

Energy Loss in Small Collision Systems

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Kolbé and WAH, PRC100 (2019) [1511.09313]

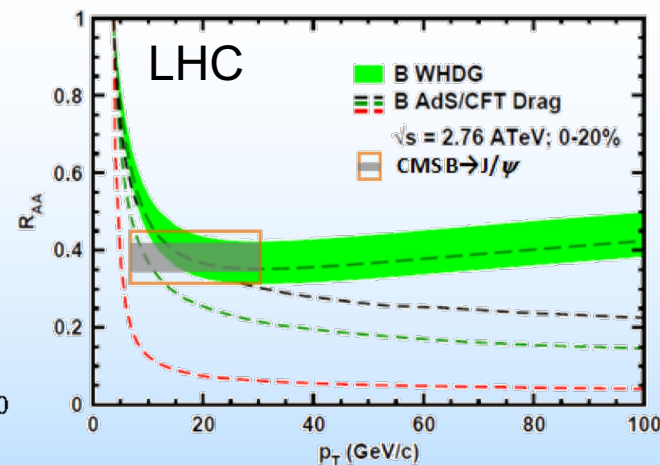
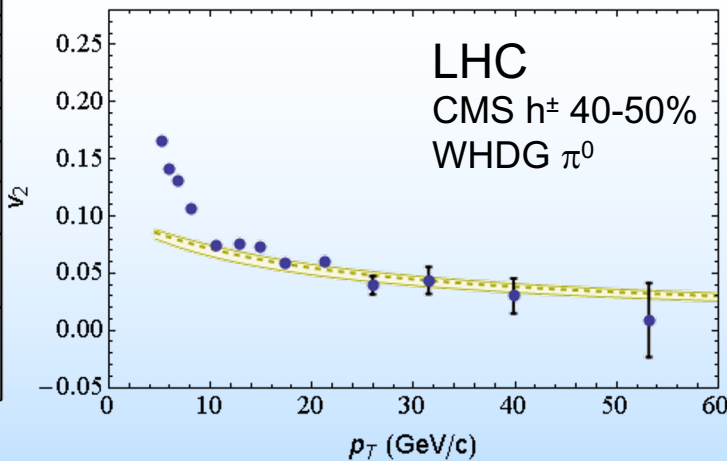
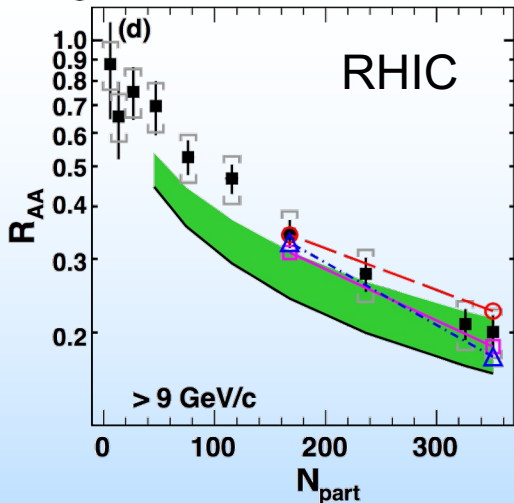
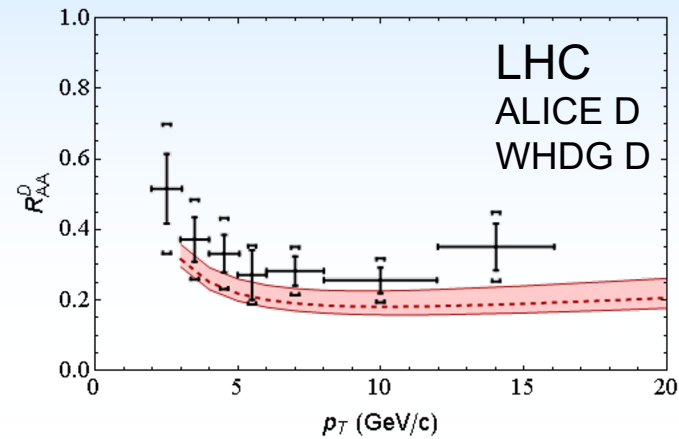
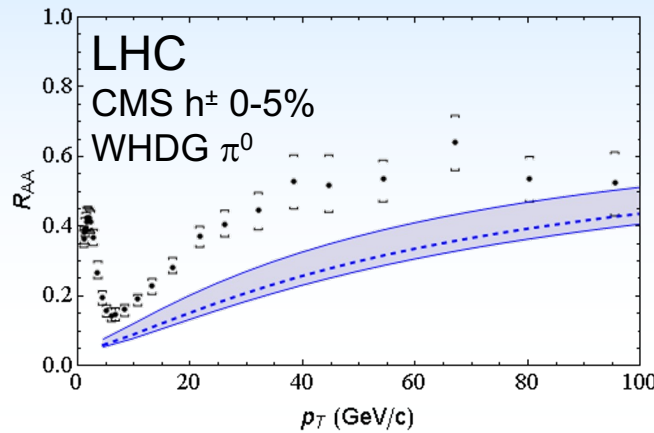
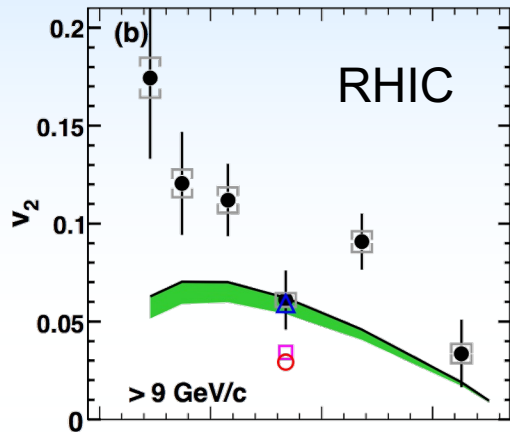
WAH and Du Plessis, PRD105 (2022) [2203.01259]

Faraday and WAH, in prep



Qual. Success of “LO” Jet Tomo in AA

- Rad + EI, realistic geom AA correct to factor ~ 2



PHENIX PRL105 (2010)

CMS, Eur.Phys.J. C72 (2012)
CMS, PRL109 (2012)

ALICE, JHEP1209 (2012) 112
CMS, JHEP 1205 (2012) 063



3/28/23

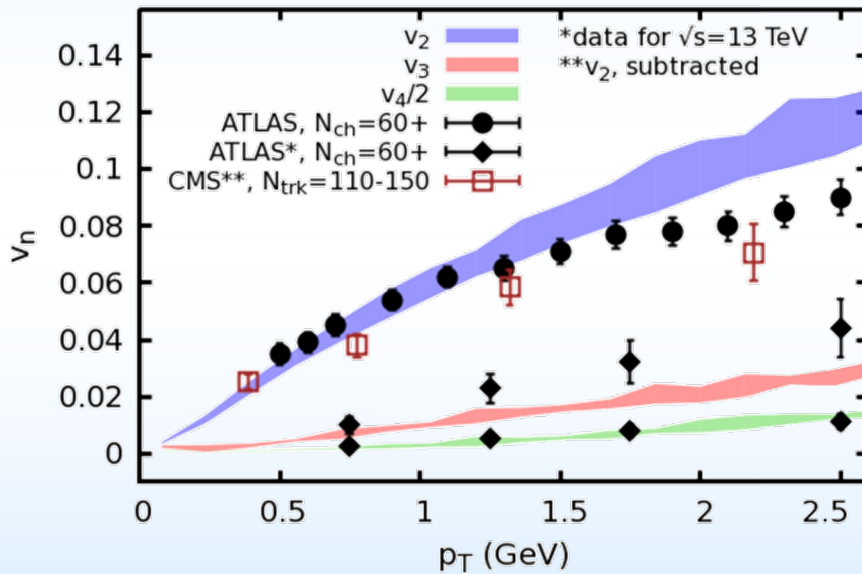
HP2023

2

Correlations and Suppression in Small Systems

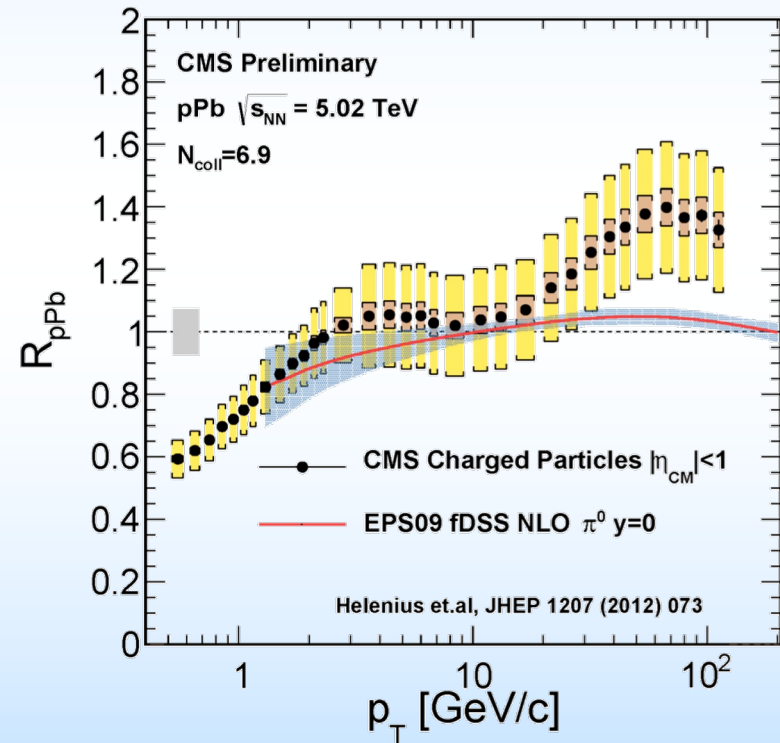
- Flow (?) in p+p

superSONIC for p+p, $\sqrt{s}=5.02$ TeV, 0-1%



Weller and Romatschke, PLB B 774 (2017)

