic dynamics (**not discussed**)
→ see talks by Victor and Nicolas
- Systematic uncertainties in transport of hydrodynamics
→ η/s , but of which hydrodynamics?
- Missing contributions of thermal noise
→ why does average hydrodynamics work so well?
- Sorites paradox
→ what can we learn from mesoscopic cold atom systems?

Postdoc position <https://academicjobsonline.org/ajo/jobs/27891>