- E-loss (?) in pPb



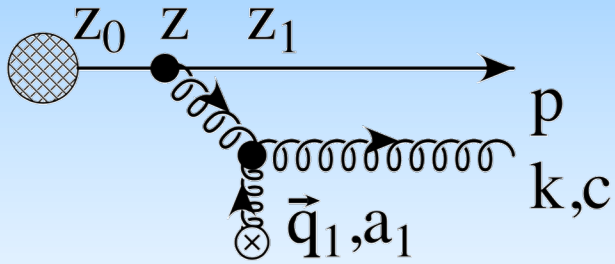
CMS, HP'13



pQCD E-Loss in pA

- Take seriously possibility of E-loss in small systems
- Wish to apply DGLV to small systems
 - However, DGLV assumes an ordering of scales:
 $1/\mu_{\text{Debye}} \ll \lambda_{\text{mfp}} \ll \tau_{\text{form}} \ll L_{\text{pathlength}}$
- Desire: derive the 1/L corrections to DGLV 1st order in opacity
 - Requires reevaluation of all 11 diagrams, no longer neglecting certain poles originating from the Yukawa scattering potential
 - Calculated by Isobel Kolbé





Neglected Poles

$$M_{1,0,1} = \int \frac{d^4 q_1}{(2\pi)^4} iJ(p+k-q_1) e^{i(p+k-q_1)x_0} \Lambda_1(p, k, q_1) V(q_1) e^{iq_1 x_1} \times \\ \times i\Delta_M(p+k-q_1) (-i)\Delta_{m_g}(k-q_1)$$

$$\text{Res}(\bar{q}_1) \approx -v(-\omega_0 - \tilde{\omega}_m, \mathbf{q}_1) \frac{e^{i(\omega_0 + \tilde{\omega}_m)(z_1 - z_0)}}{E^+ k^+ (\omega_1 + \tilde{\omega}_m)}$$

$$\text{Res}(\bar{q}_2) \approx v(\omega_1 - \omega_0, \mathbf{q}_1) \frac{e^{i(\omega_0 - \omega_1)(z_1 - z_0)}}{E^+ k^+ (\omega_1 + \tilde{\omega}_m)}$$

$$\text{Res}(-i\mu_1) \approx \frac{4\pi\alpha_s e^{-\mu_1(z_1 - z_0)}}{(-2i\mu_1)E^+ k^+ (-i\mu_1)^2}$$



Summed, Squared Result

$$\begin{aligned}
 \Delta E_{ind}^{(1)} = & \frac{C_R \alpha_s L E}{\pi \lambda} \int \frac{d^2 \mathbf{q}_1}{\pi} \frac{\mu^2}{(\mu^2 + \mathbf{q}_1^2)^2} \frac{d^2 \mathbf{k}}{4\pi} \int d\Delta z \bar{\rho}(\Delta z) \times \\
 & \times \left[-2 \frac{(\mathbf{k} - \mathbf{q}_1)}{(\mathbf{k} - \mathbf{q}_1)^2 + M^2 x^2 + m_g^2} \times (1 - \cos\{(\omega_1 + \tilde{\omega}_m)\Delta z\}) \right. \\
 & \times \left(\frac{\mathbf{k}}{m_g^2 + \mathbf{k}^2 + x^2 M^2} - \frac{(\mathbf{k} - \mathbf{q}_1)}{(\mathbf{k} - \mathbf{q}_1)^2 + M^2 x^2 + m_g^2} \right) \\
 & + \frac{1}{2} e^{-\mu \Delta z} \left\{ \left(\frac{\mathbf{k}}{m_g^2 + \mathbf{k}^2 + x^2 M^2} \right)^2 \times \right. \\
 & \times \left(1 - \frac{2C_R}{C_A} \right) \left(1 - \cos\{(\omega_0 - \tilde{\omega}_m)\Delta z\} \right) \\
 & + \frac{\mathbf{k}}{m_g^2 + \mathbf{k}^2 + x^2 M^2} \cdot \frac{(\mathbf{k} - \mathbf{q}_1)}{(\mathbf{k} - \mathbf{q}_1)^2 + M^2 x^2 + m_g^2} \times \\
 & \left. \left. \times \left(\cos\{(\omega_0 - \tilde{\omega}_m)\Delta z\} - \cos\{(\omega_0 - \omega_1)\Delta z\} \right) \right\} \right]
 \end{aligned}$$

DGLV

“Small L”
Correction



Analyzing the Short Path Correction

- **Correction:**

- $\Rightarrow 0$ as $\Delta z \Rightarrow 0$
(destructive interference)

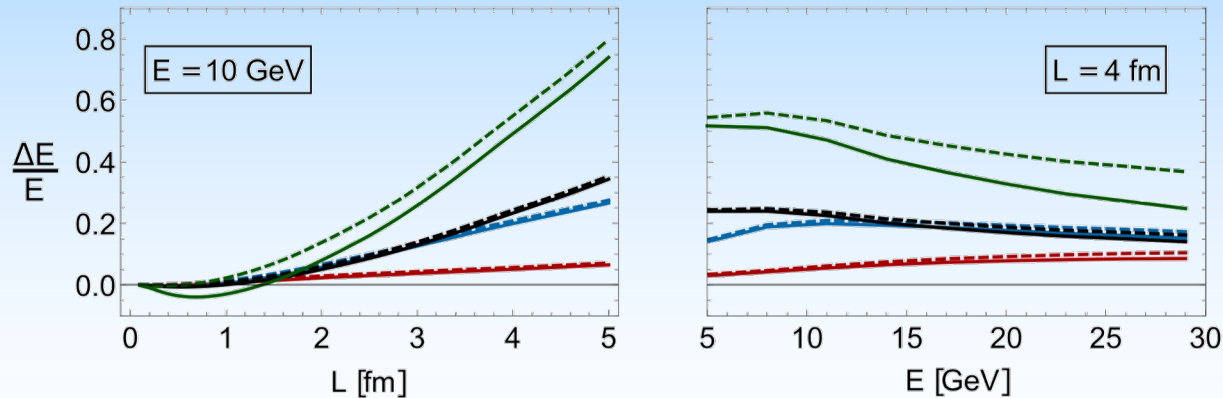
- $\Rightarrow 0$ as $\mu \Rightarrow \infty$
(all paths are long)

- breaks color triviality

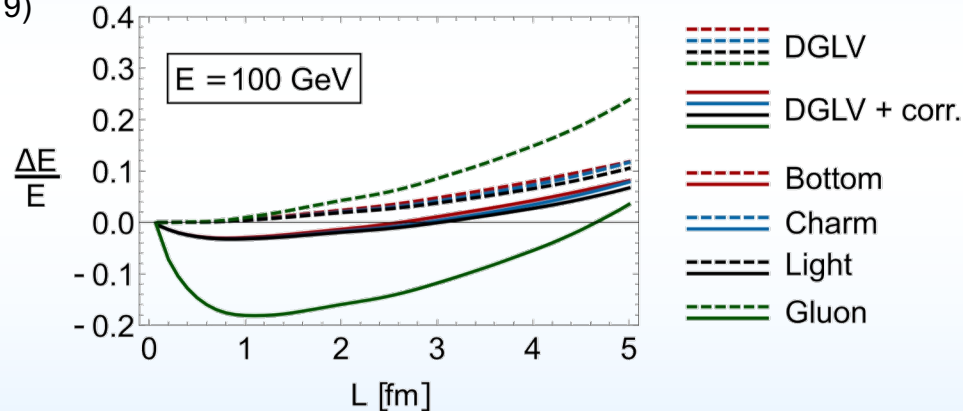
$$\begin{aligned} \Delta E_{ind}^{(1)} = & \frac{C_R \alpha_s L E}{\pi \lambda} \int \frac{d^2 \mathbf{q}_1}{\pi} \frac{\mu^2}{(\mu^2 + \mathbf{q}_1^2)^2} \frac{d^2 \mathbf{k}}{4\pi} \int d\Delta z \bar{\rho}(\Delta z) \times \\ & \times \left[-2 \frac{(\mathbf{k} - \mathbf{q}_1)}{(\mathbf{k} - \mathbf{q}_1)^2 + M^2 x^2 + m_g^2} \times (1 - \cos\{(\omega_1 + \tilde{\omega}_m)\Delta z\}) \right. \\ & \times \left(\frac{\mathbf{k}}{m_g^2 + \mathbf{k}^2 + x^2 M^2} - \frac{(\mathbf{k} - \mathbf{q}_1)}{(\mathbf{k} - \mathbf{q}_1)^2 + M^2 x^2 + m_g^2} \right) \\ & + \frac{1}{2} e^{-\mu \Delta z} \left\{ \left(\frac{\mathbf{k}}{m_g^2 + \mathbf{k}^2 + x^2 M^2} \right)^2 \times \right. \\ & \times \left(1 - \frac{2C_R}{C_A} \right) \left(1 - \cos\{(\omega_0 - \tilde{\omega}_m)\Delta z\} \right) \\ & + \frac{\mathbf{k}}{m_g^2 + \mathbf{k}^2 + x^2 M^2} \cdot \frac{(\mathbf{k} - \mathbf{q}_1)}{(\mathbf{k} - \mathbf{q}_1)^2 + M^2 x^2 + m_g^2} \times \\ & \left. \left. \times (\cos\{(\omega_0 - \tilde{\omega}_m)\Delta z\} - \cos\{(\omega_0 - \omega_1)\Delta z\}) \right\} \right] \end{aligned}$$



Num. Investigation of Correction



Kolbé and WAH, PRC100 (2019)



- Surprise 1: Correction leads to reduction in E-loss
- Surprise 2: Affects *all* pathlengths L
 - Due to integrating over all distances to scattering Δz in $[0, L]$
- Surprise 3: Correction grows with p_T
- Need to think more carefully re small systems?



Pheno Investigation

- Reasons to be excited:
 - Destructive interference $\Rightarrow R_{pA} > 1$ similar to data?
 - E (vs $\ln E$) dependence \Rightarrow faster rise in $R_{AA}(p_T)$?
- Modest ambition:
 - How important is short L corr. vs original WHDG in AA? pA ? pp ?
- Serious comparison to data? Caveats:
 - Literally hot off the presses (plots from $>1am$ today)
 - No nPDFs, no NLO, no NLL, no Bayesian, no new parameter extraction, mapping of hydro to bricks, time dependence of expansion, KKP/DSS π FFs not under full control for large p_T , limited by publicly available hydro backgrounds
- Seek **qualitative** comparison; is correction important to include? Other important previously neglected physics?



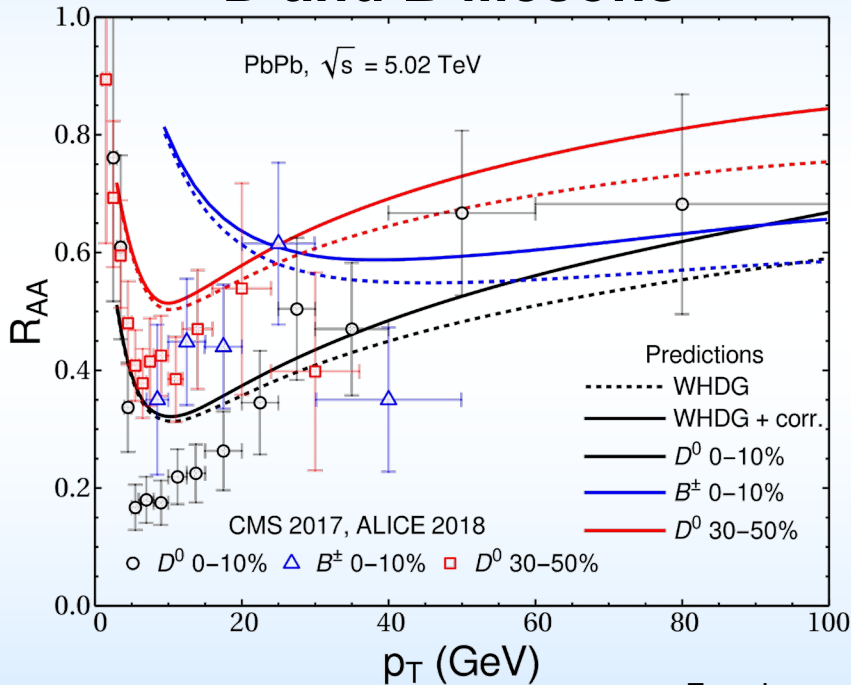
Cole Faraday

Effect on R_{AA}

- Very preliminary first results:

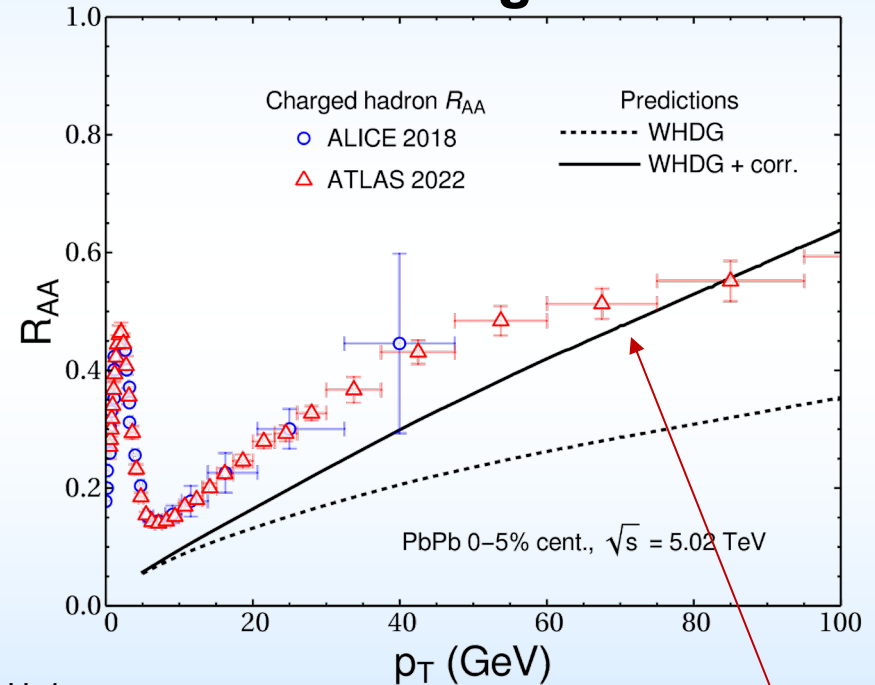
- Full Rad + Coll + Geom (fluc. IP-G into VISHNU)

- **D and B Mesons**



Faraday and WAH, in prep

- **Pions vs Charged Hadrons**



FFs becoming questionable

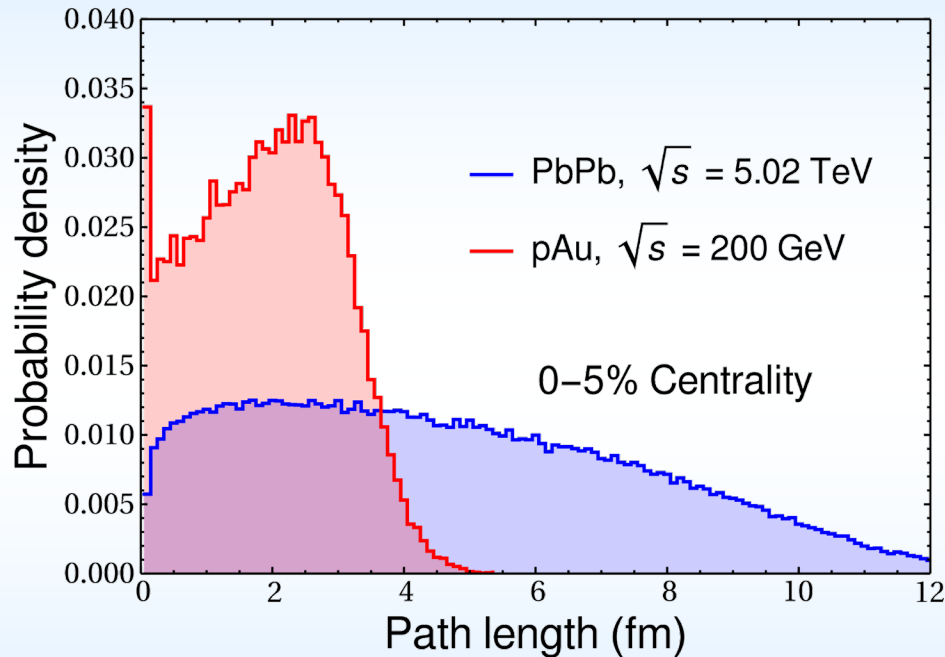
- Short path length correction to ΔE_{rad} :

- reduces suppression
- increases energy dependence

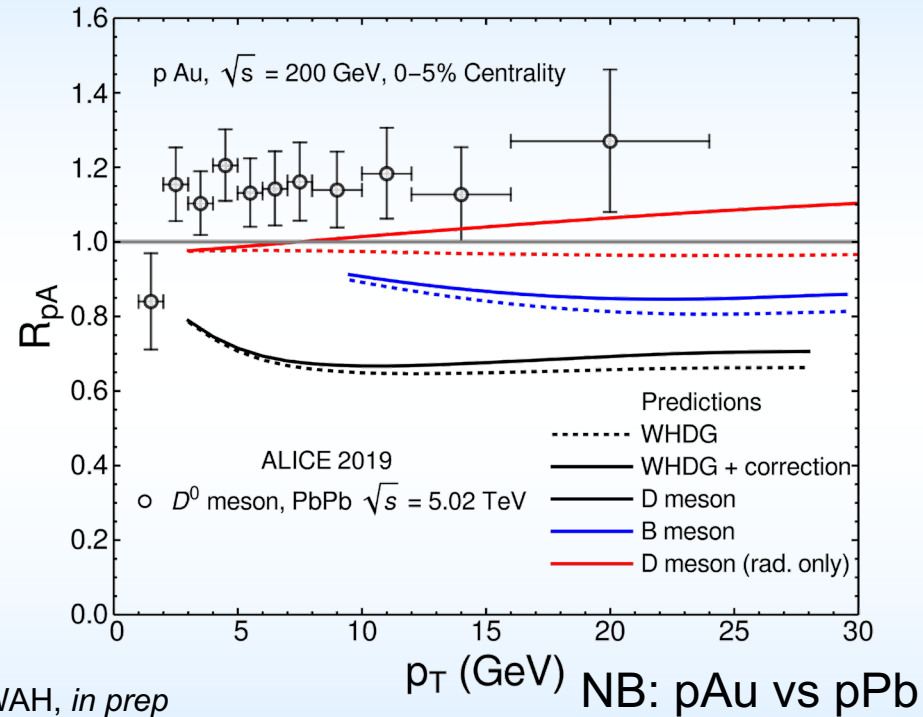


Effect on R_{pA}

- Non-trivial pathlengths in pA collisions



Faraday and WAH, *in prep*



– Average elastic energy loss inappropriate for small systems

– Tantalizing hints of $R_{pA} > 1$ and rising

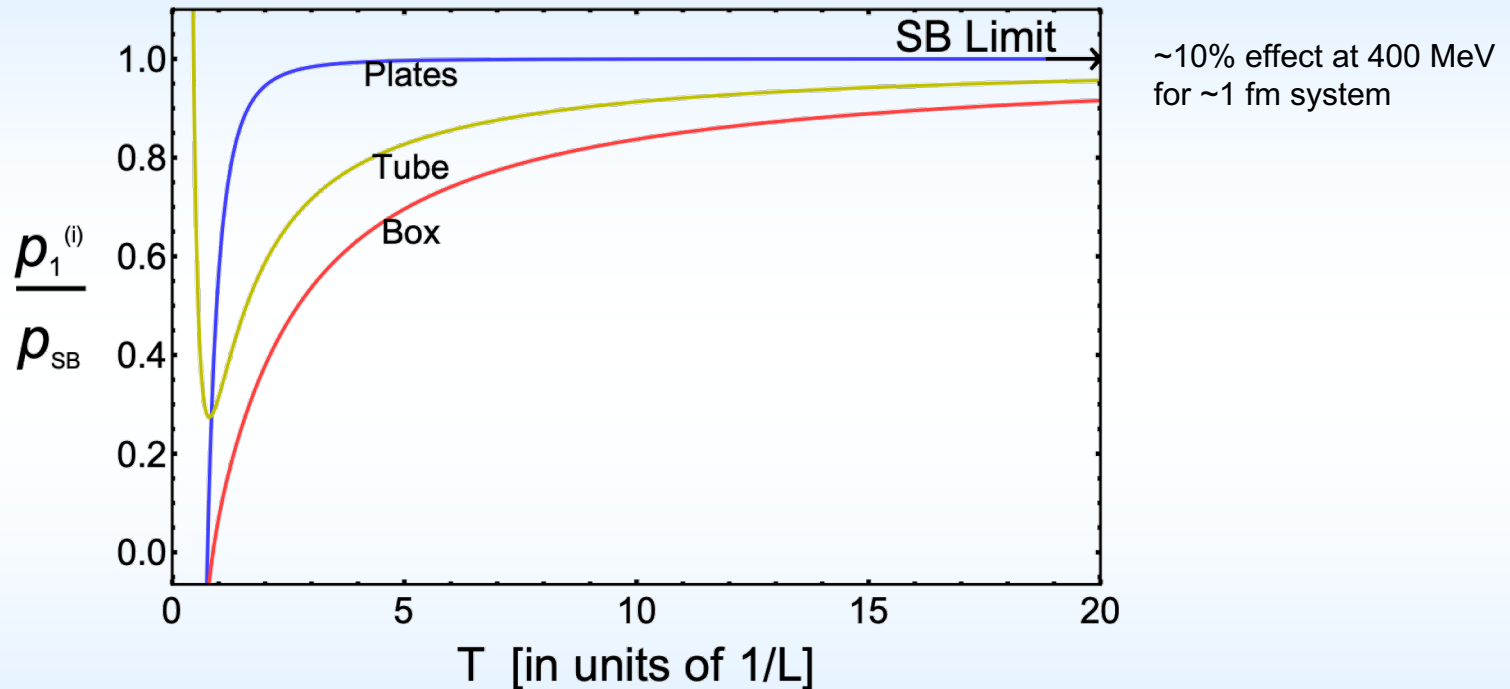


Thermodynamics of Small Systems



Does Finite Size Affect Thermodyn.?

- Test using free scalar field theory



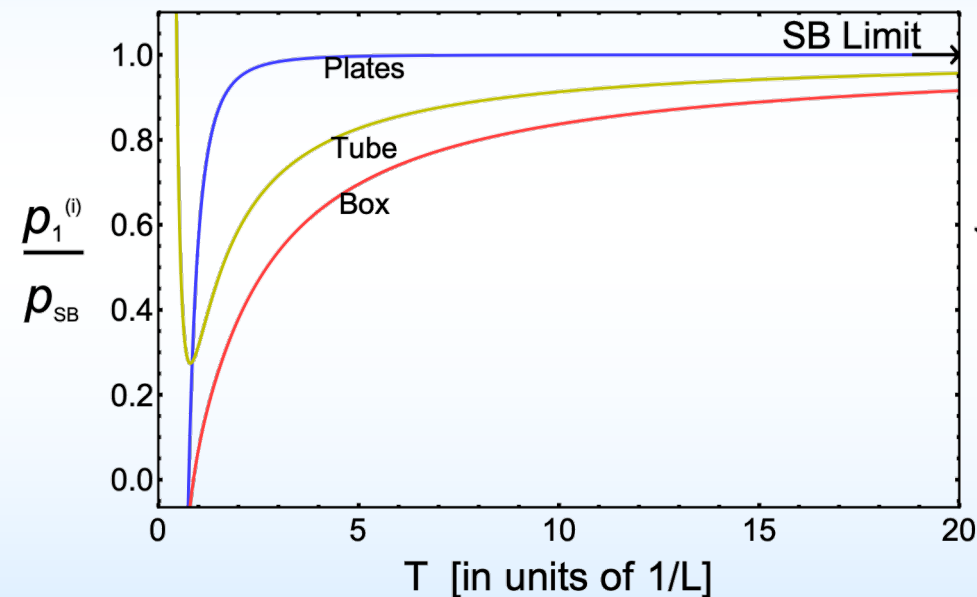
Mogliacci, Kolbé, and WAH, PRD102 (2020)

- p decreases significantly as T decreases for fixed L , converging to $T = 0$ Casimir effect

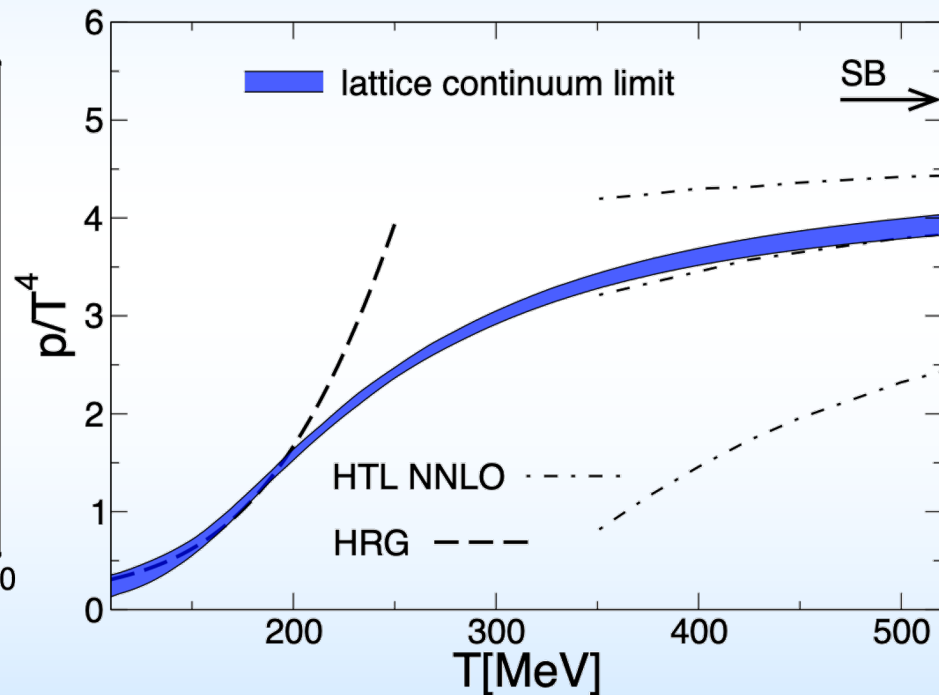


Does Finite Size Affect Thermodyn.?

- Provocative qualitative T dependence:



Mogliacci, Kolbé, and WAH, PRD102 (2020)

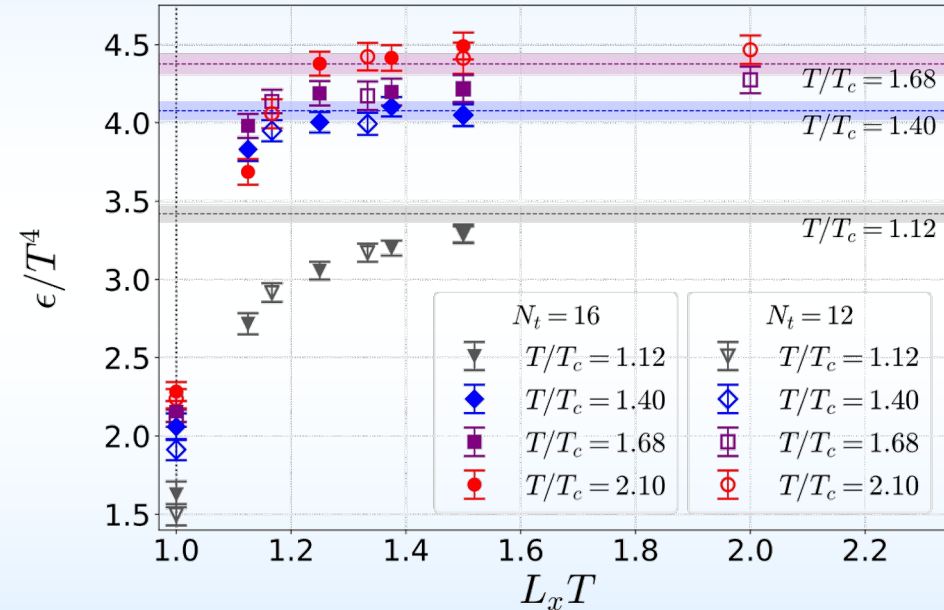
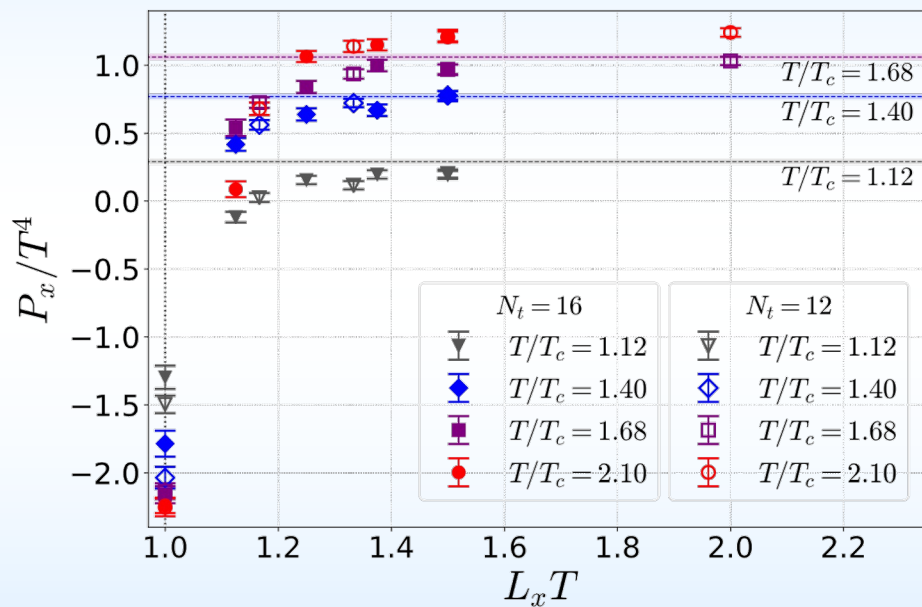


Borsanyi et al., PLB730 (2014)



FS Effects in Quenched QCD

- Thermodynamic quantities on anisotropic lattice with periodic boundary conditions



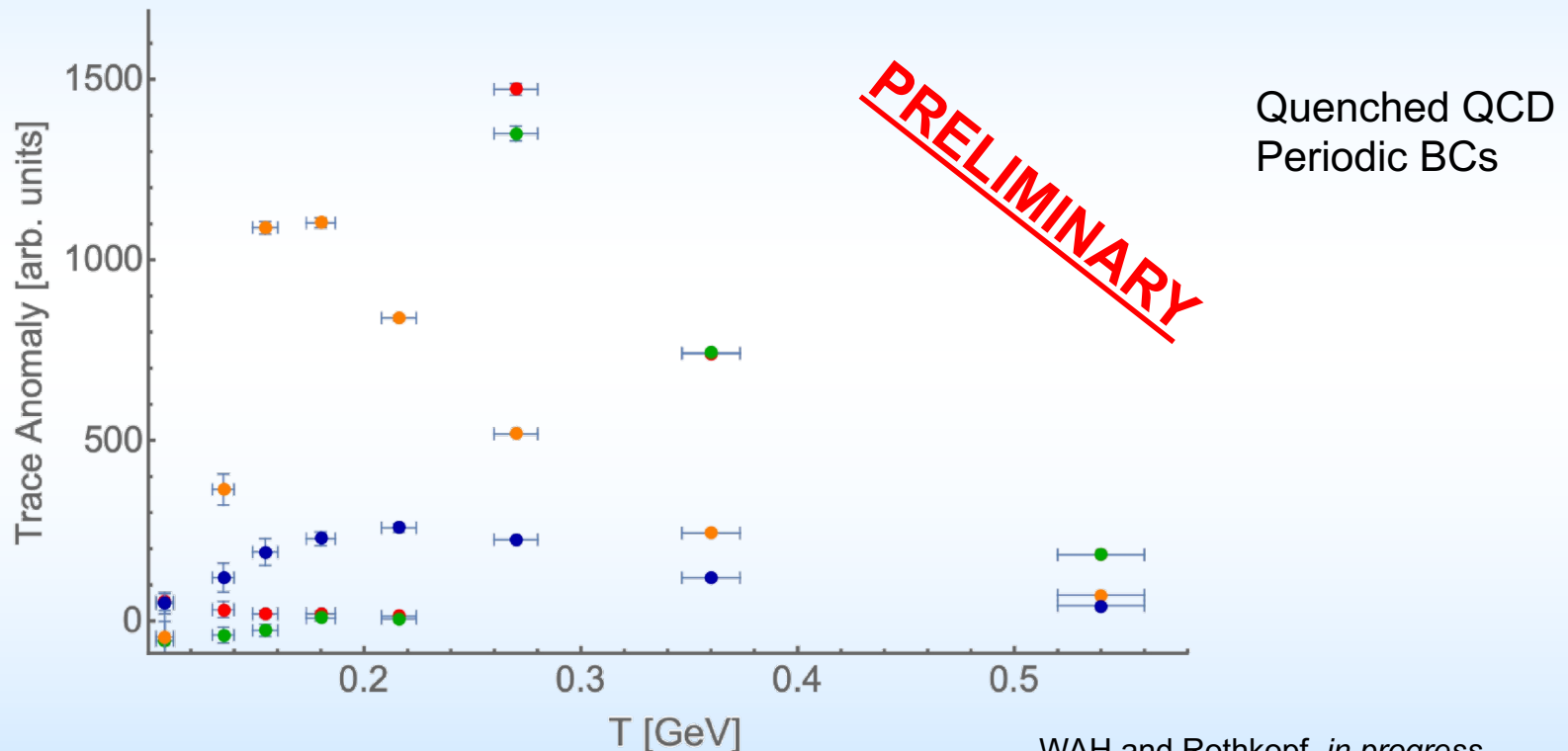
Kitazawa, Mogliacci, Kolbé, and WAH, PRD99 (2019)

– Reduction of p, e with T^*L qualitatively similar to free, massless scalar theory



What about Trace Anomaly Δ ?

- $\Delta \Rightarrow c_s \Rightarrow \eta/s$
- From lattice QCD:



– Decreasing system size:

- decreases Δ
- washes out phase transition



Analytic Results for Δ

- $\Delta = 0$ for massless, free scalar theory

- Even when in a box!

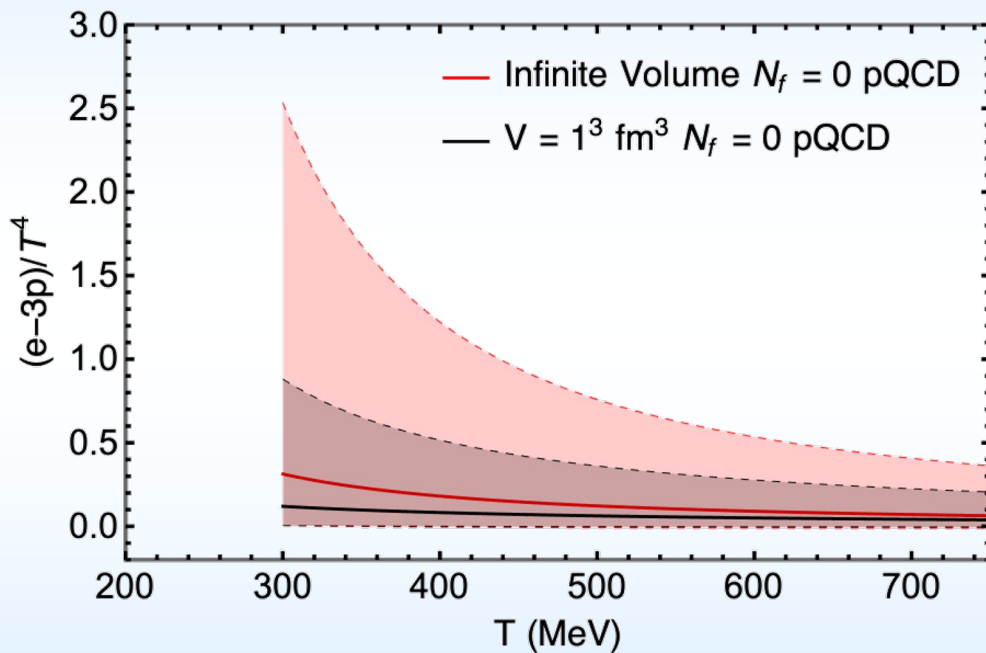
- $\Delta = 0$ for HTL QCD

- Even to g^5 !

- $\Delta \neq 0$ *only* when coupling runs

- Estimate: use HTL QCD with lattice scalar running coupling

- λ decreases and Δ *significantly* decreases from F.S.



WAH and Rothkopf, *in progress*



Analytic FS Δ in QCD

- What happens in QCD?
 - Does α_s grow or shrink as $L \Rightarrow 0$?
- Non-trivial conceptual issues:
 - How to regularize and renormalize?
 - Dim reg difficult to generalize in F.S. setting
 - Torons
- First examine F.S. effects for running coupling in $2 \Rightarrow 2$ scattering in ϕ^4 theory



Define New Regularization Scheme

- Denominator regularization (den reg) instead of dimensional regularization (dim reg)
 - **Number of dimensions fixed**
 - Feynman x combine propagators
 - Analytically continue the power of the single denominator
 - Introduce fictitious scale μ to maintain dimensions

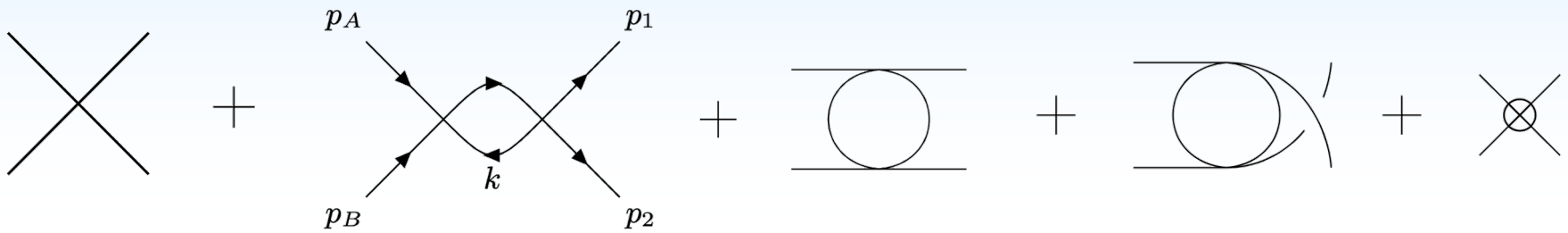
$$V(p^2; \mu, \epsilon) = -\frac{1}{2} \int_0^1 dx \int \frac{d^{4-\epsilon} k}{(2\pi)^{4-\epsilon}} \frac{i\mu^\epsilon}{(k^2 - \Delta^2)^2}$$


Use instead: $\Rightarrow -\frac{1}{2} \int_0^1 dx \int \frac{d^4 k}{(2\pi)^4} \frac{i(-\mu^2)^{2\epsilon}}{(k^2 - \Delta^2)^{2+\epsilon}}$



2 => 2 at NLO in ϕ^4 in Finite System

- Feynman Diagrams:**



- Define** $(-i\lambda)^2 iV(p^2) \equiv$ 

- Impose Periodic B.C.'s**

$$V(p^2, \{L_i\}; \mu, \epsilon) = -\frac{1}{2} \int_0^1 dx \int \frac{dk^0}{2\pi} \sum_{\vec{k} \in \mathbb{Z}^3} \frac{1}{(2\pi)^3 L_1 L_2 L_3} \frac{\mu^{2\epsilon}}{[k^2 - \Delta^2]^{2+\epsilon}}$$



Capture the Divergence

- Result is a generalized Epstein Zeta fcn

$$V(p^2, \{L_i\}; \mu, \epsilon) = -\frac{1}{2} \frac{1}{2\pi} \frac{1}{(2\pi)^3 L_1 L_2 L_3} \frac{\sqrt{\pi} \Gamma(\frac{3}{2} + \epsilon)}{\Gamma(2 + \epsilon)} \int_0^1 dx \sum_{\vec{k} \in \mathbb{Z}^3} \frac{\mu^{2\epsilon}}{\left(\sum_{i=1}^3 \left(\frac{k^i}{L_i} + x p^i\right)^2 + \Delta^2\right)^{\frac{3}{2} + \epsilon}}$$

- Poisson Summation Formula: $\sum_{\vec{n} \in \mathbb{Z}^p} f(\vec{n}) = \sum_{\vec{m} \in \mathbb{Z}^p} \tilde{F}(\vec{m})$
- Yields new analytic continuation for g.E.Z:

$$\sum_{\vec{n} \in \mathbb{Z}^p} (a_i^2 n_i^2 + b_i n_i + c - i\epsilon)^{-s} = \frac{1}{a_1 \cdots a_p} \frac{1}{\Gamma(s)} \left[\pi^{p/2} \Gamma\left(s - \frac{p}{2}\right) \left(c - \sum \frac{b_i^2}{4a_i^2} - i\epsilon\right)^{\frac{p}{2} - s} \right. \\ \left. + 2\pi^s \sum'_{\vec{m} \in \mathbb{Z}^p} e^{-2\pi i \sum \frac{m_i b_i}{2a_i^2}} \left(\frac{c - \sum \frac{b_i^2}{4a_i^2} - i\epsilon}{\sum \frac{m_i^2}{a_i^2}}\right)^{\frac{p}{4} - \frac{s}{2}} K_{s - \frac{p}{2}} \left(2\pi \sqrt{\left(c - \sum \frac{b_i^2}{4a_i^2} - i\epsilon\right) \left(\sum \frac{m_i^2}{a_i^2}\right)}\right) \right]$$



Finite Size Result at NLO

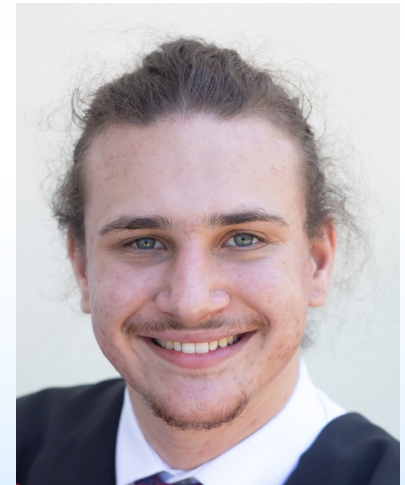
- Found the pole! And the F.S. correction!

$$V(p^2, \{L_i\}; \mu, \epsilon) = -\frac{1}{2} \frac{1}{(4\pi)^2} \int_0^1 dx \left\{ \frac{1}{\epsilon} - 1 + \ln \frac{\mu^2}{\Delta^2} + 2 \sum'_{\vec{m} \in \mathbb{Z}^3} e^{-2\pi i x \sum m_i p^i L_i} K_0 \left(2\pi |\Delta| \sqrt{\sum m_i^2 L_i^2} \right) \right\}$$

– Correction

- goes to 0 as $L_i, p \Rightarrow$ infinity
- satisfies unitarity/optical theorem
- Optical thm check highly non-trivial

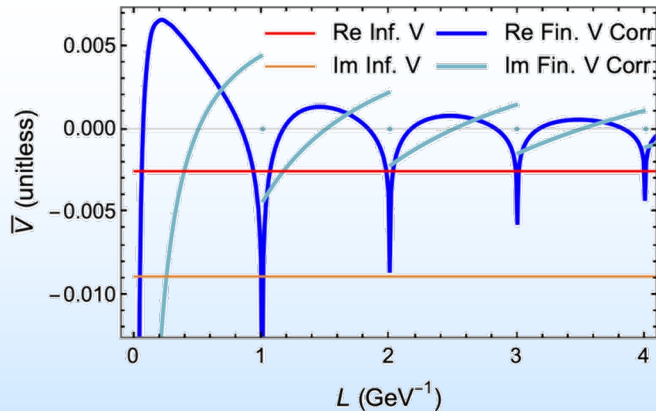
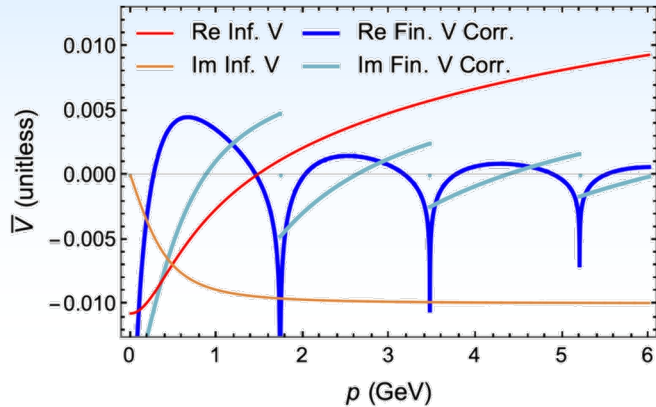
– Requires generalization of a number theory result from Hardy/Ramanujan



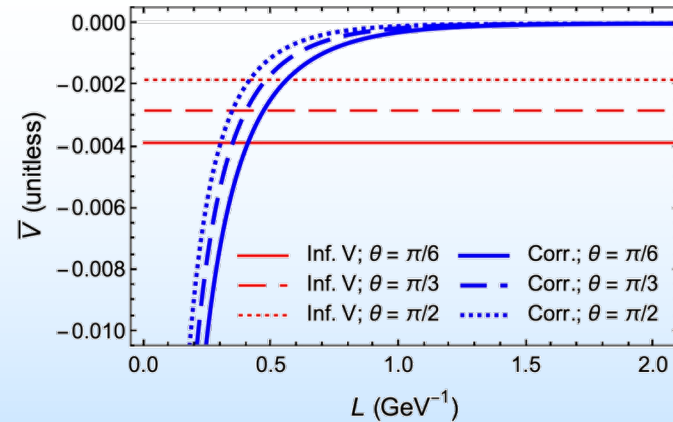
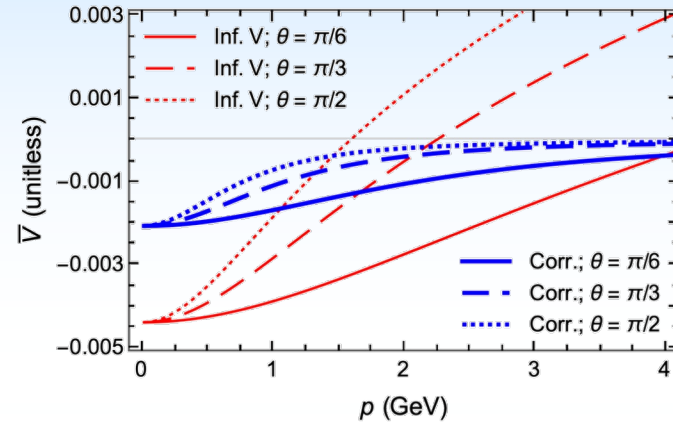
Jean du Plessis

Numerical Results with 1 Compact D

- s channel



- t channel



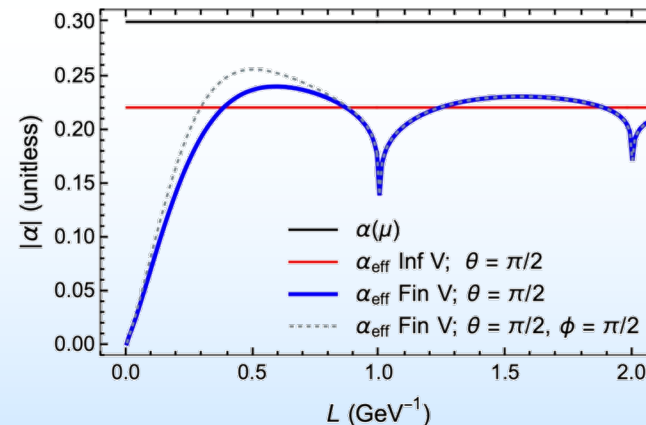
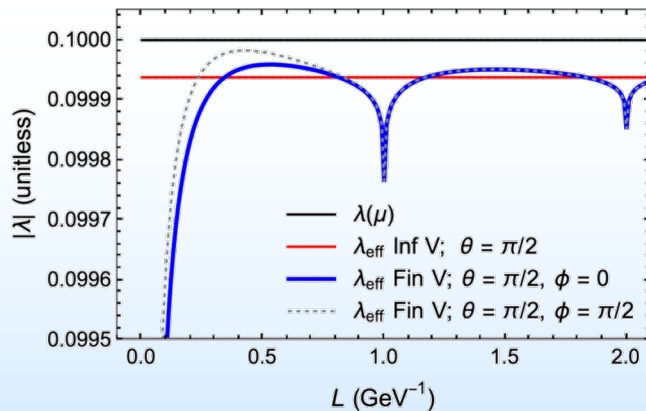
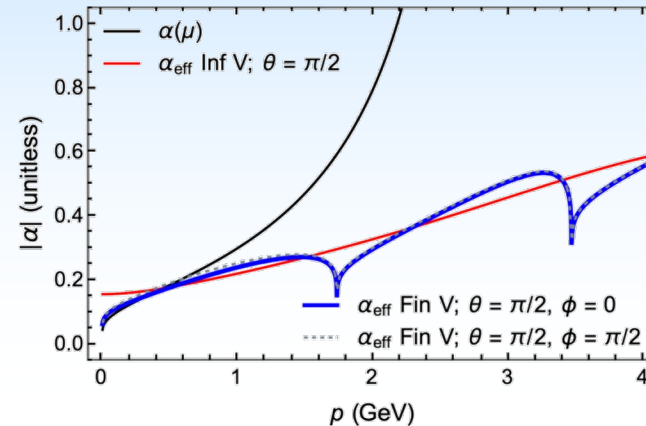
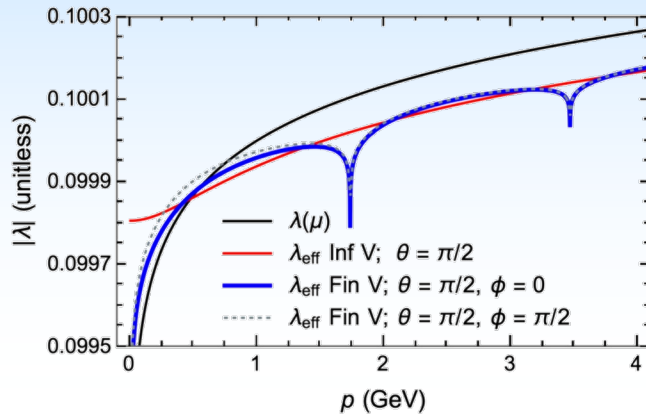
– Correction $\Rightarrow 0$ in L and p

– Non-trivial behavior from x int over BesselK



Running Coupling in 1 Compact D

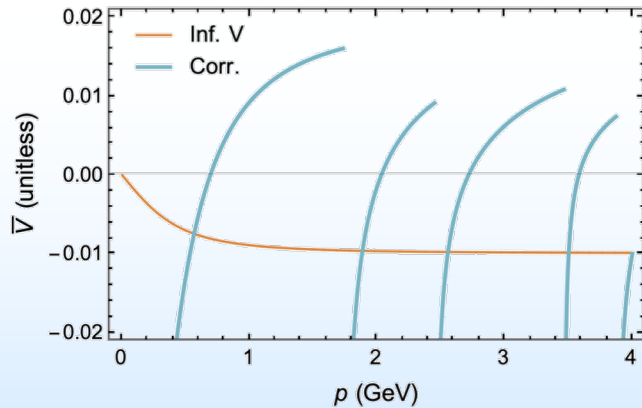
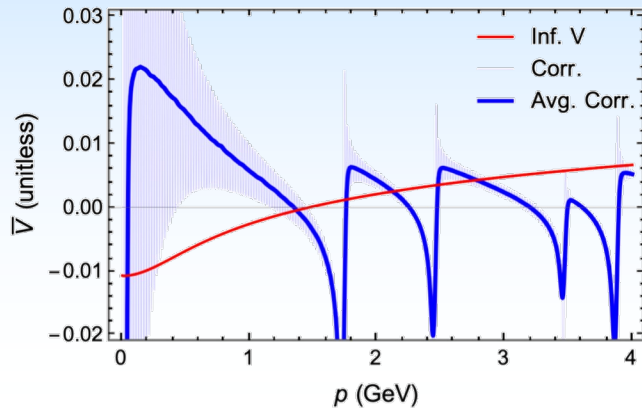
– Sum the geometric series of bubbles



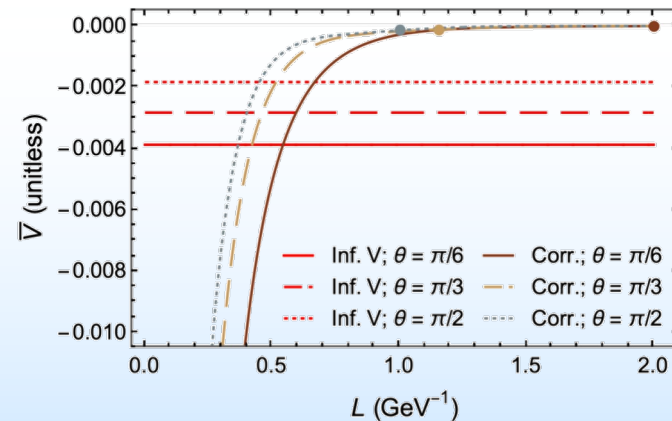
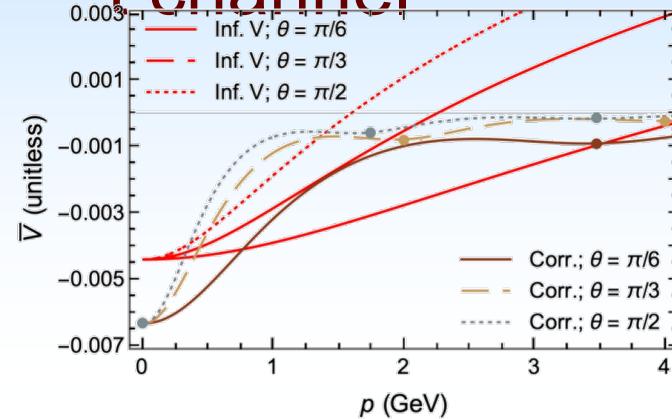
- Result depends on coupling at initial scale
- No pole in full series; cf LL from beta fcn

Numerical Results with 2 Compact D's

- s channel



- t channel

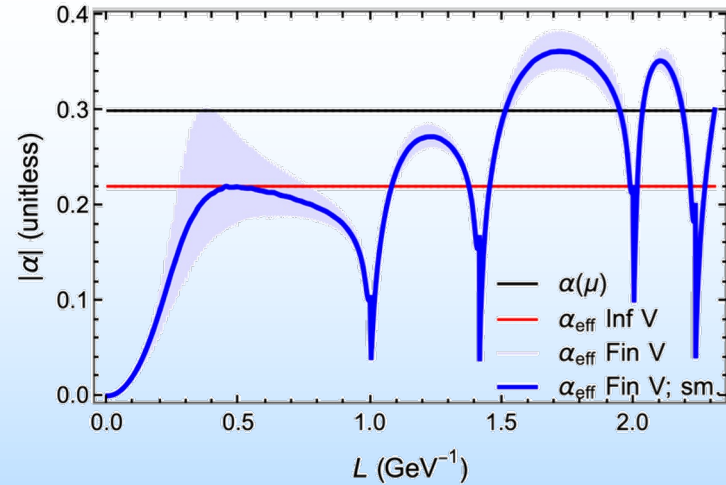
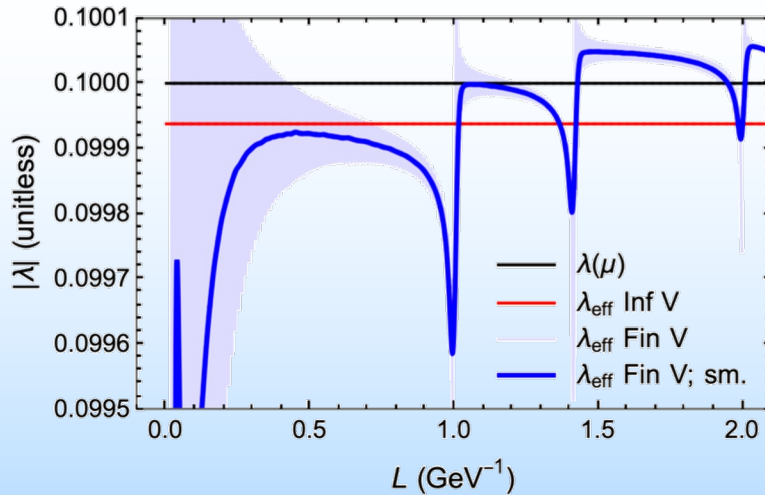
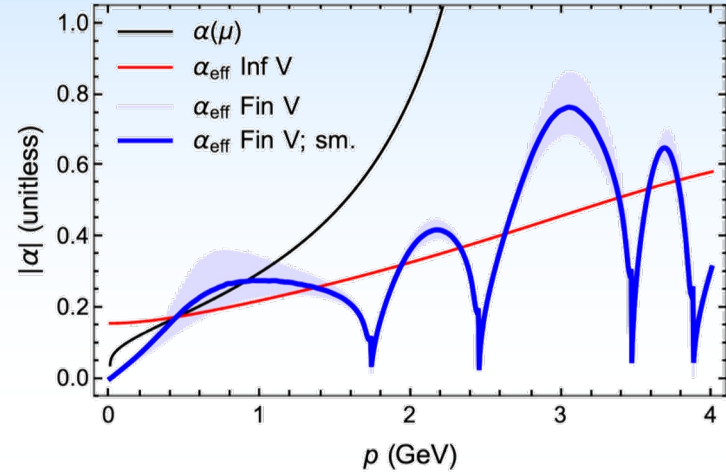
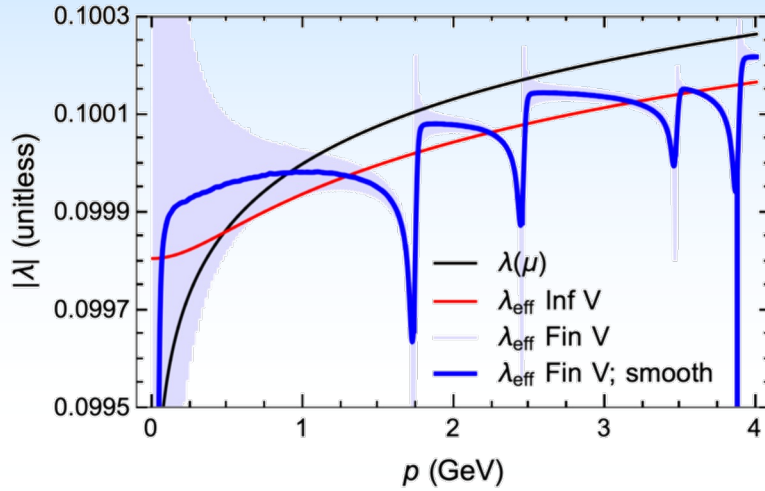


– Extremely difficult numerical problem for Re of s channel

– 3D compact is trivial



Running Coupling in 2 Compact D's



Conclusions

- Presented first (*and only*) analytic small L correction to E -loss derivation
 - Correction grows w $1/L$, E ; breaks color triviality
- First pheno results. Small L correction:
 - Reduces suppression
 - Increases rise with p_T
 - $\sim 10\%$ effect for heavy flavor R_{AA}
 - Broken color triviality \Rightarrow large effect for light hadrons
 - Significant pathlengths in pA; hints of $R_{pA} > 1$; **must** compute small pathlength corrections for elastic energy loss
- To do: see caveats; then v_n 's for AA and pA, high multiplicity pp, ...



Bonus Conclusions

- Massive, free scalar field in a box
 - Thermodynamics significantly altered; mimics QCD
 - Quenched lattice confirms qualitative physics
- NLO $2 \Rightarrow 2$ scattering in ϕ^4 in 1, 2, and 3 compact dim's
 - Analytic continuation of the generalized Epstein zeta function
 - Captured finite size corrections
 - Checked unitarity
 - Showed first results for running coupling in 1, 2, and 3 compact dim's
- A lot of interesting work to do!!



Very Preliminary v_n 's

